

PHY 504 Problem Set #3 Solutions

1. Let (q_1, \dots, q_n) be generalized coordinates,

and

$$q_i = q_i(s_1, \dots, s_n, t) \quad i=1, \dots, n$$

a coordinate transformation.

Define

$$L(s_1, \dots, s_n; t) = L(q_1(s_i), q_2(s_i), \dots; t)$$

We wish to show that if

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{q}_j} \right) - \frac{\partial L}{\partial q_j} = 0$$

then
$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{s}_j} \right) - \frac{\partial L}{\partial s_j} = 0.$$

Now
$$\dot{q}_i = \sum_{j=1}^n \frac{\partial q_i}{\partial s_j} \dot{s}_j + \frac{\partial q_i}{\partial t}$$

so
$$\frac{\partial \dot{q}_i}{\partial \dot{s}_j} = \frac{\partial q_i}{\partial s_j}$$

also
$$\frac{\partial L}{\partial s_j} = \sum_i \frac{\partial L}{\partial q_i} \frac{\partial q_i}{\partial s_j} + \frac{\partial L}{\partial \dot{q}_i} \frac{\partial \dot{q}_i}{\partial s_j}$$

$$\frac{\partial L}{\partial s_j} = \sum_i \frac{\partial L}{\partial q_i} \frac{\partial q_i}{\partial s_j} + \frac{\partial L}{\partial \dot{q}_i} \frac{\partial \dot{q}_i}{\partial s_j} = \sum_i \frac{\partial L}{\partial \dot{q}_i} \frac{\partial \dot{q}_i}{\partial s_j}$$

$$\begin{aligned} \therefore \frac{d}{dt} \frac{\partial L}{\partial \dot{s}_j} &= \frac{d}{dt} \sum_i \frac{\partial L}{\partial \dot{q}_i} \frac{\partial \dot{q}_i}{\partial s_j} = \sum_i \left(\frac{d}{dt} \frac{\partial L}{\partial \dot{q}_i} \right) \frac{\partial \dot{q}_i}{\partial s_j} + \frac{\partial L}{\partial \dot{q}_i} \left(\frac{d}{dt} \frac{\partial \dot{q}_i}{\partial s_j} \right) \\ &= \sum_i \frac{\partial L}{\partial q_i} \frac{\partial q_i}{\partial s_j} + \frac{\partial L}{\partial \dot{q}_i} \frac{\partial \dot{q}_i}{\partial s_j} = \frac{\partial L}{\partial s_j} \quad \text{Q.E.D.} \end{aligned}$$

2. (a) Let the length of the pendulum be l . Then

$$L = \frac{1}{2} m v^2 + V = \frac{1}{2} m (l \dot{\theta})^2 + mgl \cos \theta$$

Lagrange equations:

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{\theta}} = \frac{\partial L}{\partial \theta} \Rightarrow$$

$$\frac{d}{dt} (m l^2 \dot{\theta}) = -mgl \sin \theta$$

$$\text{or } \ddot{\theta} = -\frac{g}{l} \sin \theta.$$

(b) For small displacements from $\theta = 0$, $\sin \theta \approx \theta$ and

$$\ddot{\theta} \approx -\frac{g}{l} \theta$$

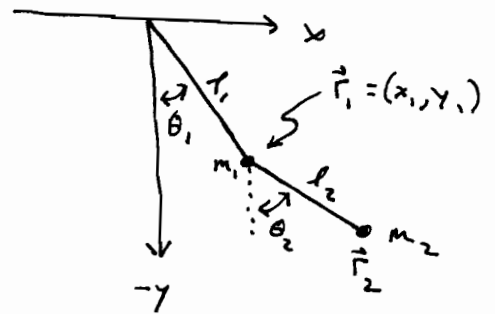
This is an oscillator of frequency $\omega = \sqrt{\frac{g}{l}}$

with general solution

$$\theta(t) = A \cos \sqrt{\frac{g}{l}} (t - t_0)$$

(c) Just copy the code into Mathematica!

3(a) As shown in the figure we let \vec{r}_1 and \vec{r}_2 be the positions of the two masses.



The Lagrangian is

$$L = T - V = \frac{1}{2} m_1 \dot{\vec{r}}_1^2 + \frac{1}{2} m_2 \dot{\vec{r}}_2^2 - m_1 g y_1 - m_2 g y_2$$

and the rigid-rod constraint fixes

$$|\vec{r}_1| = l_1 \quad \text{and} \quad |\vec{r}_2 - \vec{r}_1| = l_2$$

The constraints can be solved by choosing the generalized coordinates (θ_1, θ_2) where

$$\begin{aligned} x_1 &= l_1 \sin \theta_1 & \text{and} & & x_2 - x_1 &= l_2 \sin \theta_2 \\ y_1 &= -l_1 \cos \theta_1 & & & y_2 - y_1 &= l_2 \cos \theta_2 \end{aligned}$$

(b) Then

$$\begin{aligned} L &= \frac{1}{2} m_1 l_1^2 \dot{\theta}_1^2 + \frac{1}{2} m_2 [l_1^2 \dot{\theta}_1^2 + l_2^2 \dot{\theta}_2^2 + 2l_1 l_2 \dot{\theta}_1 \dot{\theta}_2 \cos(\theta_1 - \theta_2)] \\ &\quad + (m_1 + m_2) g l_1 \cos \theta_1 + m_2 g l_2 \cos \theta_2 \end{aligned}$$

