

# PHYSICS OF ACTIVE GALACTIC NUCLEI

## 活動銀河核の物理学

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# LECTURE 3

## AGN: INFLOWS VERSUS OUTFLOWS

Accretion and AGN power

Winds from accretion disks and driving mechanisms

Broad-line regions in AGN and outflows

AGN tori and outflows

Collimated winds: jets

Where is accretion and where is an outflow

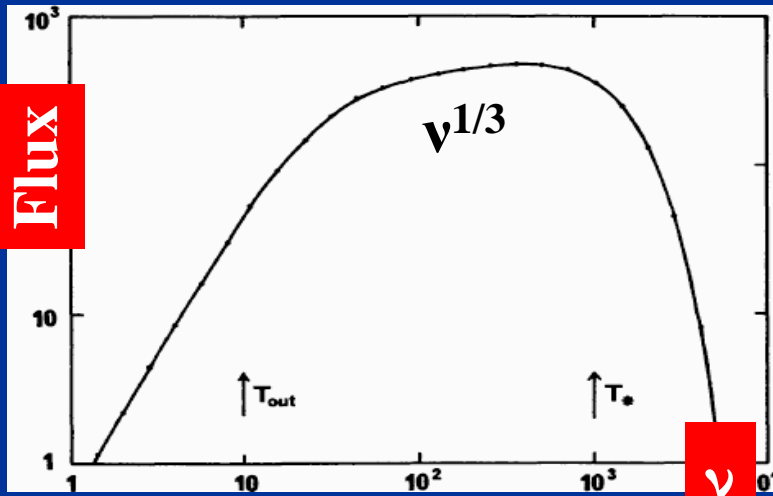
# OBSERVATIONAL EVIDENCE FOR INFLOWS AND OUTFLOWS IN AGN

## ❖ Review of AGN energetics

$$L = \eta \dot{M} c^2 \sim 6 \times 10^{45} \left( \frac{\eta}{0.1} \right) \left( \frac{\dot{M}}{1 M_{\odot}/\text{yr}} \right) \text{ erg/s} \quad \text{accretion luminosity}$$

$$L = \frac{4\pi G c m_p}{\sigma_T} M_{\bullet} \sim 1.3 \times 10^{46} \left( \frac{M_{\bullet}}{10^8 M_{\odot}} \right) \text{ erg/s} \quad \text{the Eddington luminosity}$$

$$T_{\text{bb}} \sim \left( \frac{L}{\pi R^2 \sigma} \right)^{1/4} \sim 10^5 \left( \frac{L}{10^{46} \text{ erg/s}} \right)^{1/4} \left( \frac{R}{1 \text{ light-day}} \right)^{-1/2} \text{ K} \quad \begin{array}{l} \text{effective temperature of} \\ \text{blackbody accretion disk} \\ \text{emission} \rightarrow \text{peaks at } \sim 10^{16} \text{ Hz} \end{array}$$



accretion disk  
continuum

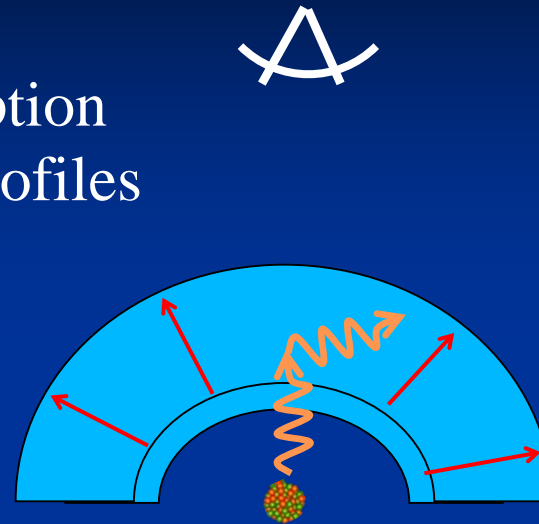
# OBSERVATIONS: INFLOWS AND OUTFLOWS IN AGN

## ❖ Analogy with hot stars

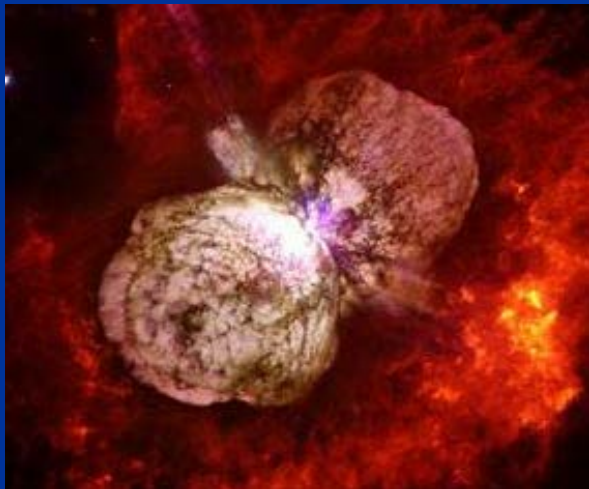
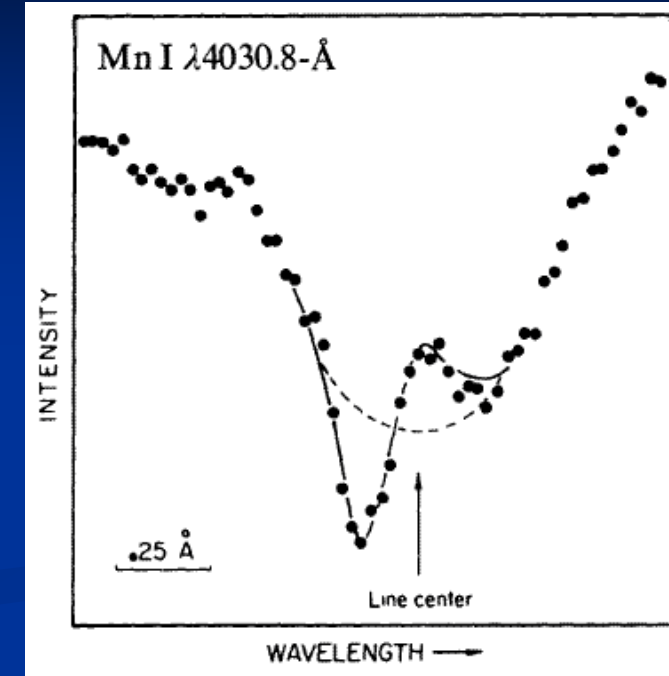
Optical/UV spectroscopy: blueshifted absorption or P Cygni line profiles

$$\mathbf{a}_{\text{rad}} = \int d\nu \frac{\kappa_{\nu} \mathbf{F}_{\nu}}{c}$$

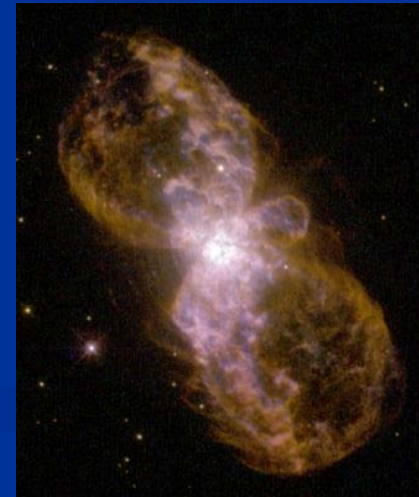
$\mathbf{F}_{\nu}$  -- stellar radiation flux  
 $\mathbf{a}_{\text{rad}}$  -- radiative acceleration  
 $\kappa_{\nu}$  -- opacity coeff.



wind from an O star



Eta Carina



# OBSERVATIONS: INFLOWS AND OUTFLOWS IN AGN

## ❖ Analogy with solar (Parker) wind

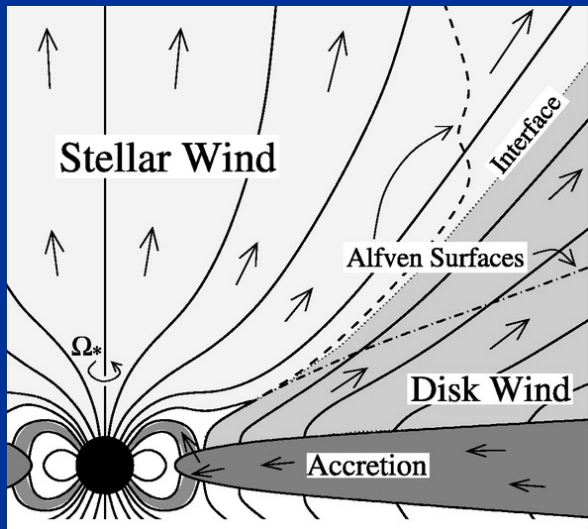
What triggers the solar wind:  
waves and turbulence,  
or magnetic reconnection?

$$\dot{M} \propto F_{\text{heat}} \propto F_X$$

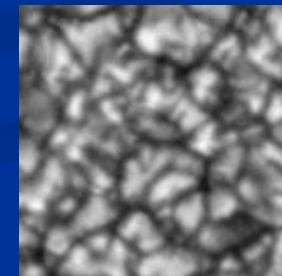
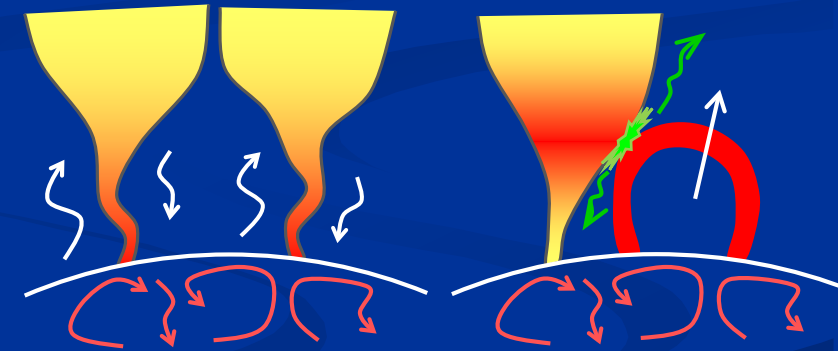
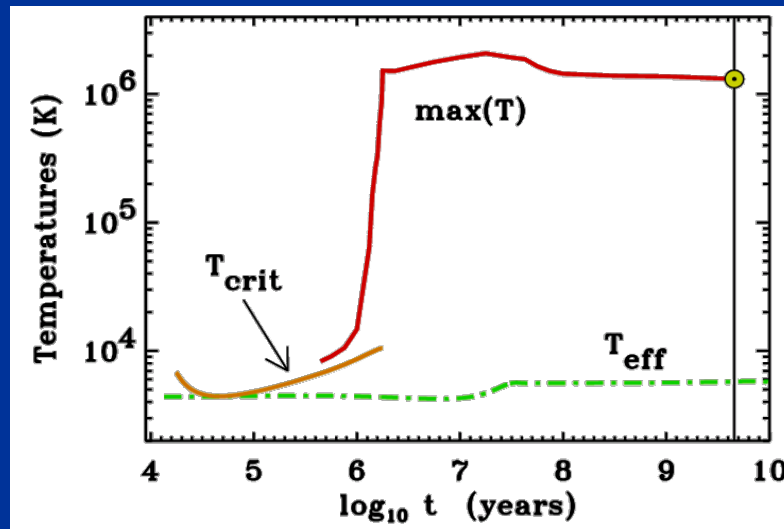


solar wind

Some stars have disks (T Tauri) and winds



(Matt & Pudritz 2005)



MHD turbulence  
→ granulation



# OBSERVATIONS: INFLOWS AND OUTFLOWS IN AGN

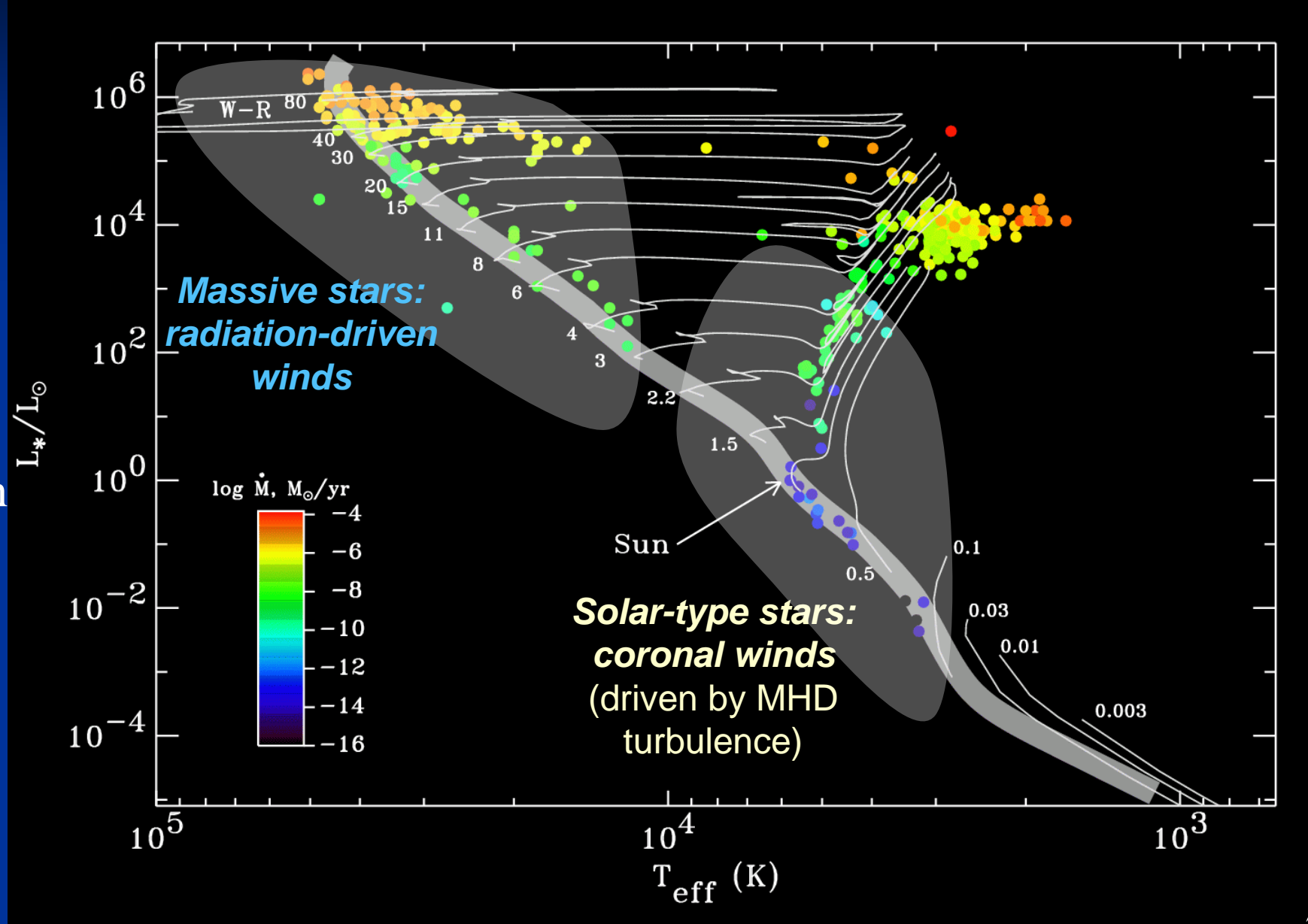
❖ Analogy with hot stars

$T$ -range for radiation-driven and coronal stellar winds

$T$ -range for geometrically-thin accretion disks in AGN is similar to stellar range!



