

PHY 504 Problem Set #8 Solutions

1. In the rotating frame,

$$\ddot{\vec{r}} = -2(\vec{\omega} \times \dot{\vec{r}}) - \vec{\omega} \times (\vec{\omega} \times \vec{r}) + m\vec{g}$$

The Coriolis and centrifugal accelerations are both in the horizontal xy -plane.

The horizontal components of the equation of motion are

$$\begin{cases} \ddot{x} = 2\omega\dot{y} + \omega^2 x \\ \ddot{y} = -2\omega\dot{x} + \omega^2 y \end{cases}$$

exactly as found in problem 1 of PS#1.

These may be solved by various techniques; however we have already found the solution, so let's use it:

$$\begin{cases} x(t) = (x_0 - \omega y_0 t) \cos \omega t + (y_0 + \omega x_0 t) \sin \omega t \\ y(t) = -(x_0 - \omega y_0 t) \sin \omega t + (y_0 + \omega x_0 t) \cos \omega t \end{cases}$$

Let's take $x_0 = r_0$, $y_0 = 0$.

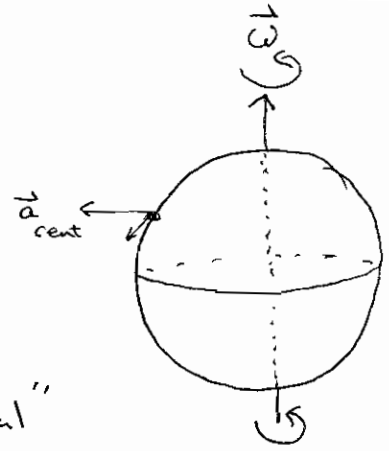
If the ball is dropped from height $h = \frac{1}{2}gt^2$,

then $t_f = \sqrt{\frac{2h}{g}}$ and

$$\begin{aligned} x(t_f) &= r_0 \cos \omega \sqrt{\frac{2h}{g}} + \omega r_0 \sqrt{\frac{2h}{g}} \sin \omega \sqrt{\frac{2h}{g}} \\ y(t_f) &= -r_0 \sin \omega \sqrt{\frac{2h}{g}} + \omega r_0 \sqrt{\frac{2h}{g}} \cos \omega \sqrt{\frac{2h}{g}} \end{aligned}$$

are the final coordinates of the ball.

2. (a) The centrifugal acceleration points radially away from the N-S axis of Earth and at 38° N latitude has a "horizontal" component pointing to the south, of magnitude



$$\begin{aligned} a_{cent,s} &= -\left[\vec{\omega} \times (\vec{\omega} \times \vec{r})\right]_s \\ &= \left[\omega^2 R \cos 38^\circ \hat{r}\right]_s \\ &= \omega^2 R \cos 38^\circ \sin 38^\circ \\ &= 0.0338 \text{ m/s}^2 \cdot 0.485 = 0.0164 \text{ m/s}^2 \\ &\quad [\text{p. 176}] \end{aligned}$$

This leads to no deflection for the bullet fired south.
horizontal

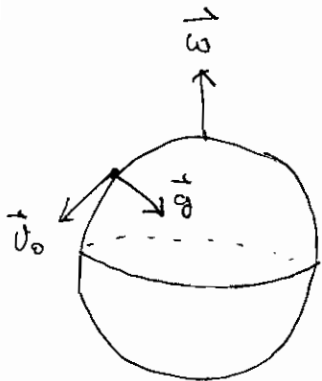
During the 10s the bullet is in the air, the effect of gravity on its velocity is significant, and must be taken into account:

$$\vec{v}(t) = \vec{v}_0 + \vec{g}t \quad \text{where } \vec{g} = (-9.8 \text{ m/s}^2) \hat{r}$$

Then

$$\begin{aligned} \vec{a}_{cor} &= -2(\vec{\omega} \times \vec{v}) = -2\vec{\omega} \times (\vec{v}_0 + \vec{g}t) \\ &= -2(\omega v_0 \sin 38^\circ \hat{E} - \omega g t \cos 38^\circ \hat{E}) \end{aligned}$$

↑ unit vector to East



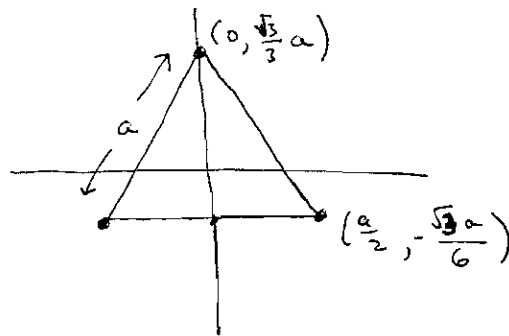
So the Coriolis acceleration leads to a net deflection
to the East of

$$\begin{aligned}\Delta X_{\text{cor}} &= -2 \left(\omega U_0 \sin 38^\circ \frac{t^2}{2} - \omega g \frac{t^3}{3} \cos 38^\circ \right) \\ &= -2 \left(7.3 \times 10^{-5} \text{ s}^{-1} \right) \left[(100 \text{ m/s}) (0.6) \frac{10^2}{2} - 9.8 \text{ m/s}^2 (0.8) \frac{(10 \text{ s})^3}{3} \right] \\ &= -0.056 \text{ m}\end{aligned}$$

The bullet will be deflected to the west by 0.056 m.