The Lagrange equation for θ_1 :

$$\begin{aligned} \frac{d}{dt} \left(\frac{\partial L}{\partial \dot{\theta}_1} \right) &= \frac{d}{dt} \left(m_1 l_1^2 \dot{\theta}_1 + m_2 l_1^2 \dot{\theta}_1 + m_2 l_1 l_2 \dot{\theta}_2 \cos(\theta_1 - \theta_2) \right) \\ &= (m_1 + m_2) l_1^2 \ddot{\theta}_1 + m_2 l_1 l_2 \left[\ddot{\theta}_2 \cos(\theta_1 - \theta_2) - \dot{\theta}_2 (\dot{\theta}_1 - \dot{\theta}_2) \sin(\theta_1 - \theta_2) \right] \\ \frac{\partial L}{\partial \theta_1} &= - (m_1 + m_2) l_1 g \sin \theta_1 - m_2 l_1 l_2 \dot{\theta}_1 \dot{\theta}_2 \sin(\theta_1 - \theta_2) \end{aligned}$$

Putting these together gives

$$(m_1 + m_2) l_1 \ddot{\theta}_1 + m_2 l_2 \ddot{\theta}_2 \cos(\theta_1 - \theta_2) + m_2 l_2 \dot{\theta}_2^2 \sin(\theta_1 - \theta_2) + (m_1 + m_2) g \sin \theta_1 = 0$$

Similarly

$$m_2 l_2 \ddot{\theta}_2 + m_2 l_1 \ddot{\theta}_1 \cos(\theta_1 - \theta_2) - m_2 l_1 \dot{\theta}_1^2 \sin(\theta_1 - \theta_2) + m_2 g \sin \theta_2 = 0$$

(c) For small angles, $\sin \theta \approx \theta$ and $\cos \theta \approx 1 - \frac{\theta^2}{2}$.

Inserting into the above equations of motion and keeping only terms linear in θ :

$$\begin{cases} (m_1 + m_2) l_1 \ddot{\theta}_1 + m_2 l_2 \ddot{\theta}_2 + (m_1 + m_2) g \theta_1 = 0 \\ m_2 l_2 \ddot{\theta}_2 + m_2 l_1 \ddot{\theta}_1 - m_2 g \theta_2 = 0 \end{cases}$$

$$4. (a) \quad L = \frac{1}{2} m \dot{x}^2 - \frac{1}{2} k x^2$$

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{x}} \right) - \frac{\partial L}{\partial x} = m \ddot{x} + kx = 0 \quad \Rightarrow \quad \ddot{x} = -\frac{k}{m} x$$

$$\Rightarrow x(t) = A \cos \sqrt{\frac{k}{m}} t + B \sin \sqrt{\frac{k}{m}} t \quad \Rightarrow \quad \omega = \sqrt{\frac{k}{m}}$$

$$(b) \quad T = \frac{2\pi}{\omega} = 2\pi \sqrt{\frac{m}{k}}$$

Boundary conditions: $x(0) = x(T) = 0 \Rightarrow x(t) = B \sin \sqrt{\frac{k}{m}} t$.

Plug into action:

$$I = \int_0^T L dt = \int_0^T \frac{1}{2} (m \dot{x}^2 - kx^2) dt$$

$$= \frac{1}{2} B^2 \int_0^T [m \omega^2 \cos^2 \omega t - k \sin^2 \omega t] dt$$

$$= \frac{1}{2} k B^2 \int_0^T (\cos^2 \omega t - \sin^2 \omega t) dt$$

$$= \frac{1}{2} k B^2 \int_0^T \cos 2\omega t dt = \frac{1}{2} k B^2 \frac{\sin 2\omega t}{2\omega} \Big|_0^T = 0$$

(c) Suppose $x(t) = \sum_{n=1}^{\infty} a_n \sin(n\omega t)$.

Then $\dot{x} = \sum n\omega a_n \cos(n\omega t)$

$$\text{and } I = \int_0^T L dt = \frac{1}{2} \int_0^T (m\dot{x}^2 - kx^2) dt$$

$$= \frac{1}{2} \int_0^T \left(m\omega^2 \left[\sum n a_n \cos(n\omega t) \right]^2 - k \left[\sum a_n \sin(n\omega t) \right]^2 \right) dt$$

$$= \frac{k}{2} \int_0^T \left[\sum_n n^2 a_n^2 \cos^2(n\omega t) + \sum_{n \neq m} n a_n a_m \cos(n\omega t) \cos(m\omega t) - \sum_n a_n^2 \sin^2(n\omega t) \right] dt$$

$$= \frac{k}{2} \sum_n \left(n^2 a_n^2 \cdot \frac{T}{2} - a_n^2 \cdot \frac{T}{2} \right) dt$$

$$= \frac{kT}{4} \sum_n (n^2 - 1) a_n^2 = \frac{kT}{2\omega} \sum_{n=1}^{\infty} (n^2 - 1) a_n^2$$

(d) Because $a_n^2 > 0$ and $n \geq 1$, $I \geq 0$ for any path, and the path studied in (b) is a minimum.