

**University of Kentucky**  
**Department of Physics and Astronomy**  
**PHY 525: Solid State Physics II**  
**Fall 2000**  
**Test 2**

Date: November 8, 2000 (Wednesday)

Time allowed: 50 minutes.

Answer all questions.

1. **Hall coefficient of a semiconductor.** (a) Starting from drift velocity approximation, show that the Hall coefficient of a semiconductor is given by

$$R_H = \frac{1}{n_i e c} \frac{1-b}{1+b}$$

where  $n_i$  is the intrinsic concentration, and  $b$  is the ratio of mobility in conduction and valance bands ( $b = \mu_n / \mu_p$ ). (b) Consider two semiconductors A and B of with same energy gap. The following table gives the ratio of electron effective mass and relaxation time in the conduction and valance bands:

	$m_n^*/m_p^*$	$\tau_n/\tau_p$
Semiconductor A	0.3	30
Semiconductor B	0.8	15

Estimate the ratio of their Hall coefficients from Drude model. (c) For a semiconductor of energy gap equivalent to a temperature of 12500K, estimate the ratio of the Hall coefficient at 300K and 200K,  $R_H(300K)/R_H(200K)$ . Assume  $\tau$  does not vary much in this temperature range.

Solution:

$$(a) \quad \hbar \left( \frac{d}{dt} + \frac{1}{\tau} \right) \vec{k} = q \left( \vec{E} + \frac{1}{c} \vec{v} \times \vec{B} \right) \quad \text{--- (0)}$$

Assume steady case,  $\frac{d\vec{v}}{dt} = 0$ .

Above equations become :

$$(0) \Rightarrow v_x = \frac{q\tau}{m} E_x + \frac{q\tau B}{mc} v_y = \begin{cases} \mu_p E_x + \frac{\mu_p B}{c} v_{py} \\ -\mu_n E_x - \frac{\mu_n B}{c} v_{ny} \end{cases} \quad \text{--- (1)}$$

$$(0) \Rightarrow v_y = \frac{q\tau}{m} E_y - \frac{q\tau B}{mc} v_x = \begin{cases} \mu_p E_y - \frac{\mu_p B}{c} v_{px} \\ -\mu_n E_y + \frac{\mu_n B}{c} v_{nx} \end{cases} \quad \text{--- (2)}$$

$$v_y = 0 \Rightarrow n_i e v_{py} - n_i e v_{ny} = 0 \text{ and } j = j_x = n_i e v_{px} - n_i e v_{nx} :$$

$$(1) \Rightarrow \begin{cases} v_{px} = \mu_p E_x \\ v_{nx} = -\mu_n E_x \end{cases} \quad \text{--- (3)}$$

$$\Rightarrow \begin{cases} j_p = n_i e \mu_p E_x \\ j_n = n_i e \mu_n E_x \end{cases} \\ \Rightarrow j = n_i e (\mu_p + \mu_n) E_x \quad \text{--- (4)}$$

$$(3) \& (4) \Rightarrow \begin{cases} v_{px} = \mu_p \frac{j}{n_i e (\mu_p + \mu_n)} \\ v_{nx} = -\mu_n \frac{j}{n_i e (\mu_p + \mu_n)} \end{cases} \quad \text{--- (5)}$$

$$(2) \Rightarrow n_i e \left( \mu_p E_y - \frac{\mu_p B}{c} v_{px} \right) - n_i e \left( -\mu_n E_y + \frac{\mu_n B}{c} v_{nx} \right) = 0 \quad (\because n_i e v_{py} - n_i e v_{ny} = 0)$$

$$\Rightarrow n_i e (\mu_p + \mu_n) E_y - \frac{n_i e B}{c} (\mu_p v_{px} + \mu_n v_{nx}) = 0$$

$$\Rightarrow n_i e (\mu_p + \mu_n) E_y - \frac{n_i e B}{c} \left( \mu_p \mu_p \frac{j}{n_i e (\mu_p + \mu_n)} - \mu_n \mu_n \frac{j}{n_i e (\mu_p + \mu_n)} \right) = 0 \quad (\text{from(4)})$$