AGN winds can be radiation-driven!

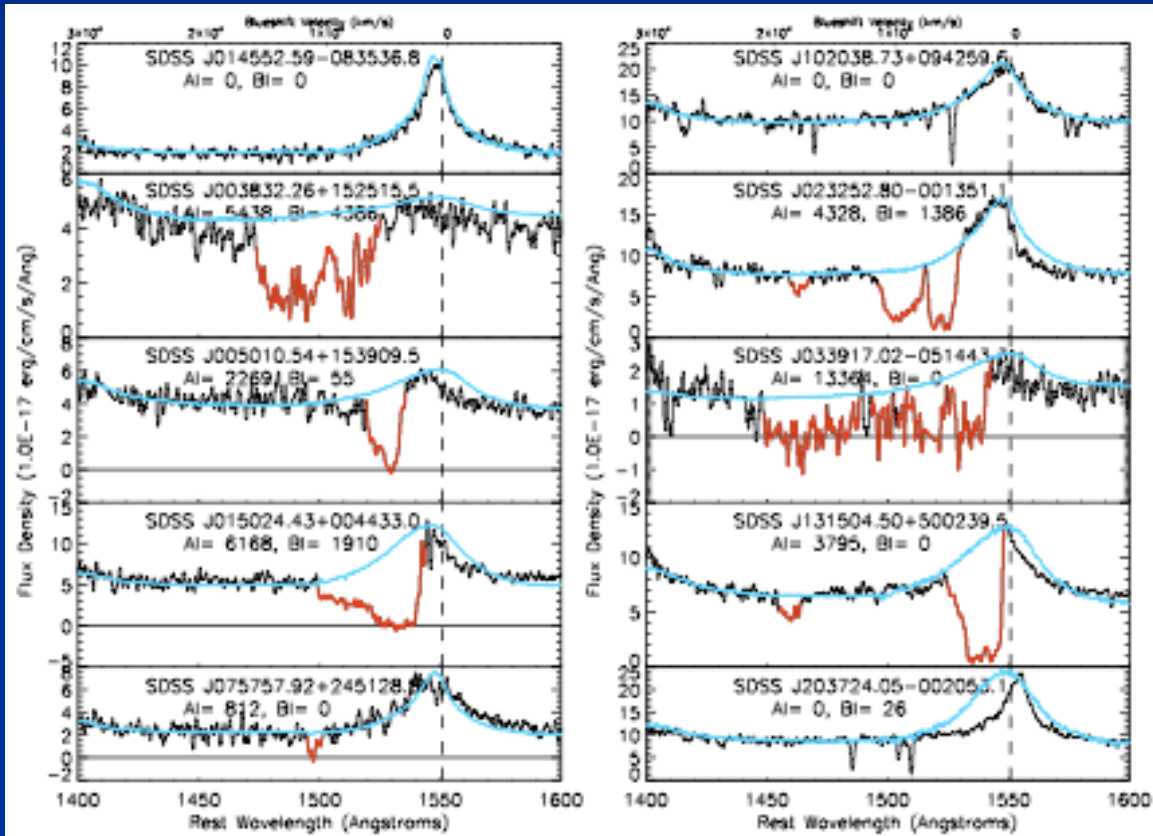


# OBSERVATIONS: INFLOWS AND OUTFLOWS IN AGN

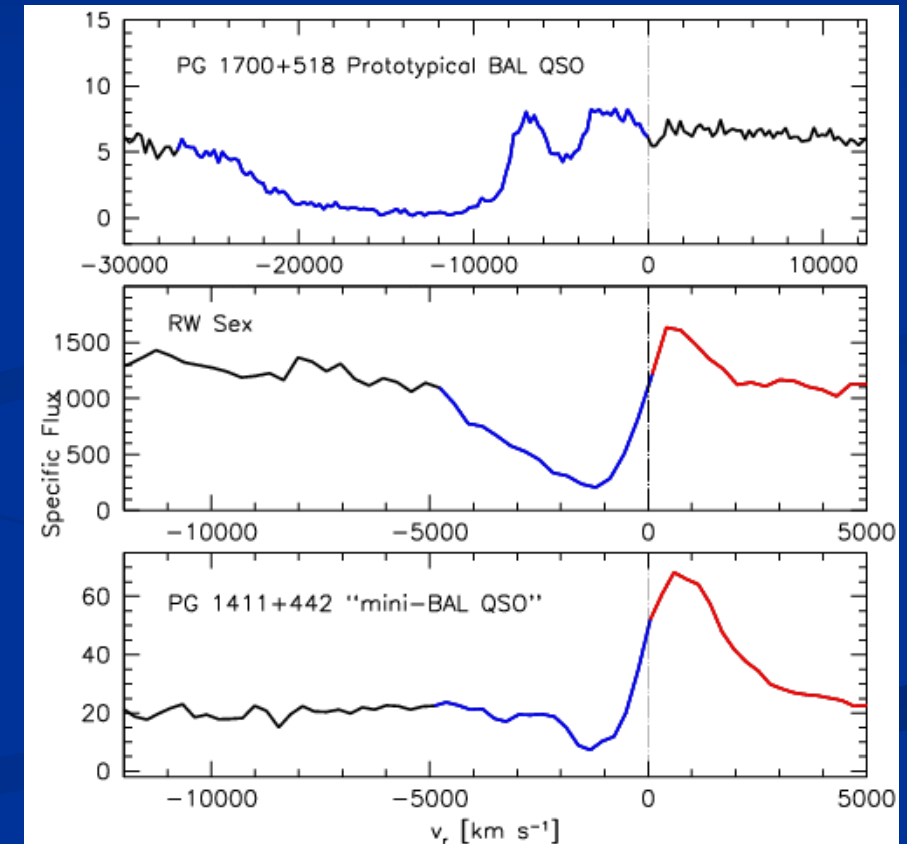
❖ Evidence for outflows: BAL QSOs

7 BAL QSOs (red throughs), 2 non-BAL QSOs

Similarity of the CIV line profile of the nova-like variable RW Sex with those of BAL QSOs



Trump et al. 2006



All BALQSOs absorption lines are blueshifted → outflows !

# OBSERVATIONS: INFLOWS AND OUTFLOWS IN AGN

❖ Evidence for outflows: high-ionization UV emission lines (HILs) in QSOs

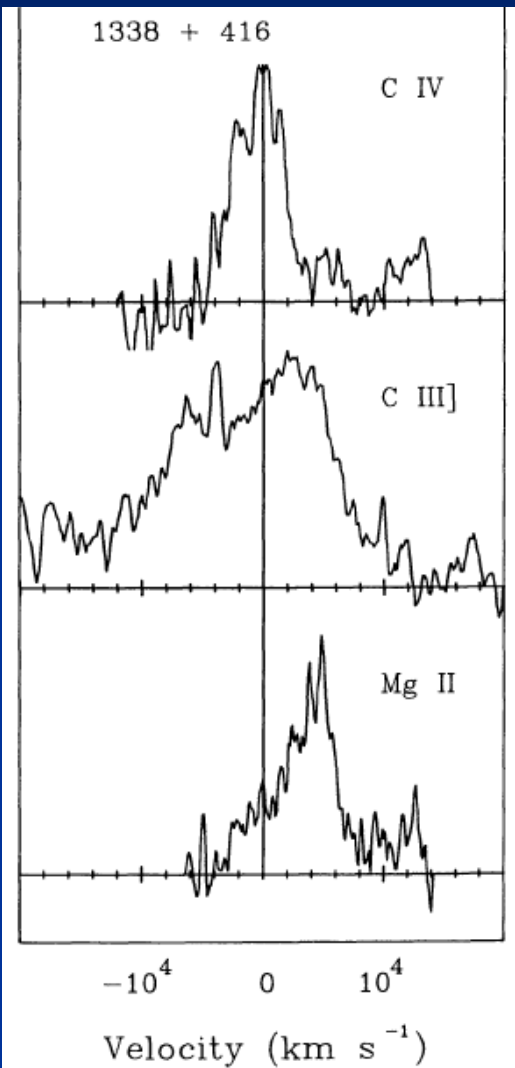
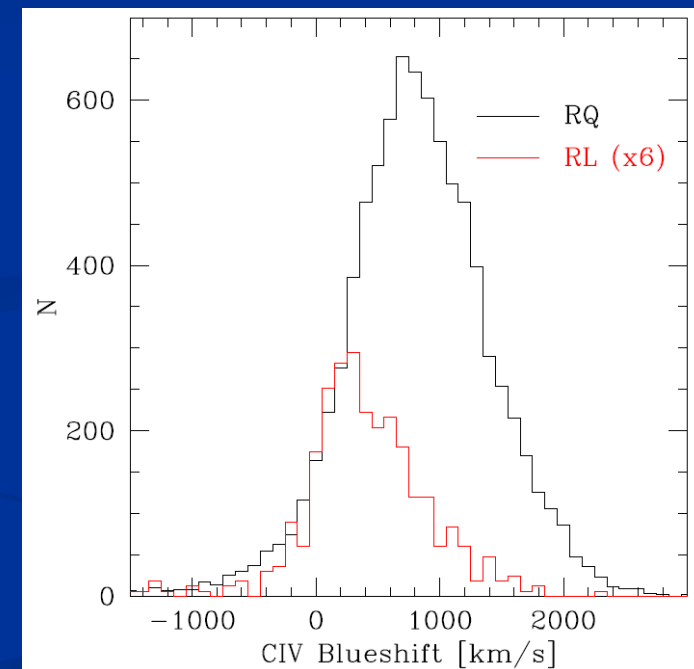


TABLE 1  
REDSHIFT DIFFERENCES

Object	Velocity of Mg II relative to C IV ( $\text{km s}^{-1}$ )	Velocity of C III] relative to C IV ( $\text{km s}^{-1}$ )
BQS		
0117+213 .....	1200	900
1008+133 .....	2500	1300
1241+176 .....	1200	900
1338+416 .....	4400	2500
1352+011 .....	1500	500
1522+101 .....	1600	200
1630+377 .....	1400	100
1634+706 .....	200	200
Mean .....	$1750 \pm 439$	$825 \pm 282$

Corbin et al. (1990)

blueshifted CIV emission line in luminous radio-quiet and radio loud QSOs



Richards et al. (2011)

**HILs in QSOs are blueshifted**  
**By ~few x 100 – 1000 km/s → these broad line produced in winds !**



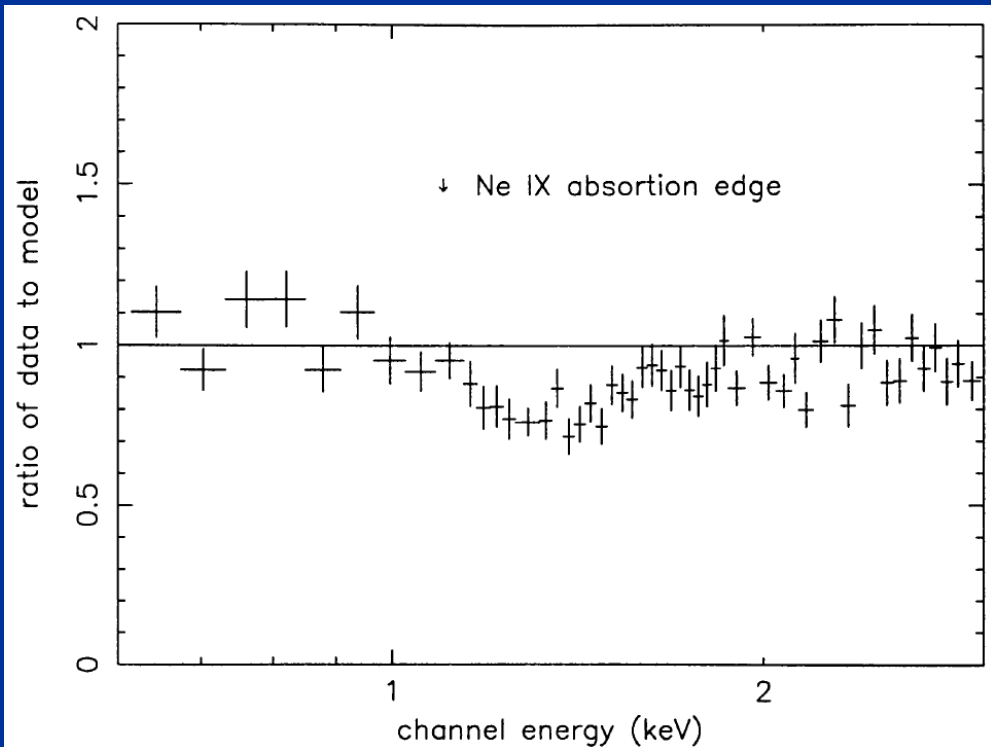
# OBSERVATIONS: INFLOWS AND OUTFLOWS IN AGN