(b) Left to the reader.

3. We will calculate \underline{I} about the center of mass



$$I_{xx} = \int_{-a/2}^{a/2} dx \int_{-\frac{\sqrt{3}a}{6}}^{\frac{\sqrt{3}a}{3} - \sqrt{3}|x|} dy \left((x^2 + y^2) - x^2 \right) \quad \text{where } \rho = \frac{m}{\text{area}} = \text{density}$$

$$= 2 \cdot \left(\frac{4m}{\sqrt{3}a^2} \right) \int_0^{a/2} dx \int_{-\frac{\sqrt{3}a}{6}}^{\frac{\sqrt{3}a}{3} - \sqrt{3}x} dy \cdot y^2$$

$$= \frac{8m}{\sqrt{3}a^2} \int_0^{a/2} dx \left[\frac{1}{3} \left(\frac{\sqrt{3}a}{3} - \sqrt{3}x \right)^3 - \frac{1}{3} \left(-\frac{\sqrt{3}a}{6} \right)^3 \right]$$

$$= \frac{8m}{\sqrt{3}a^2} \left[\frac{1}{12} \left(\frac{\sqrt{3}a}{3} - \sqrt{3}x \right)^4 \cdot \left(-\frac{1}{\sqrt{3}} \right) + \frac{1}{3} \left(\frac{\sqrt{3}a}{6} \right)^3 \cdot x \right]_0^{a/2}$$

$$= \frac{8m}{\sqrt{3}a^2} \left[\frac{1}{12} \left(\frac{\sqrt{3}a}{6} \right)^4 \left(-\frac{1}{\sqrt{3}} \right) - \frac{1}{12} \left(\frac{\sqrt{3}a}{3} \right)^4 \left(-\frac{1}{\sqrt{3}} \right) + \frac{1}{3} \left(\frac{\sqrt{3}a}{6} \right)^3 \cdot \frac{a}{2} \right]$$

$$= \frac{8m}{\sqrt{3}a^2} \left(\frac{\sqrt{3}a}{6} \right)^4 \left[\frac{1}{12} \left(-\frac{1}{\sqrt{3}} \right) - \frac{16}{12} \left(-\frac{1}{\sqrt{3}} \right) + \frac{1}{\sqrt{3}} \right]$$

$$= \frac{8m}{3a^2} \cdot \frac{9}{6^4} a^4 \cdot \frac{27}{12} = \frac{1}{24} ma^2$$

4.(a) With the edges of the cube parallel to the x -, y -, and z -axes and a corner at the origin, the diagonal components of I are equal to

$$\begin{aligned} I_{xx} &= \sum_n m (r_{(n)}^2 - x_{(n)}^2) \\ &= m \cdot [3a^2 + 3(\sqrt{2}a)^2 + (\sqrt{3}a)^2 - 4a^2] \\ &= 8ma^2 \end{aligned}$$

Symmetry also implies that the off-diagonal components are equal to

$$\begin{aligned} I_{xy} &= \sum_n m (-x_{(n)}y_{(n)}) \\ &= -2ma^2 \end{aligned}$$

So

$$I = ma^2 \begin{pmatrix} 8 & -2 & -2 \\ -2 & 8 & -2 \\ -2 & -2 & 8 \end{pmatrix}$$

The long diagonal of the cube is

the vector $\vec{d} = (1, 1, 1)$. Check if this is eigenvector:

$$\underline{I} \vec{d} = ma^2 \begin{pmatrix} 8 & -2 & -2 \\ -2 & 8 & -2 \\ -2 & -2 & 8 \end{pmatrix} \begin{pmatrix} 1 \\ 1 \\ 1 \end{pmatrix} = 4ma^2 \begin{pmatrix} 1 \\ 1 \\ 1 \end{pmatrix}$$

It is an eigenvector with eigenvalue $4ma^2$.

(b) Choosing coordinates so that the z -axis is aligned with the long axis of the cube, $\hat{d} = -\hat{z}$, \hat{z} will now be an eigenvector of \underline{I} , so that

$$\underline{I} = \begin{pmatrix} I_{xx} & I_{xy} & 0 \\ I_{yx} & I_{yy} & 0 \\ 0 & 0 & 4ma^2 \end{pmatrix}$$

↑ eigenvalue of I
in \hat{z} direction

\hat{z} is a principal axis with moment of inertia $\underline{I}_3 = 4ma^2$.

