

# PHY 504 Problem Set #2 Solutions

1. (a) Let  $\vec{R} = \frac{m_1 \vec{r}_1 + m_2 \vec{r}_2}{m_1 + m_2} =$  center of mass

$\vec{r}_{12} = \vec{r}_2 - \vec{r}_1 =$  relative position

$\vec{r}'_1 = \vec{r}_1 - \vec{R} =$  center-of-mass coordinate of #1

$\vec{r}'_2 = \vec{r}_2 - \vec{R}$

Then  $\vec{r}'_1 = \vec{r}_1 - \frac{m_1 \vec{r}_1 + m_2 \vec{r}_2}{m_1 + m_2} = \frac{m_2}{m_1 + m_2} \vec{r}_{12}$

$\vec{r}'_2 = -\frac{m_1}{m_1 + m_2} \vec{r}_{12}$

Total kinetic energy  $T = \frac{1}{2} m_1 \dot{\vec{r}}_1^2 + \frac{1}{2} m_2 \dot{\vec{r}}_2^2$

$\frac{1}{2} m_1 \dot{\vec{r}}_1^2 = \frac{m_1}{2} \left( \dot{\vec{R}} + \frac{m_2}{m_1 + m_2} \dot{\vec{r}}_{12} \right)^2 = \frac{1}{2} m_1 \dot{\vec{R}}^2 + \frac{m_1 m_2}{2(m_1 + m_2)} \dot{\vec{r}}_{12}^2 + \frac{m_1 m_2}{m_1 + m_2} \dot{\vec{R}} \cdot \dot{\vec{r}}_{12}$

$\frac{1}{2} m_2 \dot{\vec{r}}_2^2 = \frac{m_2}{2} \left( \dot{\vec{R}} - \frac{m_1}{m_1 + m_2} \dot{\vec{r}}_{12} \right)^2 = \frac{1}{2} m_2 \dot{\vec{R}}^2 + \frac{m_1^2 m_2}{2(m_1 + m_2)} \dot{\vec{r}}_{12}^2 - \frac{m_1 m_2}{m_1 + m_2} \dot{\vec{R}} \cdot \dot{\vec{r}}_{12}$

$\Rightarrow T = \frac{1}{2} (m_1 + m_2) \dot{\vec{R}}^2 + \frac{1}{2} \frac{m_1 m_2}{m_1 + m_2} \dot{\vec{r}}_{12}^2 - \frac{m_1 m_2}{m_1 + m_2} \dot{\vec{R}} \cdot \dot{\vec{r}}_{12}$   
 $= \frac{1}{2} M \dot{\vec{R}}^2 + \frac{1}{2} \mu \dot{\vec{r}}_{12}^2$  where  $\mu = \frac{m_1 m_2}{m_1 + m_2}$ .

Angular momentum

$\vec{L} = \vec{r}_1 \times m_1 \dot{\vec{r}}_1 + \vec{r}_2 \times m_2 \dot{\vec{r}}_2$

$= \left( \vec{R} + \frac{m_2}{m_1 + m_2} \vec{r}_{12} \right) \times \left( m_1 \dot{\vec{R}} + \frac{m_1 m_2}{m_1 + m_2} \dot{\vec{r}}_{12} \right) + (1 \leftrightarrow 2)$

$= \vec{R} \times (m_1 \dot{\vec{R}}) + \frac{m_1 m_2}{(m_1 + m_2)^2} \vec{r}_{12} \times \dot{\vec{r}}_{12} + \frac{m_1 m_2}{m_1 + m_2} \vec{R} \times \dot{\vec{r}}_{12} + \frac{m_1 m_2}{m_1 + m_2} \vec{r}_{12} \times \dot{\vec{R}}$   
 $+ (1 \leftrightarrow 2)$

$= \vec{R} \times M \dot{\vec{R}} + \vec{r}_{12} \times \mu \dot{\vec{r}}_{12}$

(b) In the CM frame,  $\vec{R} = \dot{\vec{R}} = 0$  so

$$T_{cm} = \frac{1}{2} \mu \dot{\vec{r}}_{12}^2 \quad \text{and} \quad \vec{L}_{cm} = \vec{r}_{12} \times \mu \dot{\vec{r}}_{12}$$

If  $\vec{L}_{cm}$  is in the  $\hat{z}$ -direction, then it is perpendicular to  $\vec{r}_{12}$  and  $\dot{\vec{r}}_{12}$ , which therefore both lie in the  $xy$ -plane. Since  $\vec{L}$  is conserved, this is true for all times, so the motion in the CM frame, as given by  $\vec{r}_{12}(t)$ , remains in the  $(xy)$ -plane.