❖ Evidence for outflows: warm absorbers in Seyfert galaxies

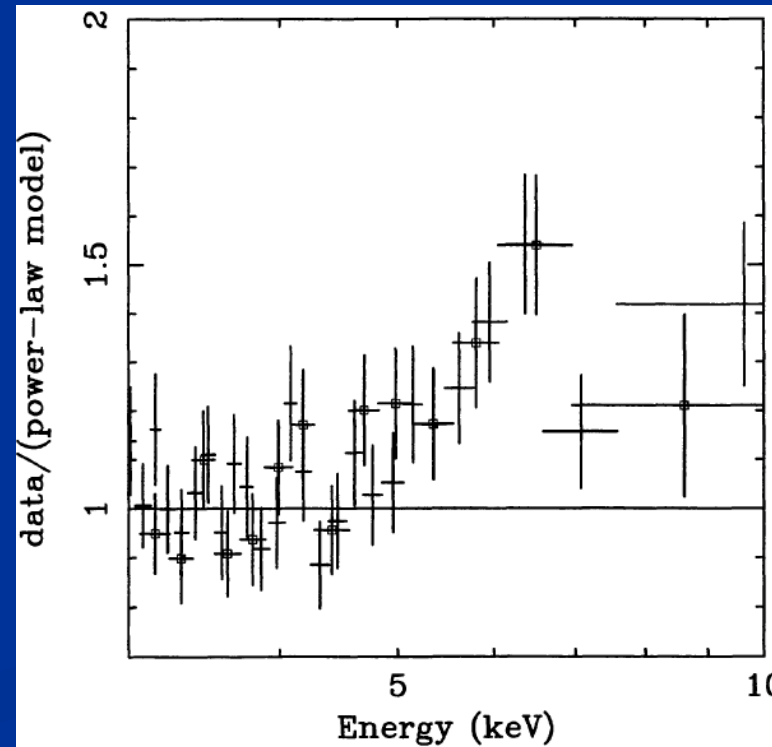
Covering factors can be as high as 0.5 ← about 50% of Seyferts show warm absorber

warm absorber fit to the ASCA satellite spectrum of NGC 3783 → Ne IX absorption line

Fe K line in Fairal 9 from ASCA



Mushotzky (1997)



Reynolds (1997)

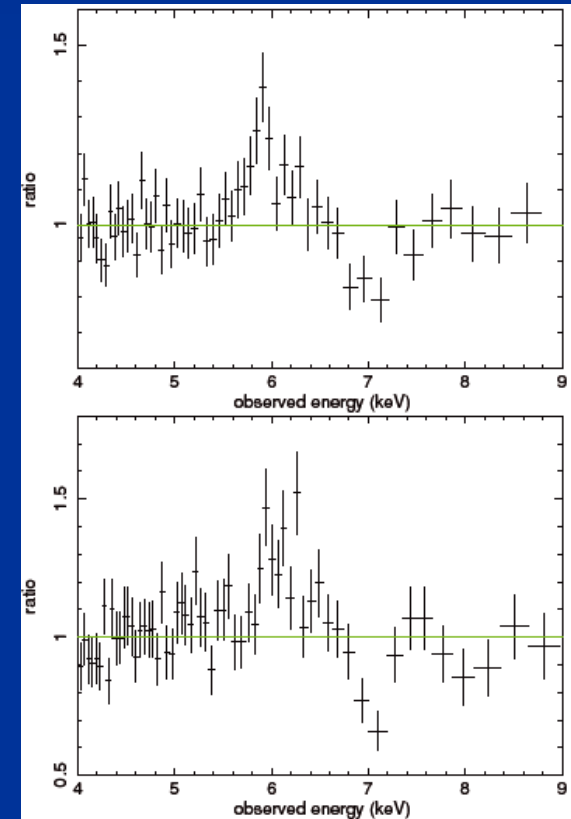
# OBSERVATIONS: INFLOWS AND OUTFLOWS IN AGN

## ❖ Evidence for outflows: ultra-fast X-ray outflows (UFOs)

Very fast outflows of highly ionized material by XMM/Newton in absorption lines of highly ionized Fe, S, Mg → **warm absorbers?**

Fe K $\alpha$  line normalized by continuum emission:

absorption line well defined, but emission line much less defined



$$L_{\text{kin}} \sim 0.1 L_{\text{bol}}, \quad v \sim 0.13c !$$

The physical implications are:

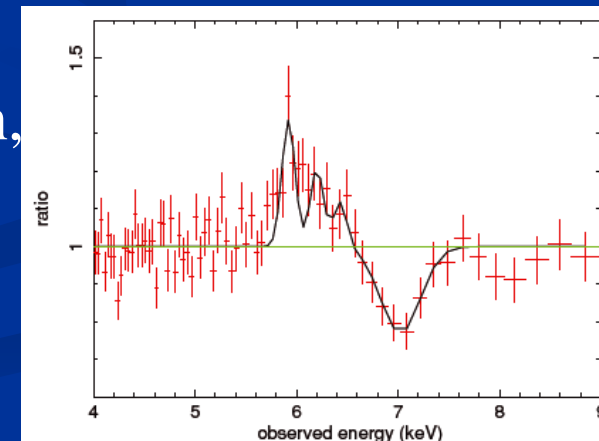
persistent (>6-7 yrs), massive wide-angle wind, covering 0.3-0.6,  
 $\dot{M}$  (wind)  $\sim \dot{M}$  (accretion), column density  $N_{\text{H}} \sim 8 \times 10^{23} \text{ cm}^{-2}$ ,  
*super-Eddington* → Compton-thick wind?

**BUT:**

Often only one line is detected: unsure identification,  
ionization/column density

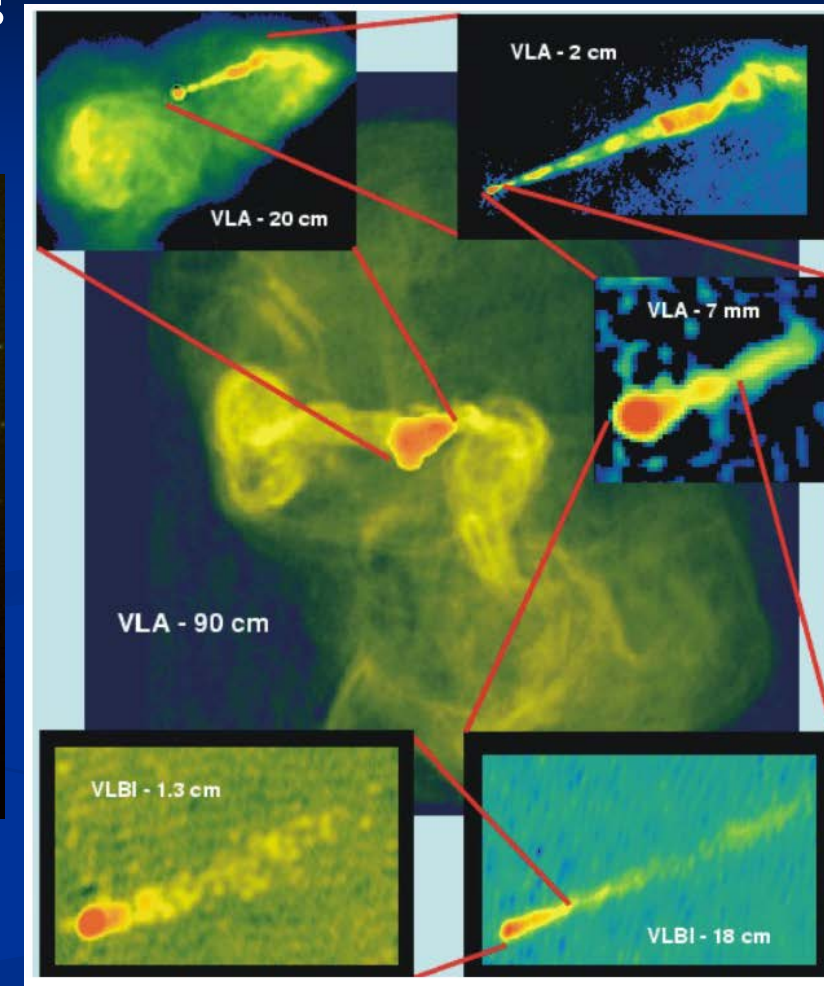
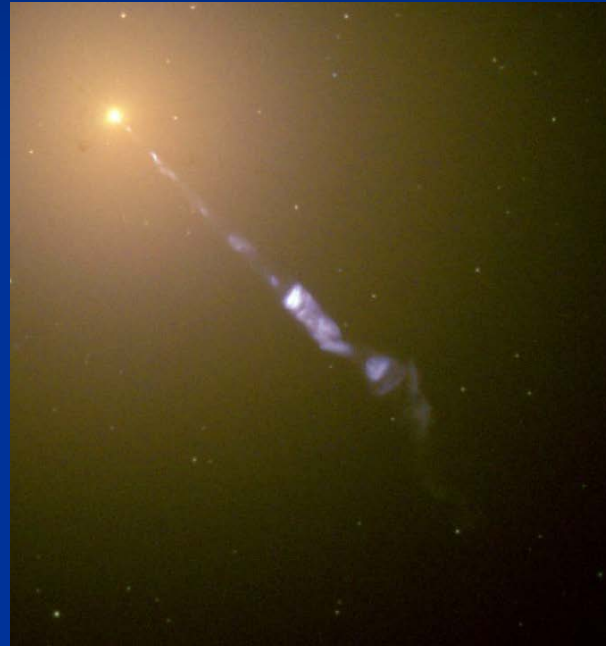
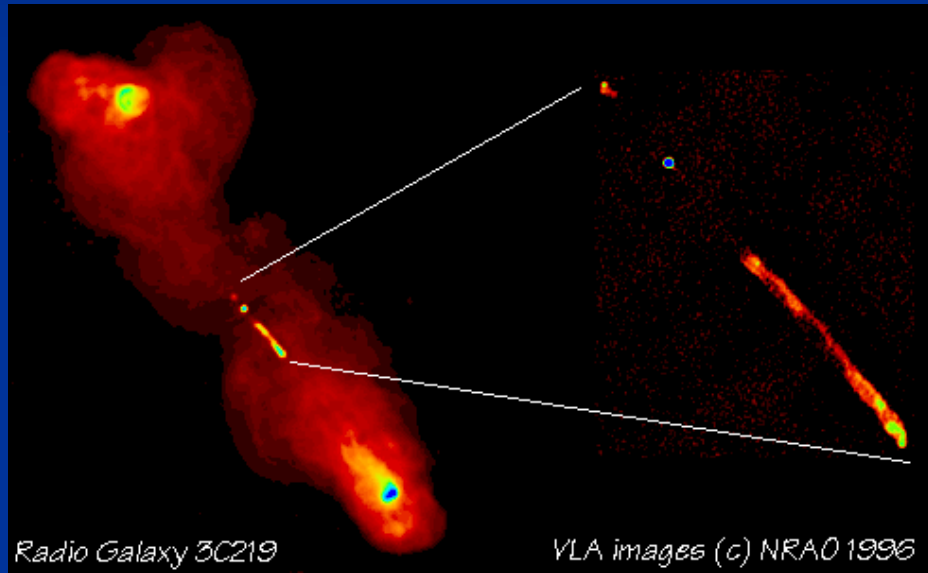
Region strongly influenced by:

background subtraction, continuum modelling,  
lower effective area/resolution



# OBSERVATIONS: INFLOWS AND OUTFLOWS IN AGN

❖ Evidence for outflows: jets in radio and other (AGN) galaxies



Jet in M87 → optical emission is  
synchrotron mechanism  
(electrons accelerated in B-field)

# OBSERVATIONS: INFLOWS AND OUTFLOWS IN AGN

## ❖ Evidence for outflows: summary

Uncollimated or partially collimated winds:

UV resonance lines in QSOs and Seyferts (line-driven winds)

BAL QSOs

UV absorbers: warm absorbers, UFOs (X-ray Ultra-Fast Outflows)

Super-Eddington winds?

Collimated winds (jets)

Powerful and not so powerful radio galaxies

(RLQ, FR II, FR I), Seyferts

LLAGN (XRB hard state compact jets)



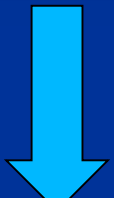
# THEORY: INFLOWS AND OUTFLOWS IN AGN

❖ Summary: radiation-driven winds

To drive a wind by radiation  $\rightarrow$  need opacity  $\rightarrow F_{\text{rad}} > F_{\text{grav}}$

resonance lines

e.g., C IV



wind  $T < 10,000$  K



dust



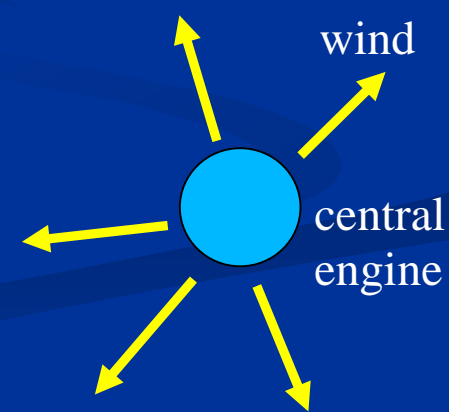
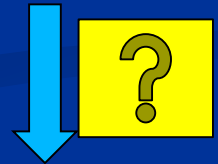
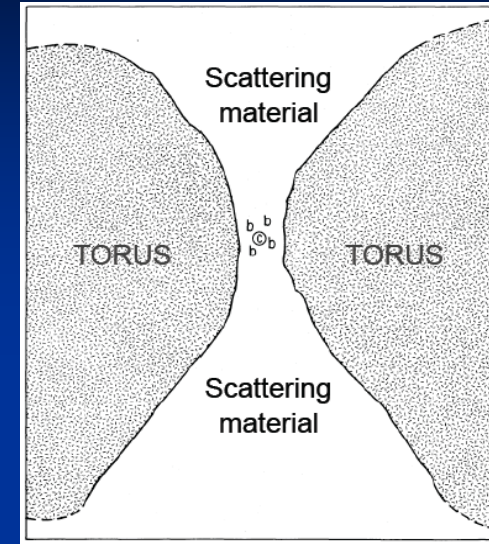
wind  $T < 2,000$  K



electron scattering



wind fully ionized  
 $T > 10,000$  K



➤ But linewidth  $\sim 10^4$  km/s  
 $\rightarrow T_{\text{gas}} \sim 10^{10}$  K!  
**No lines! No dust!**  
 ➤ What is the geometry of the wind?

