

# PHYSICS OF ACTIVE GALACTIC NUCLEI

## 活動銀河核の物理学

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**and**  
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# LECTURE 5

## FORMATION AND EVOLUTION OF AGN

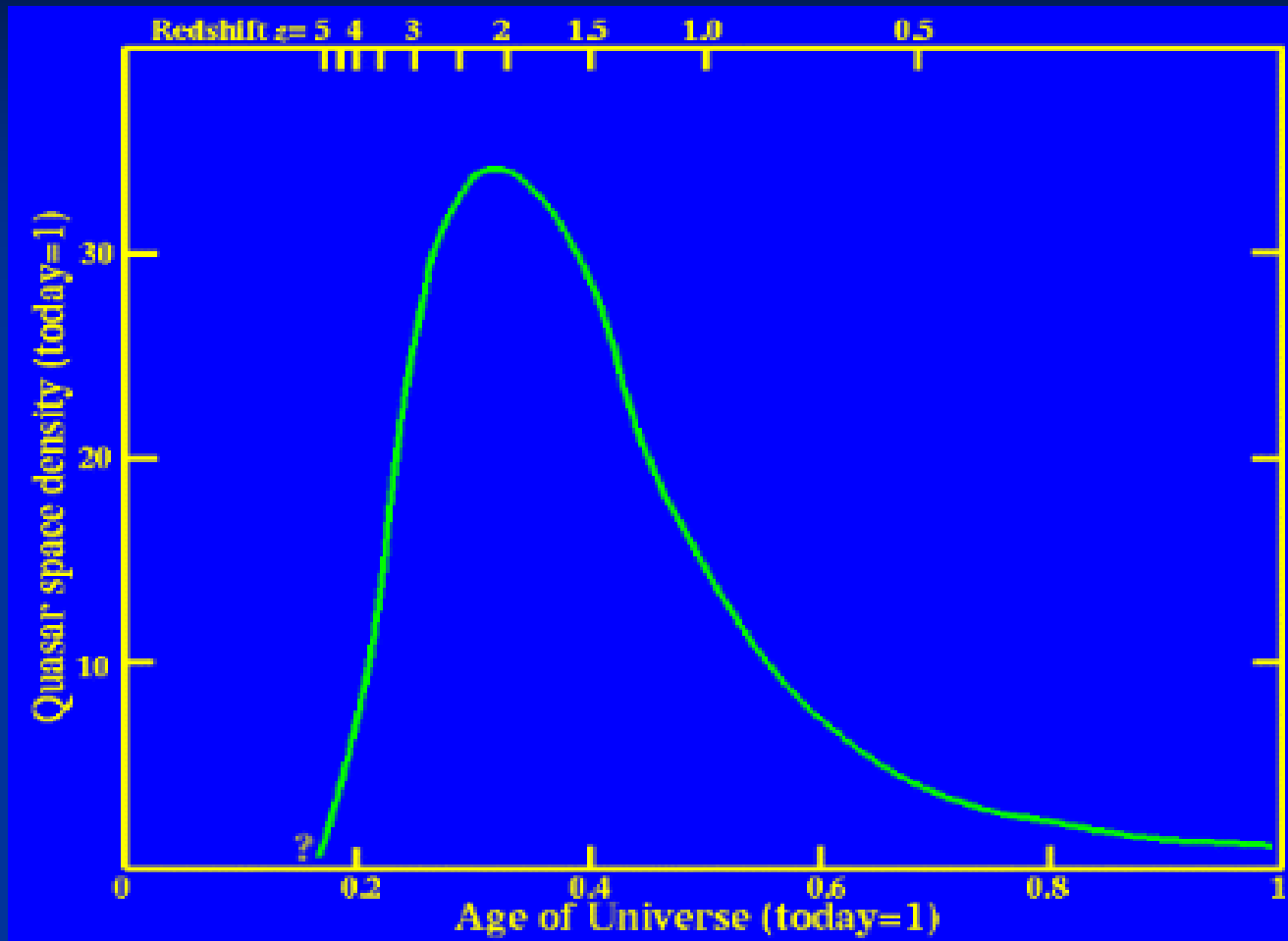
AGN across the Hubble time

Coevolution of SMBHs and galaxies

Population III stars and High-Redshift QSOs