Since $\vec{L}_z = \underline{I} \vec{\omega}$ is conserved in the space frame, in the absence of torques,

$$\frac{d}{dt} \vec{L} = \frac{d}{dt} (\underline{I} \vec{\omega}) = \frac{d}{dt} (\underline{I}_3 \omega_3) = \underline{I}_3 \dot{\omega}_3 = 0$$

if the cube rotates about the z -axis, so the cube will spin with constant ω_3 .

If a torque N_3 is applied in the z -direction,

then $\underline{I}_3 \dot{\omega}_3 = N_3$ and the angular acceleration is

$$\alpha = \dot{\omega}_3 = \frac{N_3}{\underline{I}_3} = \frac{N_3}{4ma^2}.$$

Similarly

$$I_{yy} = 2\rho \int_0^{a/2} dx \int_{-\frac{\sqrt{3}}{6}a}^{\frac{\sqrt{3}}{2}a - \sqrt{3}x} dy x^2 = 2\rho \int_0^{a/2} dx x^2 \cdot \left(\frac{\sqrt{3}}{2}a - \sqrt{3}x\right)$$
$$= 2\rho \left(\frac{\sqrt{3}}{2}a \frac{x^3}{3} - \sqrt{3} \frac{x^4}{4} \right) \Big|_0^{a/2} = \frac{1}{24} ma^2.$$

$$I_{xy} = \int dx \int dy (-xy) = 0$$

\triangle
by antisymmetry ~~about~~ of integrand under $x \rightarrow -x$.

$$I_{zz} = \iint dx dy (x^2 + y^2) = I_{xx} + I_{yy} = \frac{1}{12} ma^2$$

$$I_{zx} = I_{zy} = 0 \text{ as for } I_{xy}.$$

Finally

$$I = \begin{pmatrix} \frac{1}{24} ma^2 & 0 & 0 \\ 0 & \frac{1}{24} ma^2 & 0 \\ 0 & 0 & \frac{1}{12} ma^2 \end{pmatrix}$$

(c) The eigenvalues of I , i.e. its principal moments, are obtained by solving

$$\det(I - \lambda \mathbb{I}) = 0$$

which gives (we omit the details) $\lambda = 4ma^2, 10ma^2, 10ma^2$.

We have already seen that $\lambda = 4ma^2$ is the moment for rotation about the vertical axis.

Thus $\lambda = 10ma^2$ must be the ~~eigen~~ moments about the horizontal axes. (Since these moments are degenerate, they are unchanged under a horizontal rotation, so it doesn't matter how we choose our x - and y -axes.)

Let θ denote the angle made by the long diagonal of the cube to the vertical.

Then

$$\begin{aligned} L &= \frac{1}{2} I_1 \dot{\theta}^2 + mg l \cos \theta \\ &\approx \frac{1}{2} I_1 \dot{\theta}^2 - \frac{1}{2} mg l \theta^2 \end{aligned}$$

$l = \frac{\sqrt{3}}{2} a = \text{dist from pivot to center of mass}$

The frequency is

$$\begin{aligned} \omega &= \sqrt{\frac{k}{m}} = \sqrt{\frac{mg l}{I_1}} = \sqrt{\frac{mg \cdot \frac{\sqrt{3}}{2} a}{10ma^2}} \\ &= \sqrt{\frac{\sqrt{3}}{20} \frac{g}{a}} \end{aligned}$$

5. (a) The cube rotates about the edge in contact with the table. The moment of inertia about this axis is

$$I = 4ma^2 + 2m(\sqrt{2}a)^2 = 8ma^2.$$

The gravitational potential energy is

$$V = 8mg \left(\frac{a}{\sqrt{2}} \right) \cos \theta$$

where $\frac{a}{\sqrt{2}} \cos \theta$ is the height of the center of mass above the table when the cube is tilting at an angle θ away from its initial position.

The conserved total energy of the cube is

$$E = \frac{1}{2} I \dot{\theta}^2 + \frac{8}{\sqrt{2}} m g a \cos \theta.$$

Since the initial energy is $\frac{8}{\sqrt{2}} m g a$, we have

$$\frac{1}{2} I \dot{\theta}^2 + \frac{8}{\sqrt{2}} m g a \cos \theta = \frac{8}{\sqrt{2}} m g a$$

$$\begin{aligned} \text{or } \dot{\theta} &= \sqrt{\frac{2}{I} (m g a) \frac{8}{\sqrt{2}} \cdot (1 - \cos \theta)} \\ &= \sqrt{\frac{g}{a} (\sqrt{2} - 1)} \end{aligned}$$

when $\theta = 45^\circ$.

(b) Similar.