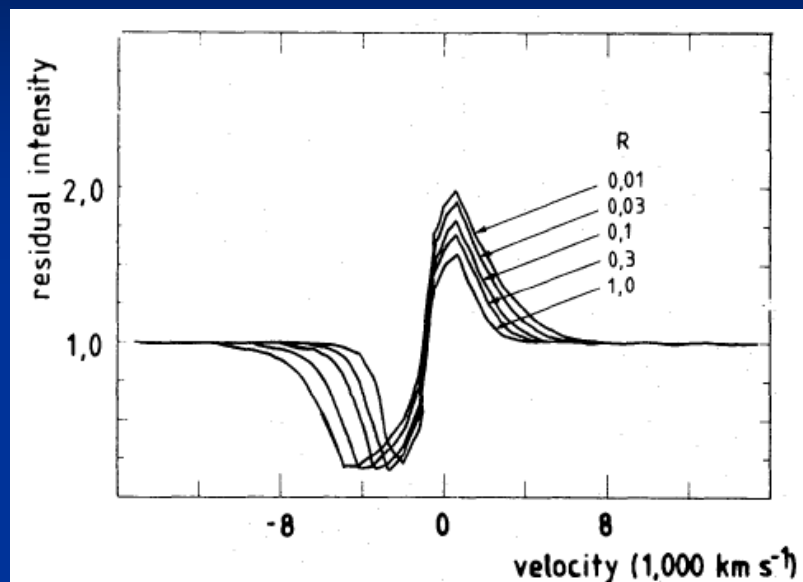
**BLR must be clumpy  
 but  
 what about wind?**

**super-Eddington  
 wind?**



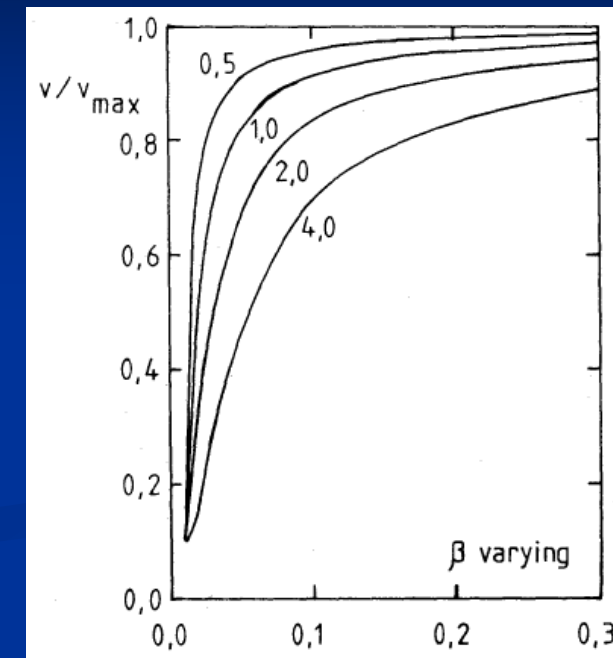
# THEORY: INFLOWS AND OUTFLOWS IN AGN

- ❖ Radiation-driven disk winds: driving spherical wind by resonance lines



line profiles similar  
to stellar line profiles

$$v(r)/v_{\max} = w_c + (1 - w_c) (r/r_{\max})^\beta$$



Drew & Giddings (1982)

dependence beyond the Lyman limit, down to X-ray wavelengths. In the optically thin approximation, even when the spectral index is increased to  $\alpha = 3$  in the ultraviolet and X-ray part of the spectrum, the mass-loss rate needs to be of the order of  $\sim 1000 M_\odot \text{ yr}^{-1}$ , if the visual luminosity of the QSO is comparable with that of 3C 273. It is anticipated that,

**Such mass loss is  
not achievable in AGN!**

# THEORY: INFLOWS AND OUTFLOWS IN AGN

❖ Radiation-driven disk winds: driving by resonance lines (Shlosman et al. 1985)

$$F_{l,\text{tot}} = \sum_l \frac{\mathcal{F}(R; \nu_l)}{c} k_l \rho \frac{1 - e^{-\tau_l}}{\tau_l}$$

radiation force (Sobolev) per unit volume summed over lines (C, N, O lines)

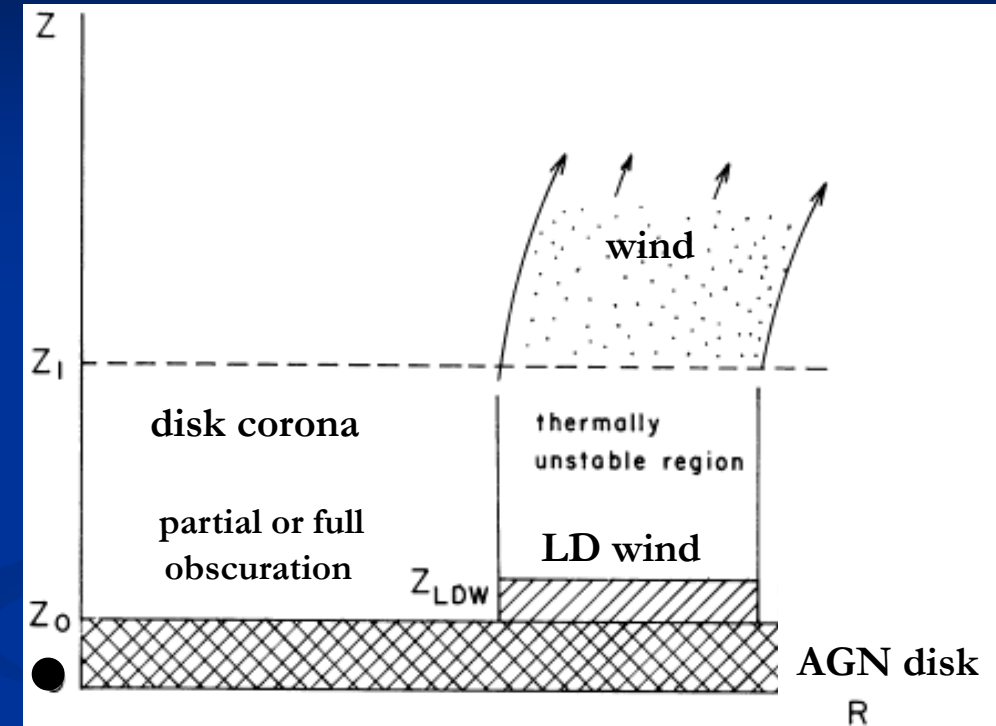
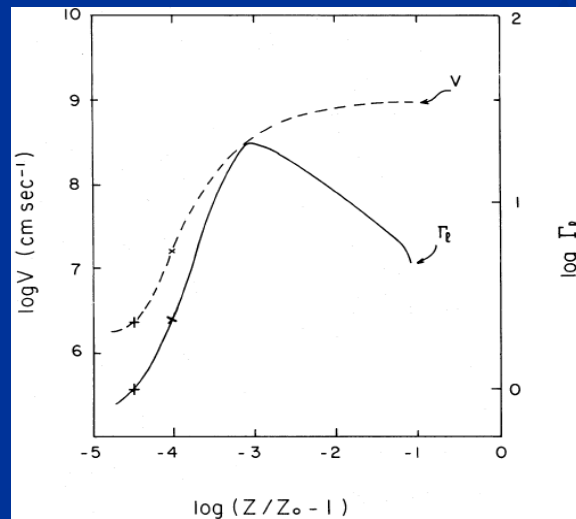
$$\tau_l = \frac{k_l \rho c}{\nu_l} \left| \frac{dv}{dz} \right|^{-1}$$

optical depth in line  $l$  depends on v-gradient!

$$\rho v_z g = \dot{M}$$

$$(v_z^2 - v_s^2) \frac{dv_z}{dz} = \frac{v_z}{z} \left[ \frac{2v_s^2 z^2}{gR_0^2} - z \frac{dv_s^2}{dz} - \frac{GM_{\text{BH}} z^2}{g^{3/2} R_0^3} (1 - \Gamma_{\text{es}} - \Gamma_l) \right]$$

$$\frac{dT}{dz} = \frac{2}{3} \frac{1}{v_z} \left[ \frac{H - \Lambda}{k_B n} - T \left( \frac{dv_z}{dz} + \frac{2zv_z}{gR_0^2} \right) \right]$$



Shlosman, Vitello & Shaviv (1985)

vertical velocity and line radiation force profiles (normalized by gravity)

**disk wind fits, But what about line profiles?**

# THEORY: INFLOWS AND OUTFLOWS IN AGN

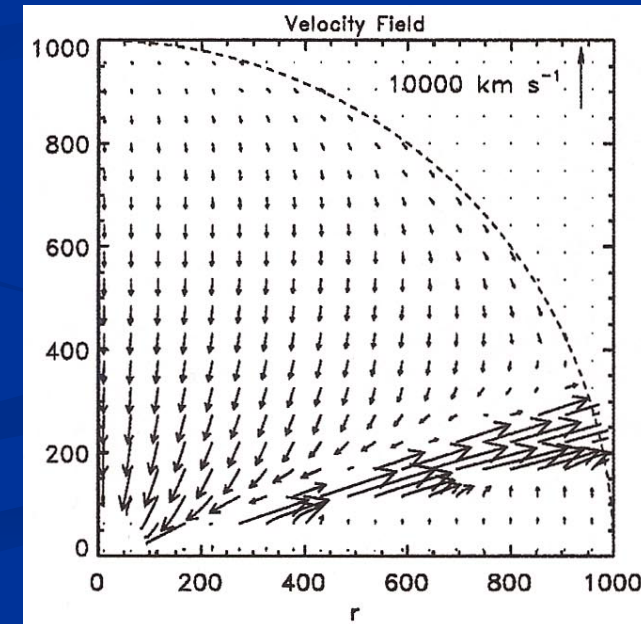
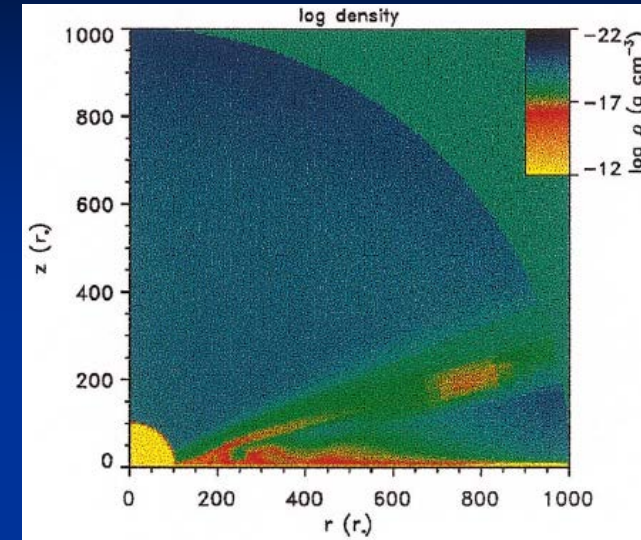
❖ Radiation-driven disk winds: driving by resonance lines (Proga et al. 2000)

$$F^{\text{rad},l}(r) = \oint_{\Omega} M(t) \left[ \hat{n} \frac{\sigma_e I(r, \hat{n}) d\Omega}{c} \right]$$

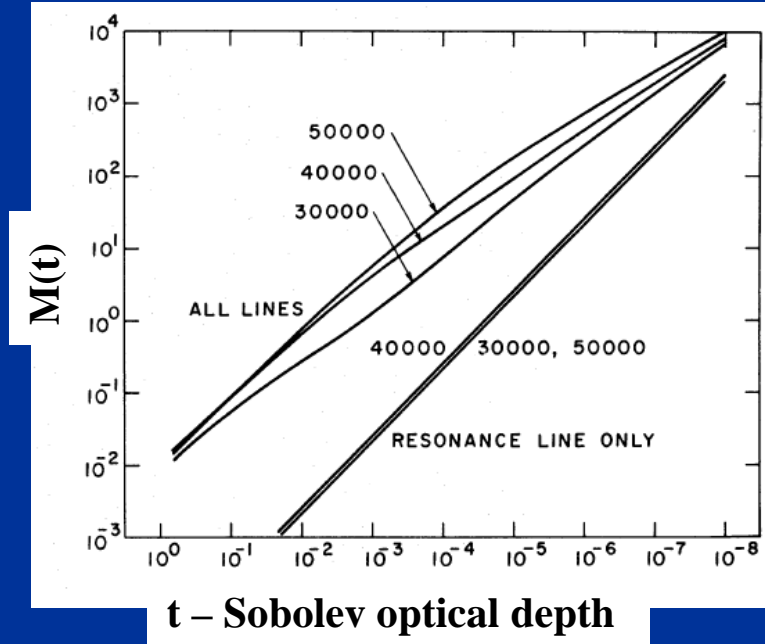
radiation force (Sobolev)  
per unit mass approximated  
by force multiplier  $M(t)$

electron scattering force

**Force multiplier has been calculated for stellar case only !**



$$M(t) = \sum_{\text{lines}} \frac{F_c \Delta v_D}{F} \min \left( \frac{1}{\beta}, \frac{1}{t} \right)$$



$t$  – Sobolev optical depth

Castor, Abbott & Klein (1973)