$$\Rightarrow E_y - \frac{Bj}{c n_i e (\mu_p + \mu_n)^2} (\mu_p^2 - \mu_n^2) = 0$$

$$\Rightarrow E_y = \frac{Bj(\mu_p - \mu_n)}{c n_i e (\mu_p + \mu_n)}$$

$$\Rightarrow E_y = \frac{Bj(1-b)}{c n_i e (1+b)} \quad (b = \frac{\mu_n}{\mu_p})$$

$$\therefore R_H = \frac{E_y}{Bj} = \frac{1}{n_i e c} \frac{1-b}{1+b}$$

(b)

Since the energy gap is the same for both semiconductors,  $\therefore n_i$  is the same for both A and B.

$$\begin{aligned} \therefore \frac{(R_H)_A}{(R_H)_B} &= \frac{\left(\frac{1-b_A}{1+b_A}\right)}{\left(\frac{1-b_B}{1+b_B}\right)} \\ \mu &= \left| \frac{e\tau}{m} \right| \propto \frac{\tau}{m} \Rightarrow b = \frac{\mu_n}{\mu_p} = \frac{\tau_n}{m_n} / \frac{\tau_p}{m_p} = \frac{(\tau_n / \tau_p)}{(m_n / m_p)} \\ \therefore \left. \begin{aligned} b_A &= \frac{30}{0.3} = 100 \\ b_B &= \frac{15}{0.8} = 18.75 \end{aligned} \right\} \Rightarrow \frac{(R_H)_A}{(R_H)_B} &= \frac{\left(\frac{1-100}{1+100}\right)}{\left(\frac{1-18.75}{1+18.75}\right)} = \frac{0.9802}{0.8987} = \underline{\underline{1.09}} \end{aligned}$$

(c)

$$R_H = \frac{1}{n_i e c} \frac{1-b}{1+b} \propto n_i^{-1} \quad (\text{assuming } b \text{ constant})$$

$$n_i^2 = np = N_C N_V \exp\left[-\frac{E_g}{k_B T}\right]$$

$$n_i = \sqrt{N_C N_V} \exp\left[-\frac{E_g}{2k_B T}\right]$$

Note that  $N_C$  and  $N_V \propto T^{3/2}$

$$\therefore n_i \propto T^{3/2} \exp\left[-\frac{E_g}{2k_B T}\right]$$

$$\begin{aligned} \frac{R_H(300K)}{R_H(200K)} &= \frac{n_i(200K)}{n_i(300K)} = \left(\frac{300}{200}\right)^{3/2} \frac{\exp\left[-\frac{12500k_B}{2k_B \times 200}\right]}{\exp\left[-\frac{12500k_B}{2k_B \times 300}\right]} \\ &= \left(\frac{300}{200}\right)^{3/2} \frac{2.681 \times 10^{-14}}{8.958 \times 10^{-10}} \\ &= \underline{\underline{5.5 \times 10^{-5}}} \end{aligned}$$

2. **Graded diode.** In the depletion layer (width  $2w$ ) of a graded p-n junction, the doping level varies linearly with position:  $N_D - N_A = kx$  for  $-w \leq x \leq +w$ . If the p- and n-semiconductors have doping level of  $N_A$  and  $N_D$  respectively. Let the effective density of state of the conduction and valance band be  $N_C$  and  $N_V$  respectively. Also let the energy gap of the intrinsic semiconductor be  $E_g$ . Assume all impurities are fully ionized. (a) Calculate chemical potential  $\mu_n$  and  $\mu_p$  in terms of  $N_D$ ,  $N_A$ ,  $N_C$ ,  $N_V$ , and  $E_g$ , from this calculate the potential difference  $\Delta\phi$  across the depletion layer. (b) What should be the electric field  $E$  at the two sides of the depletion layer? Find the electric field  $E(x)$  within the depletion layer. (c) Find the electric potential  $\phi(x)$  within the depletion layer. (d) Determine the layer width  $2w$  in term of  $\Delta\phi$ .