In polar coordinates,  $x_{12} = r \cos \theta$  and  $y_{12} = r \sin \theta$ , we have

$$T_{cm} = \frac{1}{2} \mu (\dot{r}^2 + r^2 \dot{\theta}^2).$$

and

$$\begin{aligned} L_{cm}^z &= \mu (x_{12} \dot{y}_{12} - y_{12} \dot{x}_{12}) \\ &= \mu (r \cos \theta \cdot (\dot{r} \sin \theta + r \dot{\theta} \cos \theta) - r \sin \theta (\dot{r} \cos \theta - r \dot{\theta} \sin \theta)) \\ &= \mu r^2 \dot{\theta} \end{aligned}$$

(d) For a rigid body,  $r = \text{const.}$  so  $\dot{r} = 0$ .

$$\text{Then } T_{cm} = \frac{1}{2} \mu r^2 \dot{\theta}^2 = \frac{L^2}{2\mu r^2} \quad \text{and} \quad E = \frac{L^2}{2\mu r^2} + V(r).$$

The motion will consist of the two particles rotating about their CM at constant  $\dot{\theta}$ , with the CM moving at constant velocity.

2. (a) Newton's 2<sup>nd</sup> Law:  $m\ddot{\vec{r}} = -\vec{\nabla}V$ .

Convert this to polar coordinates:

$$\begin{cases} \dot{x} = \dot{r} \cos \theta - r \dot{\theta} \sin \theta \\ \dot{y} = \dot{r} \sin \theta + r \dot{\theta} \cos \theta \end{cases} \Rightarrow \begin{cases} \ddot{x} = \ddot{r} \cos \theta - 2\dot{r}\dot{\theta} \sin \theta - r\dot{\theta}^2 \cos \theta - r\ddot{\theta} \sin \theta \\ \ddot{y} = \ddot{r} \sin \theta + 2\dot{r}\dot{\theta} \cos \theta - r\dot{\theta}^2 \sin \theta + r\ddot{\theta} \cos \theta \end{cases}$$

Also

$$\frac{\partial V}{\partial x} = \frac{\partial V}{\partial r} \frac{\partial r}{\partial x} + \frac{\partial V}{\partial \theta} \frac{\partial \theta}{\partial x} = \frac{\partial V}{\partial r} \frac{\partial \sqrt{x^2+y^2}}{\partial x} + \frac{\partial V}{\partial \theta} \frac{\partial \tan^{-1}\left(\frac{y}{x}\right)}{\partial x}$$

$$= \frac{\partial V}{\partial r} \frac{x}{r} + \frac{\partial V}{\partial \theta} \left( \frac{-y}{x^2+y^2} \right) = \frac{\partial V}{\partial r} \cos \theta - \frac{\partial V}{\partial \theta} \frac{\sin \theta}{r}$$

$$\frac{\partial V}{\partial y} = \frac{\partial V}{\partial r} \sin \theta + \frac{\partial V}{\partial \theta} \frac{\cos \theta}{r}$$

Plug all this back into Newton:

$$(1) \quad m\ddot{x} = m(\ddot{r} \cos \theta - 2\dot{r}\dot{\theta} \sin \theta - r\dot{\theta}^2 \cos \theta - r\ddot{\theta} \sin \theta) = -\frac{\partial V}{\partial r} \cos \theta + \frac{\partial V}{\partial \theta} \frac{\sin \theta}{r}$$

$$(2) \quad m\ddot{y} = m(\ddot{r} \sin \theta + 2\dot{r}\dot{\theta} \cos \theta - r\dot{\theta}^2 \sin \theta + r\ddot{\theta} \cos \theta) = -\frac{\partial V}{\partial r} \sin \theta - \frac{\partial V}{\partial \theta} \frac{\cos \theta}{r}$$

Multiply (1) by  $\cos \theta$ , (2) by  $\sin \theta$ , and add:

$$\boxed{m(\ddot{r} - r\dot{\theta}^2) = -\frac{\partial V}{\partial r}}$$

Multiply (1) by  $\sin \theta$ , (2) by  $-\cos \theta$ , and add:

$$-2\dot{r}\dot{\theta} - r\ddot{\theta} = \frac{1}{r} \frac{\partial V}{\partial \theta} = 0 \quad (\text{since } V(r) \text{ is central})$$

or  $\boxed{2\dot{r}\dot{\theta} + r\ddot{\theta} = 0}$ . This is equivalent

to conservation of  $L$ :  $\frac{d}{dt}(mr^2\dot{\theta}) = 0$ .

(b) Using results from problem 1,

$$E = T + V = \frac{1}{2} m (\dot{r}^2 + r^2 \dot{\theta}^2)$$

$$L = m r^2 \dot{\theta}$$

The second equation implies  $\dot{\theta} = \frac{L}{m r^2}$ , so

$$E = \frac{1}{2} m \left( \dot{r}^2 + r^2 \left( \frac{L}{m r^2} \right)^2 \right) + V = \frac{1}{2} m \dot{r}^2 + \frac{L^2}{2 m r^2} + V(r).$$

(c) Just differentiate equations in (b) with respect to  $t$ .

$$(d) \quad E = \frac{1}{2} m \left( \frac{dr}{dt} \right)^2 + \frac{L^2}{2 m r^2} + V(r)$$

$$\Rightarrow \quad \frac{dr}{dt} = \sqrt{\frac{2}{m} [E - V(r)] - \frac{L^2}{m^2 r^2}}$$