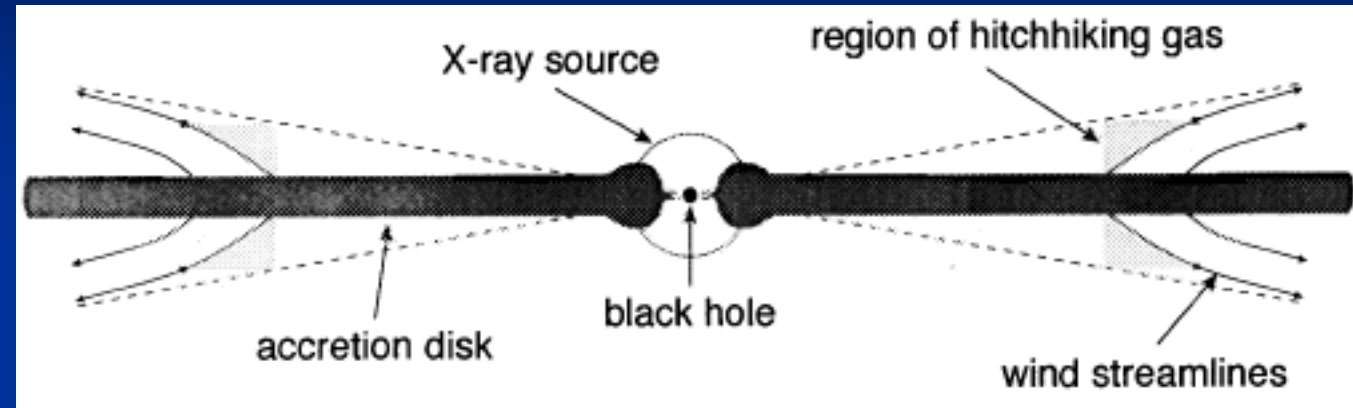
Proga et al. (2000)

# THEORY: INFLOWS AND OUTFLOWS IN AGN

❖ **Radiation-driven disk winds:** driving by resonance lines (Murray et al. 1995)

Continuous wind → no clumps

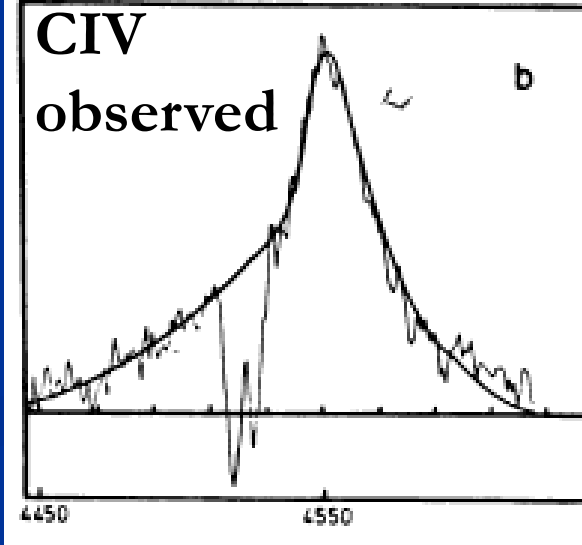
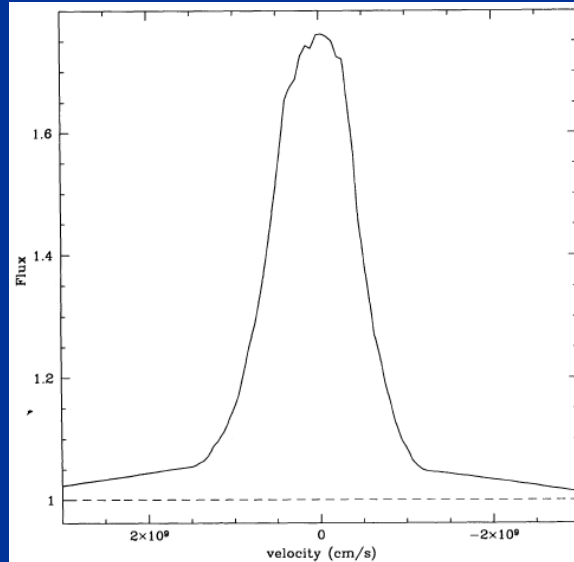
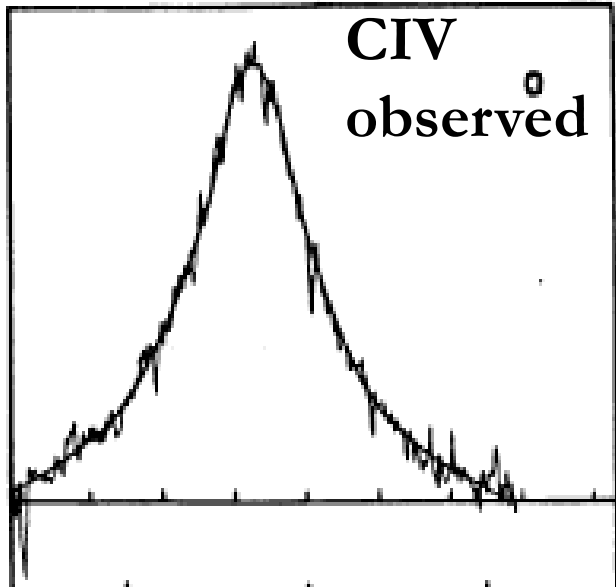
Radiation force (Sobolev)  
per unit mass approximated  
by **force multiplier  $M(t)$**   
as in Proga et al. (2000)



Wind streamlines make  $5^\circ$  with the disk surface

*B. J. Wilkes*

**CIV**  
**observed**



**CIV line shapes  
don't fit....  
geometry wrong?**



# THEORY: INFLOWS AND OUTFLOWS IN AGN

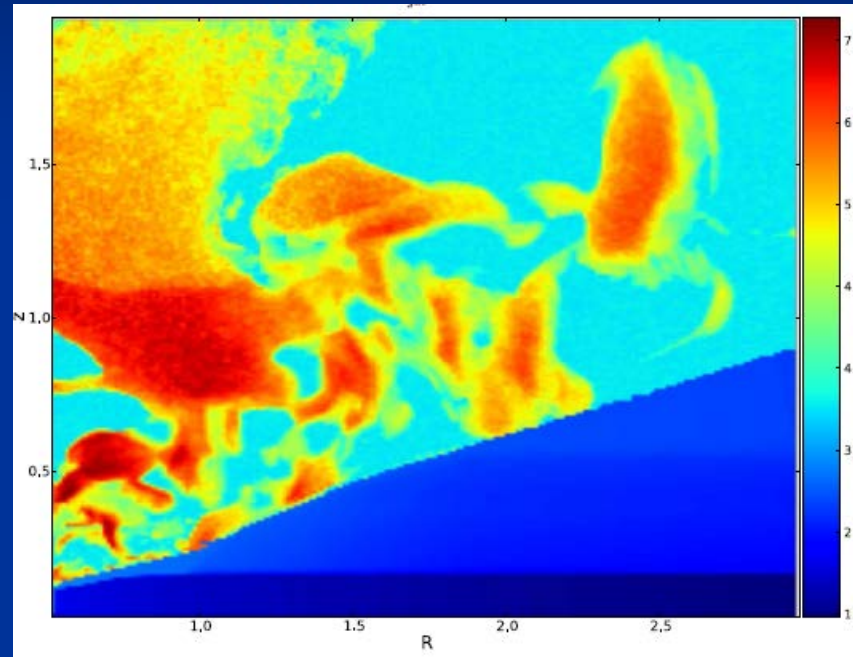
❖ Radiation-driven disk winds: driving by dust

Single fluid approximation  
→ no dust evolution

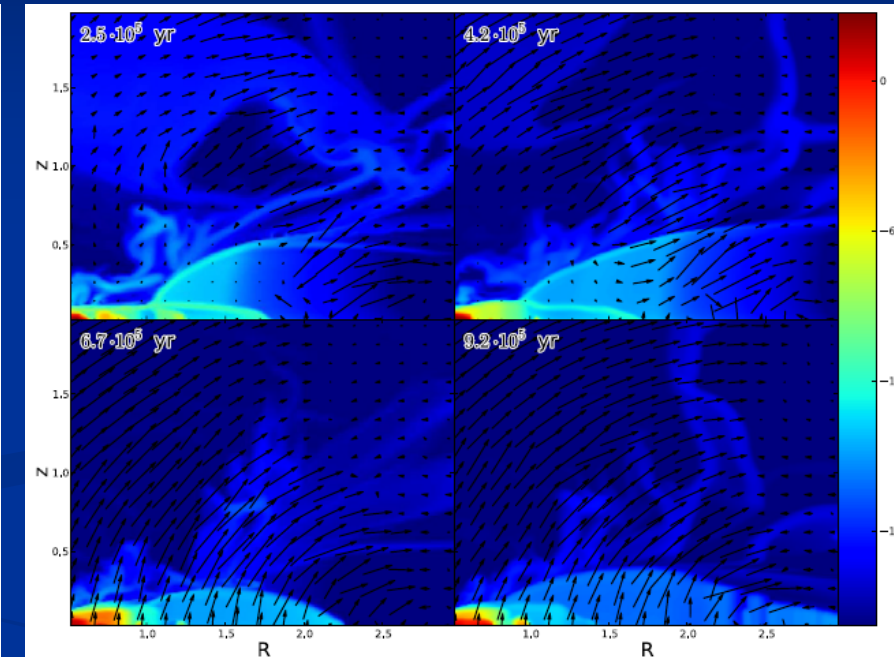
$$\mathbf{F} = -D \nabla E$$

Radiation flux calculated  
using flux-limited diffusion  
approximation

color → gas  $T$



color → gas density



Dorodnitsyn & Kallman (2012)

**But do dust grains survive both the acceleration and associated temperature? → need self-consistent treatment!**



# THEORY: INFLOWS AND OUTFLOWS IN AGN

## ❖ Theory: summary of MHD

Mass conservation

$$\nabla \cdot (\rho \mathbf{v}) = 0$$

z- component of momentum conservation (Euler eq.)

$$\rho (\mathbf{v} \cdot \nabla) v_z = -\frac{\partial p}{\partial z} - \frac{\rho \partial \Phi}{\partial z} - \frac{1}{8\pi} \frac{\partial B^2}{\partial z} + \frac{1}{4\pi} (\mathbf{B} \cdot \nabla) B_z$$

Energy

$$\rho T \left( \frac{dS}{dt} \right) = \rho T v_p \bullet \nabla S = Q$$

Perfect gas

$$P = \rho \frac{k_B}{\mu m_p} T$$

Ohm's law

$$\eta_m J_\phi e_\phi = \mathbf{v}_p \times \mathbf{B}_p$$

Induction

$$\nabla \cdot \left( \frac{\nu'_m}{r^2} \nabla r B_\phi \right) = \nabla \cdot \frac{1}{r} (B_\phi \mathbf{u}_p - B_p \Omega r)$$

# THEORY: INFLOWS AND OUTFLOWS IN AGN

❖ **Outflows:** magneto-centrifugal winds → Blandford & Payne (1982) solution ( $\rho_{\text{disk}} \sim r^{-3/2}$ )

Include inertia and assume MHD conditions  $\mathbf{E} + (1/c)\mathbf{v} \times \mathbf{B} = 0$

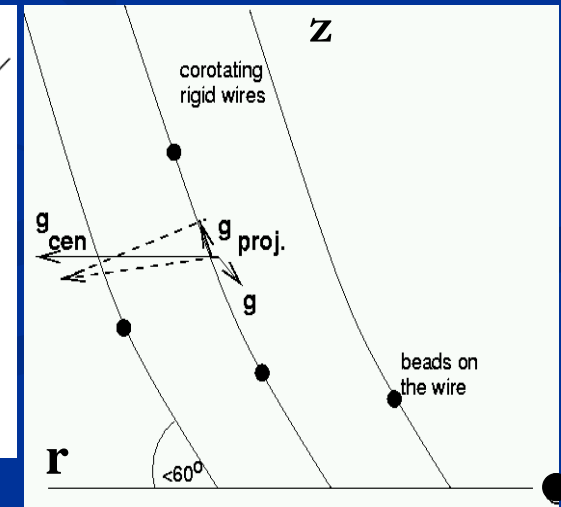
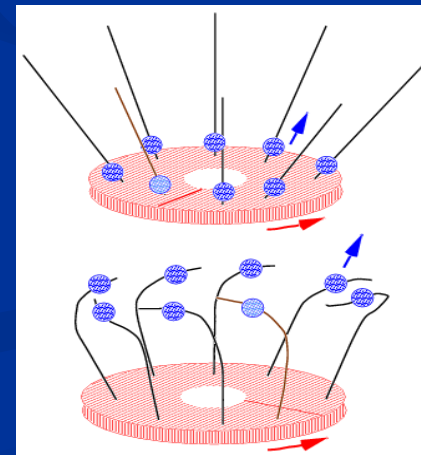
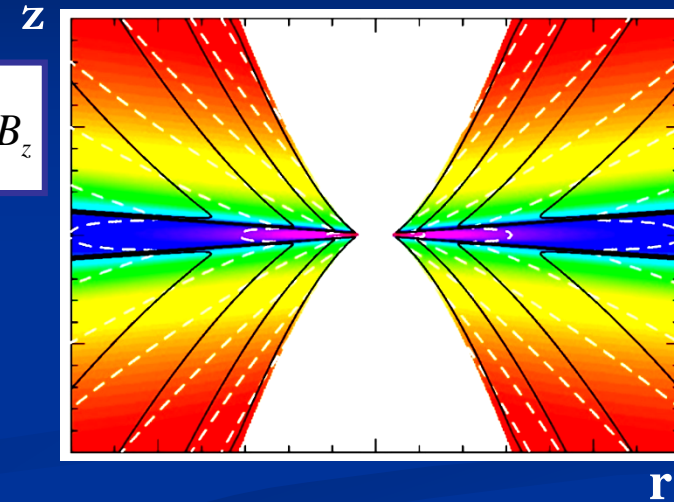
Stationary axisymmetric MHD flow  
Euler equation

$$\rho(\mathbf{v} \cdot \nabla)v_z = -\frac{\partial P}{\partial z} - \rho \frac{\partial \Phi}{\partial z} - \frac{1}{8\pi} \frac{\partial B^2}{\partial z} + \frac{1}{4\pi} (\mathbf{B} \cdot \nabla)B_z$$