Solution:

(a)

Let us first consider the n - type semiconductor. Since the impurities are fully ionized.

$$\begin{aligned} \therefore n \approx N_D &= N_C \exp\left[-\frac{E_g - \mu_n}{k_B T}\right] \Rightarrow \mu_n - E_g = k_B T \ln\left[\frac{N_D}{N_C}\right] \\ &\Rightarrow \mu_n = E_g - k_B T \ln\left[\frac{N_C}{N_D}\right] \end{aligned}$$

Similarly, for the p - type semiconductor :

$$\begin{aligned} p \approx N_A &= N_V \exp\left[-\frac{\mu_p}{k_B T}\right] \Rightarrow -\mu_p = k_B T \ln\left[\frac{N_A}{N_V}\right] \\ &\Rightarrow \mu_p = k_B T \ln\left[\frac{N_V}{N_A}\right] \\ e\Delta\phi \approx \mu_n - \mu_p &= E_g - k_B T \ln\left[\frac{N_C}{N_D}\right] - k_B T \ln\left[\frac{N_V}{N_A}\right] \\ &= E_g - k_B T \ln\left[\frac{N_C N_V}{N_D N_A}\right] \end{aligned}$$

(b)

Electric field at  $x = \pm w$  should be zero.

Applied Gauss' law to the static charges at the depletion layer.

$$\epsilon \oint \vec{E} \cdot d\vec{S} = 4\pi Q_{\text{enclosed}} \Rightarrow \epsilon dEA = 4\pi \rho A dx$$

$$\Rightarrow dE = \frac{4\pi \epsilon k x dx}{\epsilon} \quad [\rho = e(N_D - N_A)]$$

$$\Rightarrow E = \frac{2\pi \epsilon k x^2}{\epsilon} + C$$

$$E \text{ at } x = \pm w \text{ is } 0 \Rightarrow \frac{2\pi \epsilon k w^2}{\epsilon} + C = 0$$

$$\Rightarrow C = -\frac{2\pi \epsilon k w^2}{\epsilon}$$

$$\therefore E = \frac{2\pi \epsilon k x^2}{\epsilon} - \frac{2\pi \epsilon k w^2}{\epsilon} = \frac{2\pi \epsilon k}{\epsilon} (x^2 - w^2)$$

(c)

$$d\phi = -E dx \Rightarrow d\phi = -\frac{2\pi \epsilon k}{\epsilon} (x^2 - w^2) dx$$

$$\Rightarrow \phi = \frac{2\pi \epsilon k}{\epsilon} \left(-\frac{1}{3}x^3 + w^2 x\right) + A$$

$$\text{Let } \phi = 0 \text{ at } x = -w \Rightarrow \frac{2\pi \epsilon k}{\epsilon} \left(\frac{1}{3}w^3 - w^3\right) + A = 0$$

$$\Rightarrow A = \frac{4\pi \epsilon k}{3\epsilon} w^3$$

$$\begin{aligned} \therefore \phi &= \frac{2\pi \epsilon k}{\epsilon} \left(-\frac{1}{3}x^3 + w^2 x\right) + \frac{4\pi \epsilon k}{3\epsilon} w^3 \\ &= \frac{2\pi \epsilon k}{\epsilon} x \left(w^2 - \frac{1}{3}x^2\right) + \frac{4\pi \epsilon k}{3\epsilon} w^3 \end{aligned}$$

(d)

At  $x = +w$ ,  $\varphi = \Delta\varphi$ .

$$\begin{aligned}\therefore \frac{2\epsilon\pi k}{\epsilon} w(w^2 - \frac{1}{3}w^2) + \frac{4\epsilon\pi k}{3\epsilon} w^3 = \Delta\varphi &\Rightarrow \frac{4\epsilon\pi k}{3\epsilon} w^3 + \frac{4\epsilon\pi k}{3\epsilon} w^3 = \Delta\varphi \\ &\Rightarrow \frac{8\epsilon\pi k}{3\epsilon} w^3 = \Delta\varphi \\ &\Rightarrow w = \left[ \frac{3\epsilon\Delta\varphi}{8\epsilon\pi k} \right]^{\frac{1}{3}}\end{aligned}$$