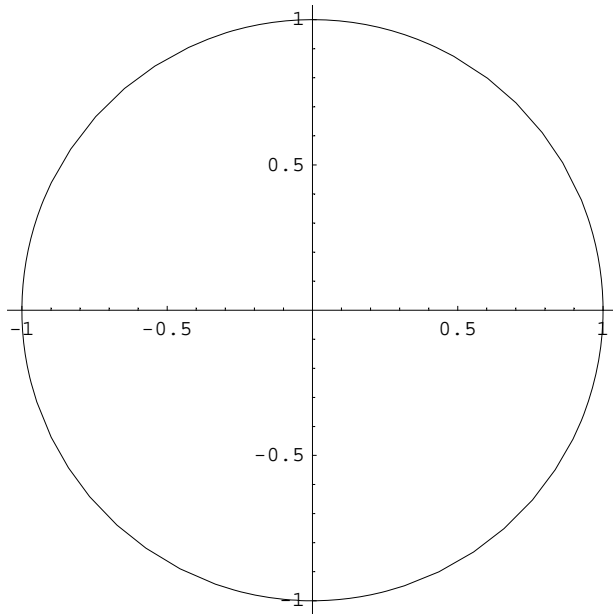
$$\Rightarrow \quad dt = \frac{dr}{\sqrt{\frac{2}{m} (E - V(r)) - \frac{L^2}{m^2 r^2}}}$$

$$\Rightarrow \quad t(r) = \int_{r_0}^r \frac{dr'}{\sqrt{\frac{2}{m} (E - V) - \frac{L^2}{m^2 r'^2}}} + t_0$$

```
In[11]:= Clear[x, y, theta]
          x[theta_] = Cos[theta]
          y[theta_] = Sin[theta]
          ParametricPlot[{x[theta], y[theta]}, {theta, 0, 2*Pi}, AspectRatio -> Automatic]
```

```
Out[12]= Cos[theta]
```

```
Out[13]= Sin[theta]
```



```
Out[14]= - Graphics -
```

3. (b) In 2(d) we found

$$t(r) = \int_{r_0}^r \frac{dr}{\sqrt{\frac{2}{m}(E - V(r)) - \frac{L^2}{m^2 r^2}}} + t_0 = \int_{r_0}^r \frac{m r dr}{\sqrt{2m r^2 (E - \frac{1}{2}(r-a)^2) - L^2}} + t_0$$

This will be integrable in terms of elementary functions only if  $a=0$ , which we assume.

Then substituting  $u=r^2$ ,

$$\begin{aligned} t(r) - t_0 &= \int_{r_0^2}^{r^2} \frac{m/2 du}{\sqrt{2mu(E - \frac{u}{2}) - L^2}} = \frac{\sqrt{m}}{2} \int_{r_0^2 - E}^{r^2 - E} \frac{dv}{\sqrt{(E^2 - L^2/m) - v^2}} \\ &= \frac{\sqrt{m}}{2} \sin^{-1} \left( \frac{v}{\sqrt{E^2 - L^2/m}} \right) \Bigg|_{r_0^2 - E}^{r^2 - E} \end{aligned}$$

or

$$t(r) = \frac{\sqrt{m}}{2} \sin^{-1} \frac{r^2 - E}{\sqrt{E^2 - L^2/m}}$$

where in the last line we chose  $t_0$  to make the constant term vanish.

(c) See attached plot.

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In[100]:=
```

```
  m = 1
```

```
  L = 1
```

```
  e = 2
```

```
  t[r_] = Sqrt[m] / 2 * (ArcSin[ (r^2 - e) / Sqrt[e^2 - L^2 / m] ])
```

```
Out[100]=
```

```
  1
```

```
Out[101]=
```

```
  1
```

```
Out[102]=
```

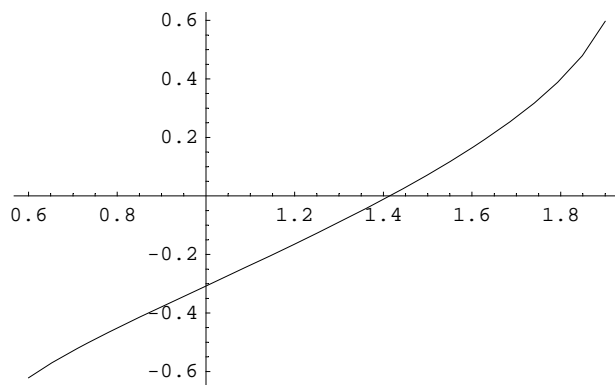
```
  2
```

```
Out[103]=
```

```
   $\frac{1}{2} \text{ArcSin}\left[\frac{-2 + r^2}{\sqrt{3}}\right]$ 
```

```
In[117]:=
```

```
  Plot[t[r], {r, .6, 1.9}]
```



```
Out[117]=
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```
  - Graphics -
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