Self-similar solution  $\mathbf{r} = [r_0 \xi(\chi), \phi, r_0 \chi]$   $\mathbf{v} = [\xi'(\chi) f(\chi), g(\chi), f(\chi)] (GM / r_0)^{1/2}$

Solutions scale with spherical radius along a given direction

Centrifugal acceleration (gas clouds on B-lines act as “beads on a wire”): a wind is launched when the inclination angle of magnetic lines to the disk is  $< 60^\circ$   
After launch the flow is dominated by the toroidal magnetic field imposed by rotation  
Collimation along the magnetic axis



# THEORY: INFLOWS AND OUTFLOWS IN AGN

❖ **Outflows:** magneto-centrifugal winds → Blandford & Payne (1982) solution ( $\rho_{\text{disk}} \sim r^{-3/2}$ )

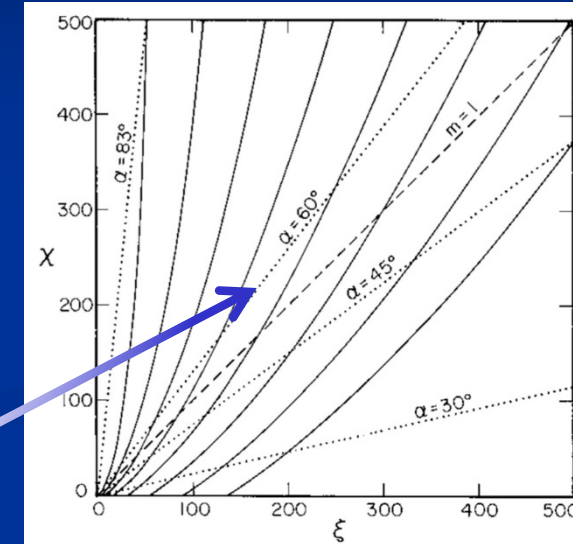
Disk winds: 2D → remove one degree of freedom → 1D ordinary differential equation

Self-similarity assumptions:

- cylindrical  $z$ : pre-supposes a collimated vertical jet structure
- cylindrical  $r$ : accretion disk structure, not jets
- spherical  $\theta$ : spherical wind (NO collimation)
- spherical  $r$ : only choice with equations that allow collimation

→ Blandford & Payne (1982):  $r$ -self-similarity;  $\theta$  structure same for every field line

reduces MHD to only two ordinary differential equations



$$\mathbf{r} = [r_0 \xi(\chi), \phi, r_0 \chi]$$

# THEORY: INFLOWS AND OUTFLOWS IN AGN

- ❖ **Outflows:** magneto-centrifugal winds → merging radiation & MHD winds  
generalized Blandford/Payne model (Emmering, Blandford & Shlosman 1992 solution)

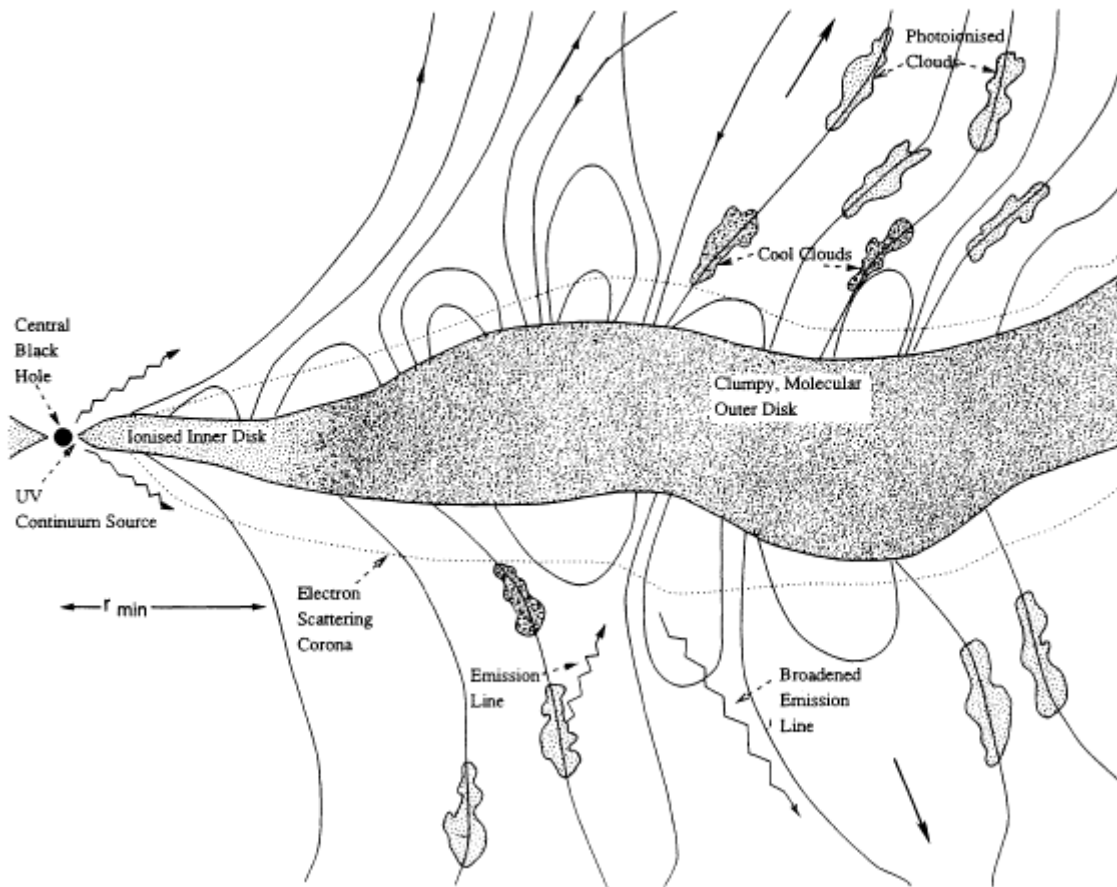
$$\rho_{\text{disk}} \sim r^{-\alpha}$$

$$\dot{M} r_0 v_0 \sim K B_A^2 r_A^3$$

**B-field extracts angular momentum → low  $\dot{M}$  wind !**

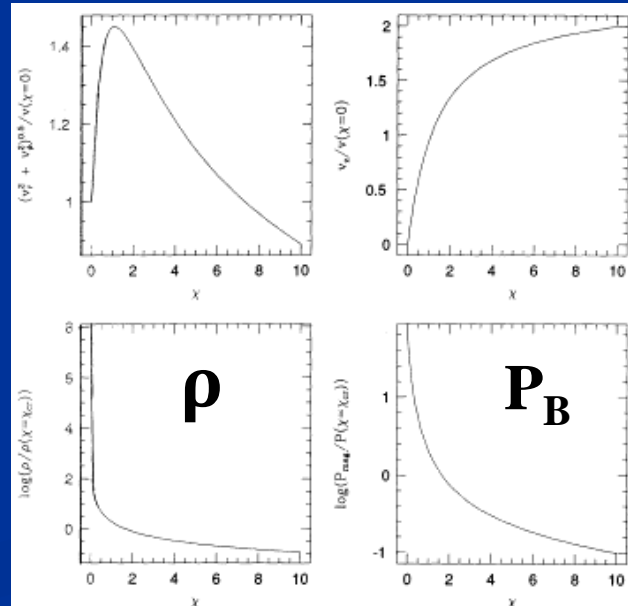
**Electron-scattering corona widens the line wings**

EMMERING, BLANDFORD, & SHLOSMAN



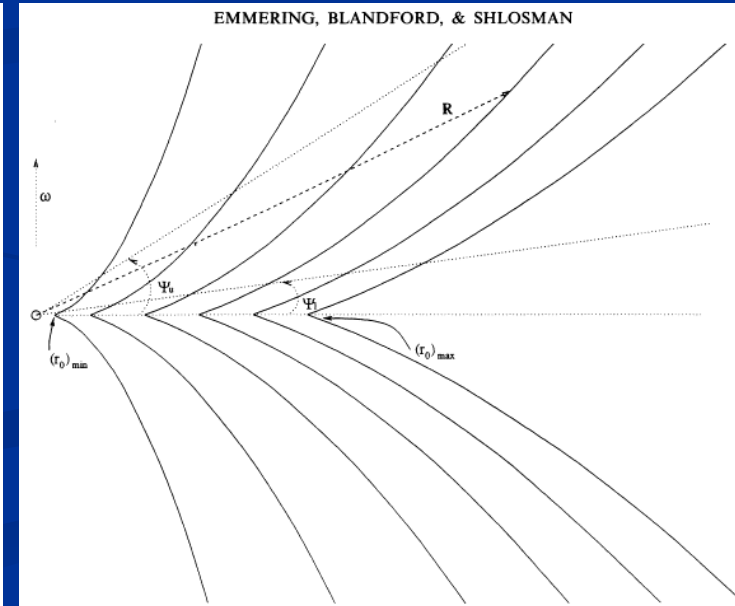
$v \perp z$  axis

$v \parallel z$  axis



**self-similarity of B-lines and  $r$**

EMMERING, BLANDFORD, & SHLOSMAN



# THEORY: INFLOWS AND OUTFLOWS IN AGN

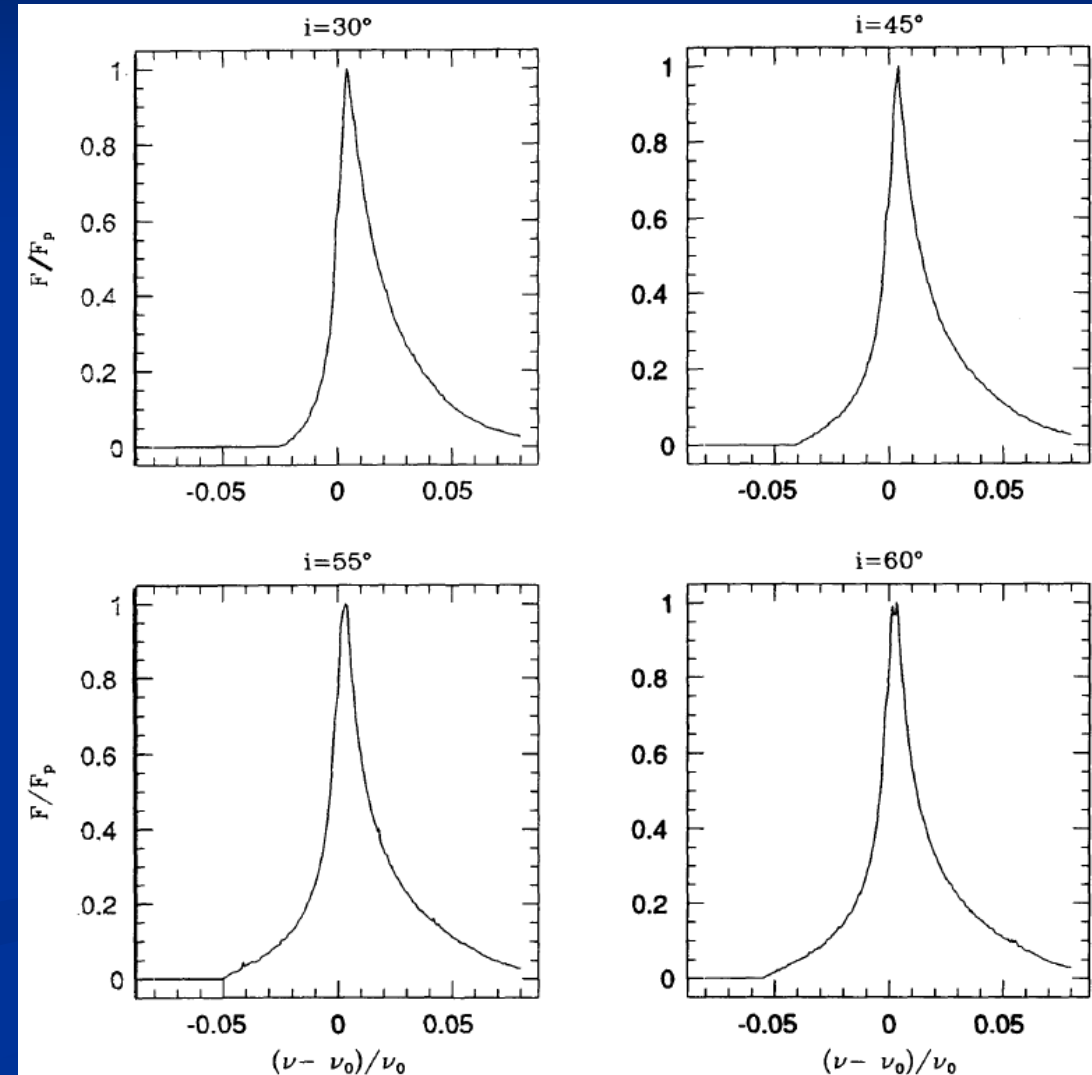
❖ **Outflows:** magneto-centrifugal winds → merging radiation & MHD winds

Generalized Blandford/Payne model (Emmering, Blandford & Shlosman 1992 solution)

So, disk winds produce characteristic triangular shapes of emission lines in AGN → match as good as for spherical outflows!

model CIV lines

CIV emission line profiles for various inclinations to z-axis



CIV emission line profiles for various gas emissivities

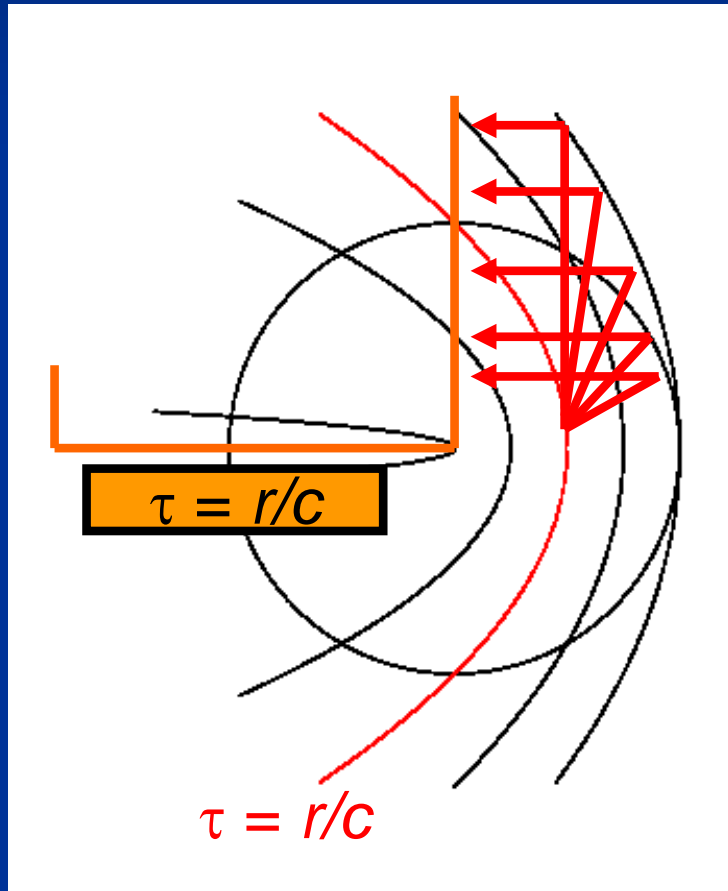
CIV emission line profiles with observed asymmetry



# THEORY: INFLOWS AND OUTFLOWS IN AGN

❖ **Outflows:** magneto-centrifugal winds → reverberation mapping of BLR

Formation of the broad line region (BLR)



isodelay surfaces

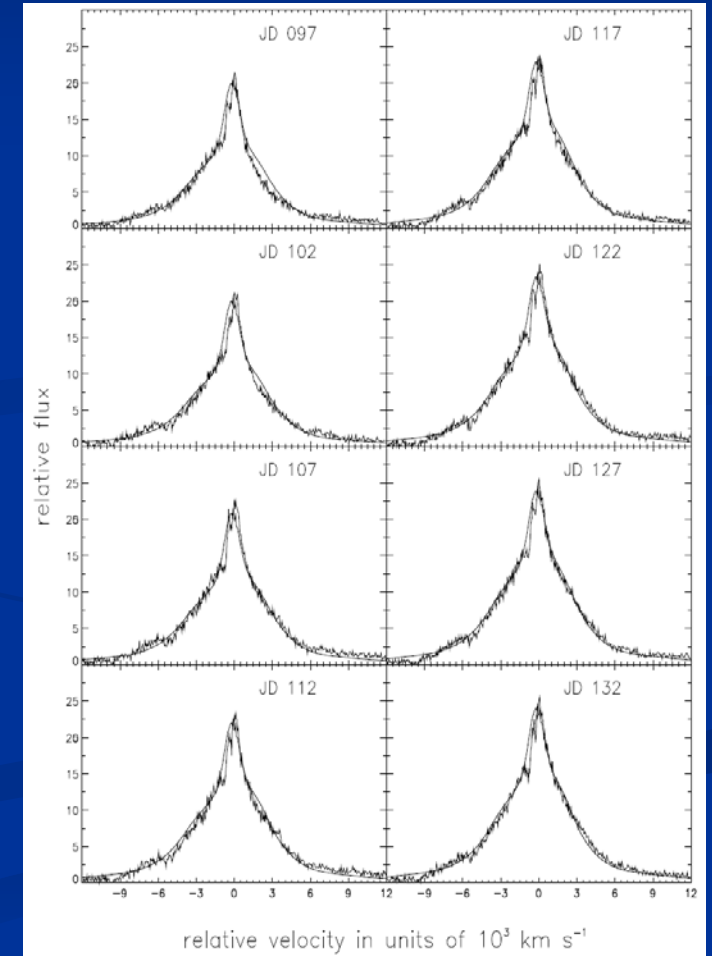
**echo mapping**

All points on an “isodelay surface” have the same extra light-travel time to the observer, relative to photons from the continuum source.

**$M_{\bullet} \sim 3 \times 10^7 M_{\odot}$**

in Seyfert 1 galaxy NGC5548

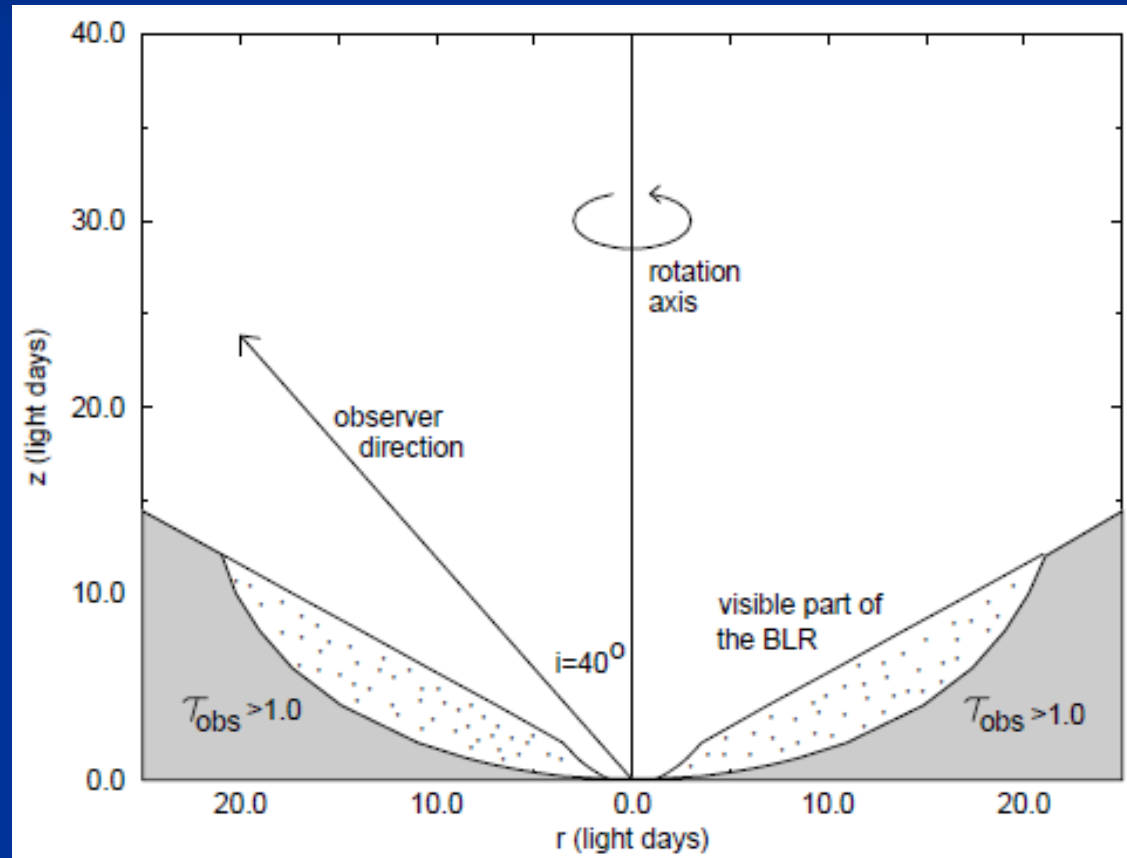
C IV emission line profile evolution in Seyfert 1 galaxy NGC5548



# THEORY: INFLOWS AND OUTFLOWS IN AGN

❖ **Outflows:** magneto-centrifugal winds

Formation of the broad line region (BLR)



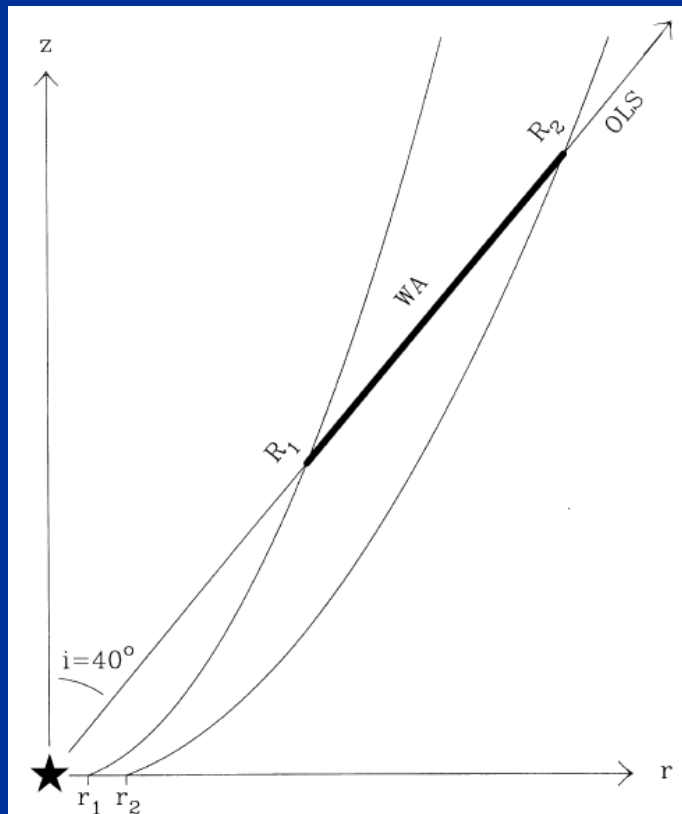
Bottorff, Korista, Shlosman & Blandford (1997)

# THEORY: INFLOWS AND OUTFLOWS IN AGN

❖ **Outflows:** magneto-centrifugal winds

Multi-component warm absorbers in NGC 5548

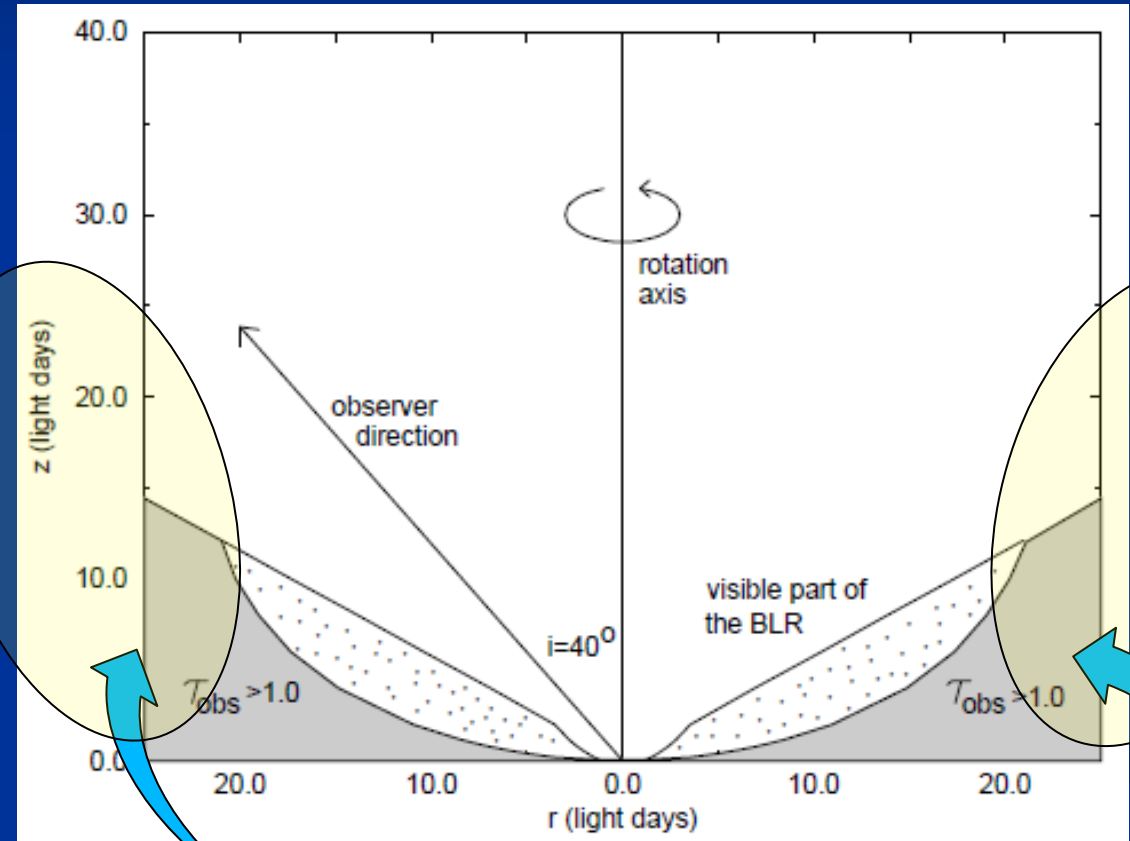
observer's orientation in NGC 5548



MODEL SOLUTION FOR THE CONTINUOUSLY DISTRIBUTED WARM ABSORBER IN NGC 5548

Parameter	log(value) <sup>a</sup>
$R_{\min}$ .....	17.75
$R_{\max}$ .....	20.59
$n_{\min}$ .....	6.455
$n_{\max}$ .....	2.253
$N(\text{H})$ .....	21.76
$N(\text{H I})$ .....	15.94
$N(\text{C IV})$ .....	14.70
$N(\text{N V})$ .....	15.13
$N(\text{O VII})$ .....	17.94
$N(\text{O VIII})$ .....	18.22
$\epsilon$ .....	-2.716

<sup>a</sup> Values in cgs units.



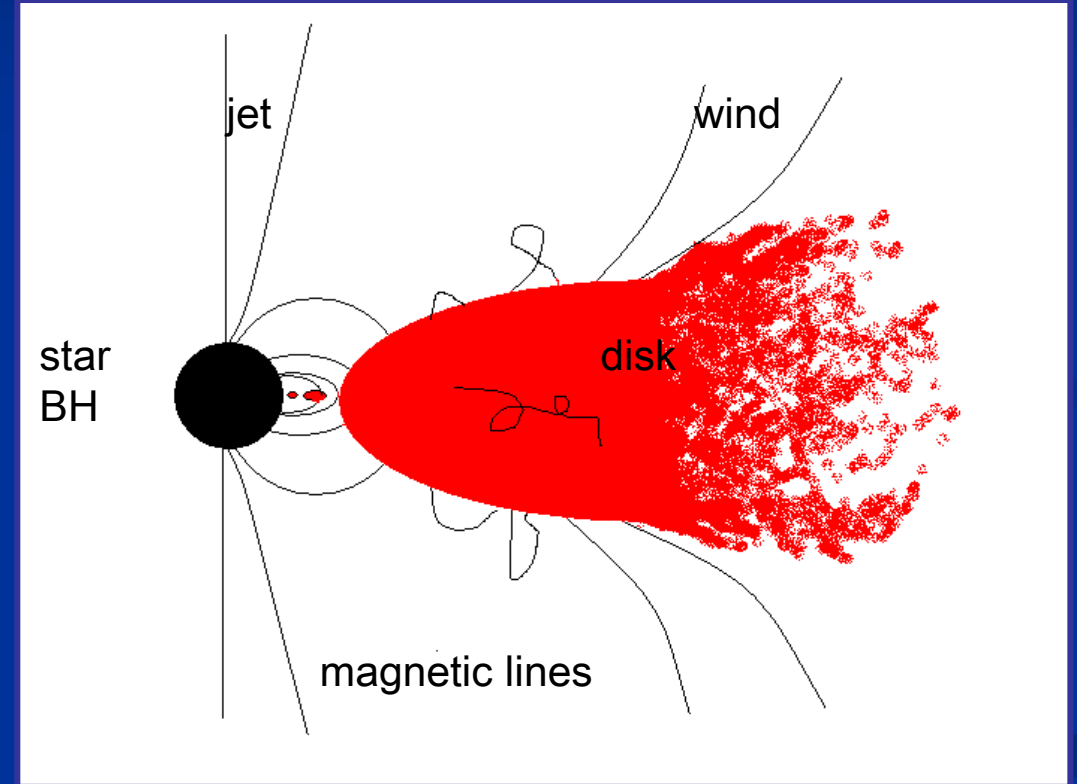
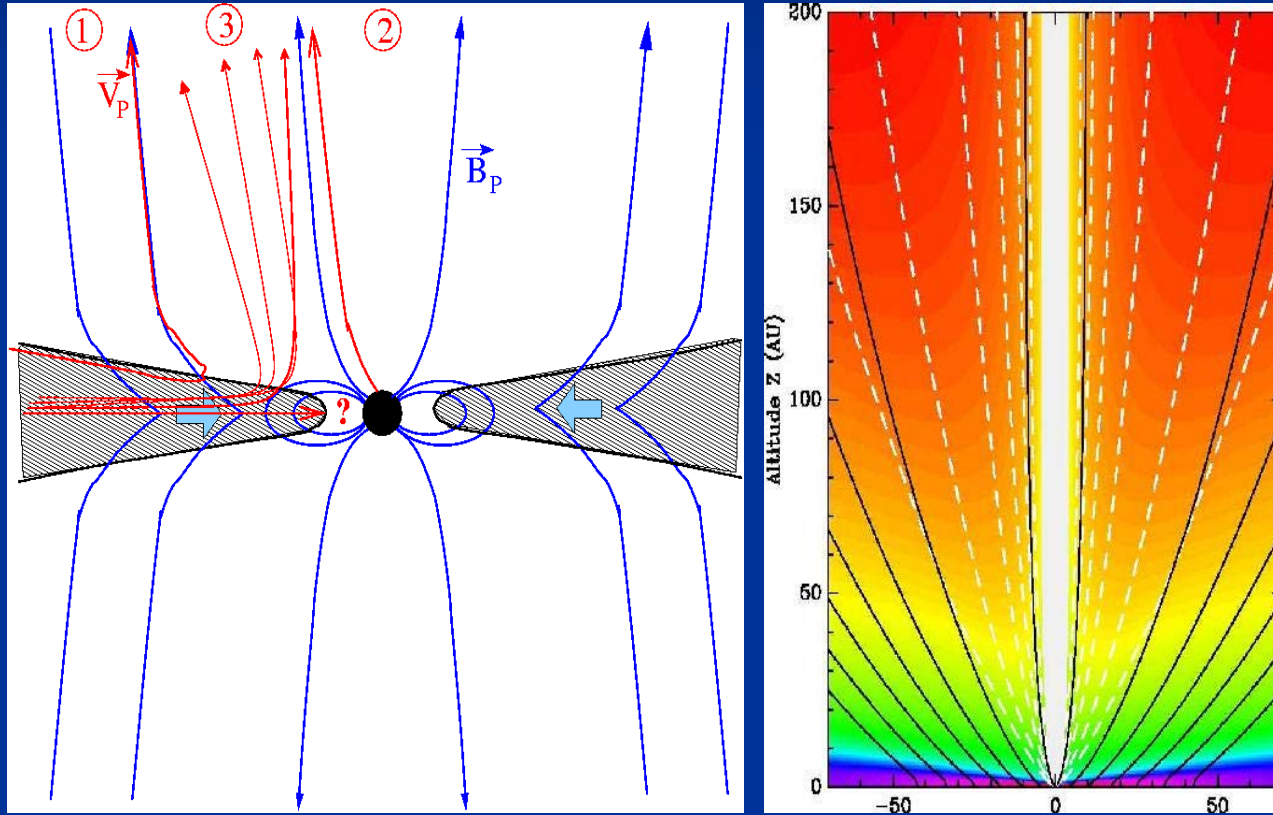
warm absorbers



# THEORY: INFLOWS AND OUTFLOWS IN AGN

❖ **Outflows:** collimated MHD winds  $\rightarrow$  jets

Accretion disk-driven jets



Accretion disk driven jets  $\rightarrow$  velocity distribution  
at the wind base is that of a Keplerian disk  
Blandford & Payne model: inertia  $\rightarrow$  poloidal B  $\rightarrow$  toroidal B

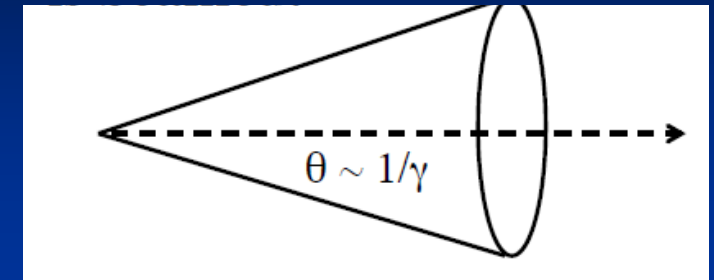


# THEORY: INFLOWS AND OUTFLOWS IN AGN

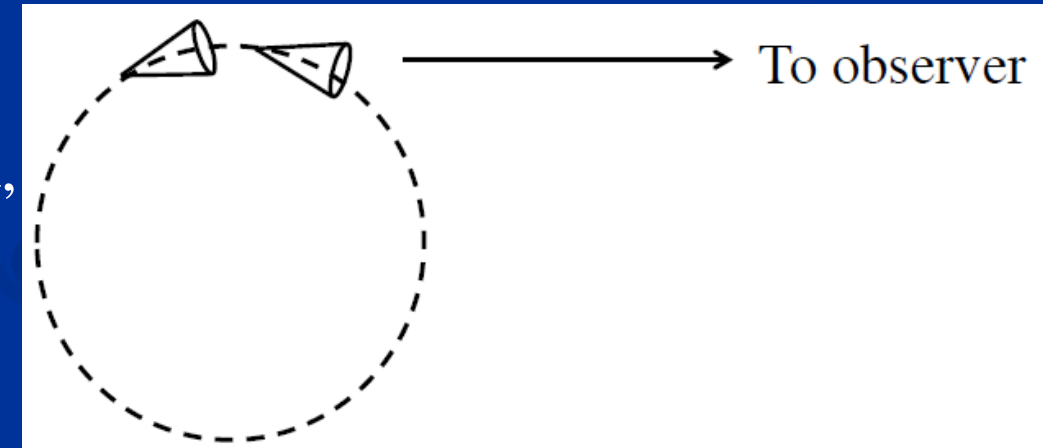
❖ **Outflows:** emission mechanism in jets  $\rightarrow$  synchrotron radiation

If electrons are moving at  $v \sim c \rightarrow$  radiation is beamed

Particle moving with Lorentz factor  $\gamma$  toward observer emits  
Into cone of opening angle  $\theta \sim \gamma^{-1}$



We only see radiation from a small portion of the orbit,  
when the cone points toward us  
 $\rightarrow$  but many electrons!



# THEORY: INFLOWS AND OUTFLOWS IN AGN

❖ **Outflows:** collimated MHD winds → jets

Alternative acceleration mechanisms:

Twin-exhaust scheme

(Blandford & Rees 1972)

Radiation pressure in accretion funnels

(FRT 1985)

Electrodynamic effects in accretion funnels and

Poynting flux jets

(Lovelace 1976, Blandford 1976)

Magneto-centrifugal acceleration

(Blandford & Payne 1982)

# THEORY: INFLOWS AND OUTFLOWS IN AGN

❖ **Outflows:** collimated MHD winds → jets

Alternative acceleration mechanisms: tapping the rotational energy of black hole

Blandford & Znajek (1977)

Blandford and Znajek(1977) found a *stationary solution* for monopole magnetospheres of slowly rotating black holes. It exhibited outflows of energy and angular momentum

Black hole rotational energy ( $a = 1$ ):  $E_b = 0.29 M_b c^2 \simeq 10^{54} \left( \frac{M_b}{2M_\odot} \right) \text{ erg}$

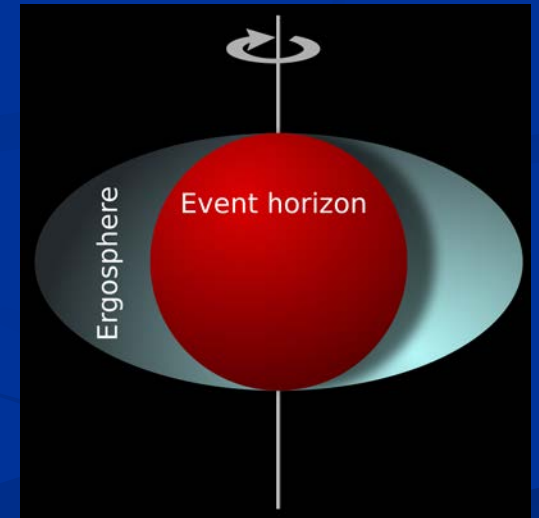
Power of the Blandford-Znajek mechanism:  $L_{BZ} \simeq 3.6 \times 10^{50} a^2 \left( \frac{M}{2M_\odot} \right)^2 \left( \frac{\Psi}{10^{27}} \right)^2 \text{ erg/s}$

$a$  - spin parameter of the black hole ( $0 < a < 1$ ),

$\Psi$  - the magnetic flux of black hole.

$\Psi = 10^{27} \text{ G cm}^2$  is the highest value observed in magnetic stars:

Ap, white dwarfs, neutron stars (magnetars).



Kerr black hole

**Efficiency of Blandford-Znajek mechanism ?**

# THEORY: INFLOWS AND OUTFLOWS IN AGN

❖ **Outflows:** collimated MHD winds → jets

Tapping the rotational energy of black hole: Blandford & Znajek (1977)

What is the condition for activation of the BZ-mechanism with finite inertia of plasma?

MHD waves must be able to escape from the black hole ergosphere !?

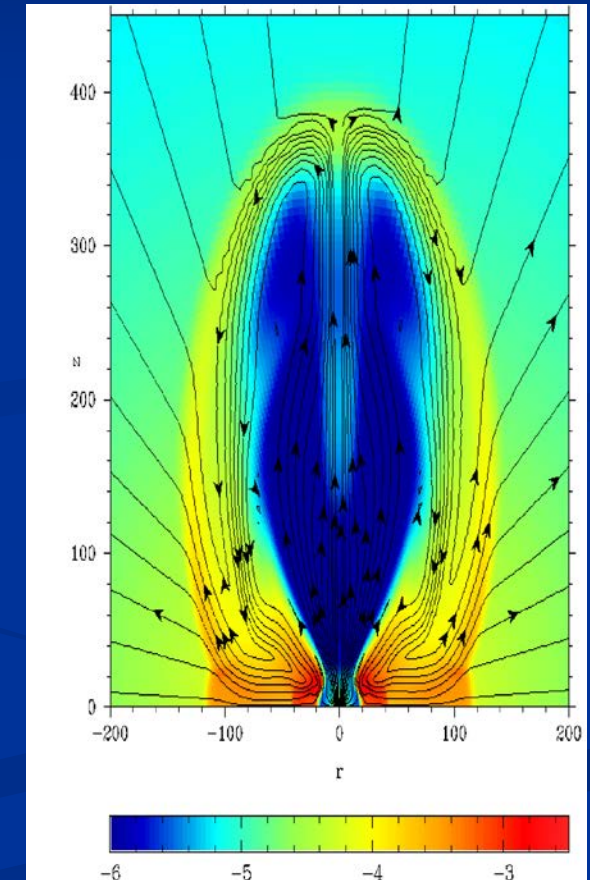
Alfven speed  $v_a > v_{ff}$  free fall

$$c_a^2 = B^2 / 4\pi\rho,$$

$$v_{ff}^2 = 2GM/r$$

Apply at the ergosphere,  $r = 2r_s = 2GM/c^2$  →  $B^2 > 4\pi\rho c$

**The energy density of magnetic field must exceed that of matter for the BZ-mechanism to be activated!**



# OBSERVATIONAL EVIDENCE FOR INFLOWS AND OUTFLOWS IN AGN

Evidence for inflows:

To be discussed on Thursday  
as a **FUELING** issue



# CONCLUSIONS FOR LECTURE 3

- ❖ **Active Galactic Nuclei (AGN)** are powered by accretion processes, but there are clear and objective difficulties to detect this accretion flow
- ❖ On the other hand, UV and some X-ray emission and absorption lines point to powerful and diverse outflows from the accretion disks in AGN
- ❖ There is a clear preference, both observationally and theoretically to the presence of accretion disks in AGN, as opposite to spherical outflows
- ❖ MHD winds have preference over radiation-driven winds in AGN, because they are capable of extracting angular momentum, which radiation is inefficient in this process
- ❖ MHD is probably collimates some of the wind into powerful jets, sometimes relativistic