The Making of the Black Hole Seeds: Alternatives

# AGN ACROSS THE HUBBLE TIME

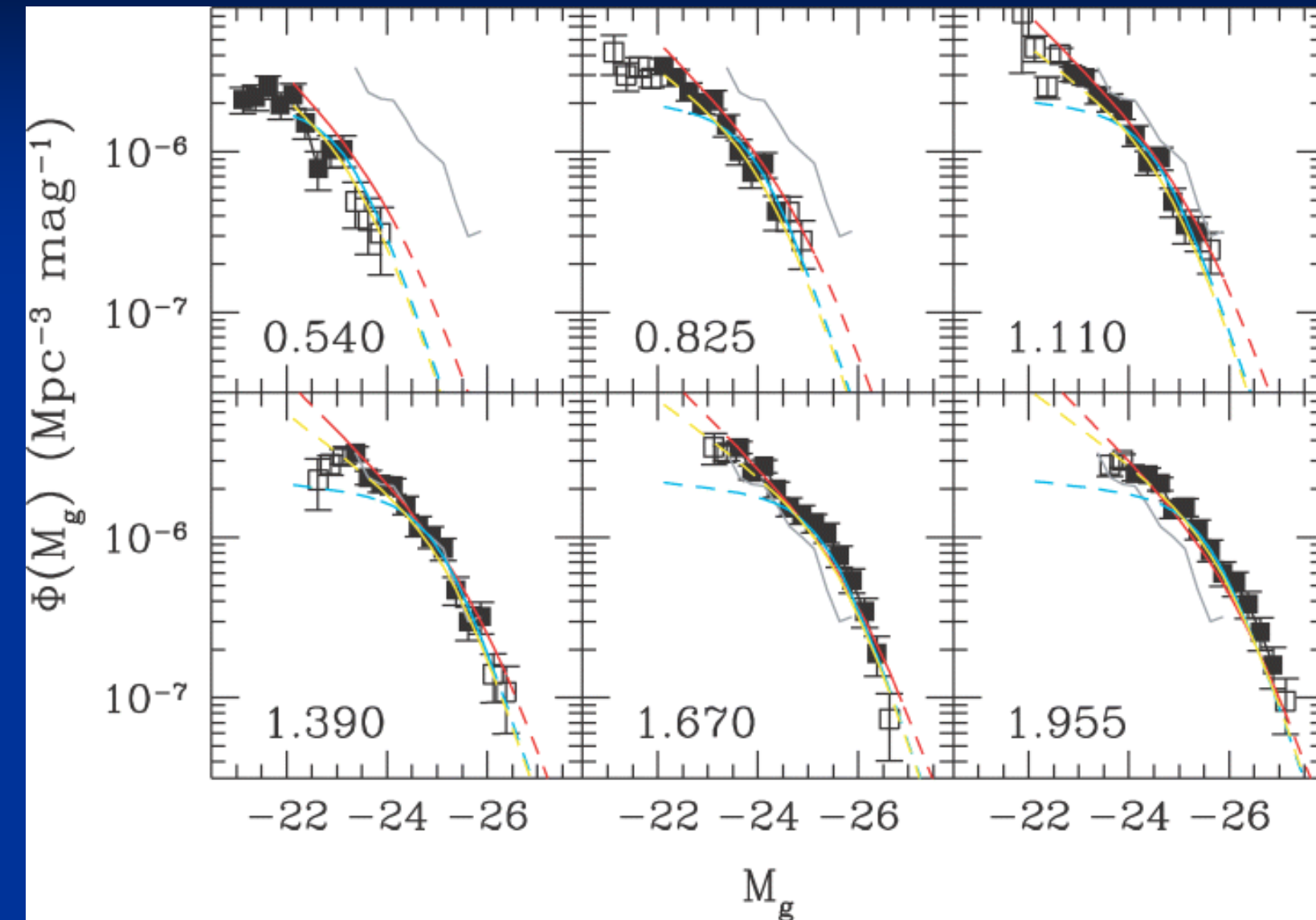


The quasar redshift distribution peaks around  $z \sim 2$

**Why around  $z \sim 2$ ?**

**This could be related to the formation of galaxies and the star-formation history**

# AGN ACROSS THE HUBBLE TIME



Richards et al (2005)

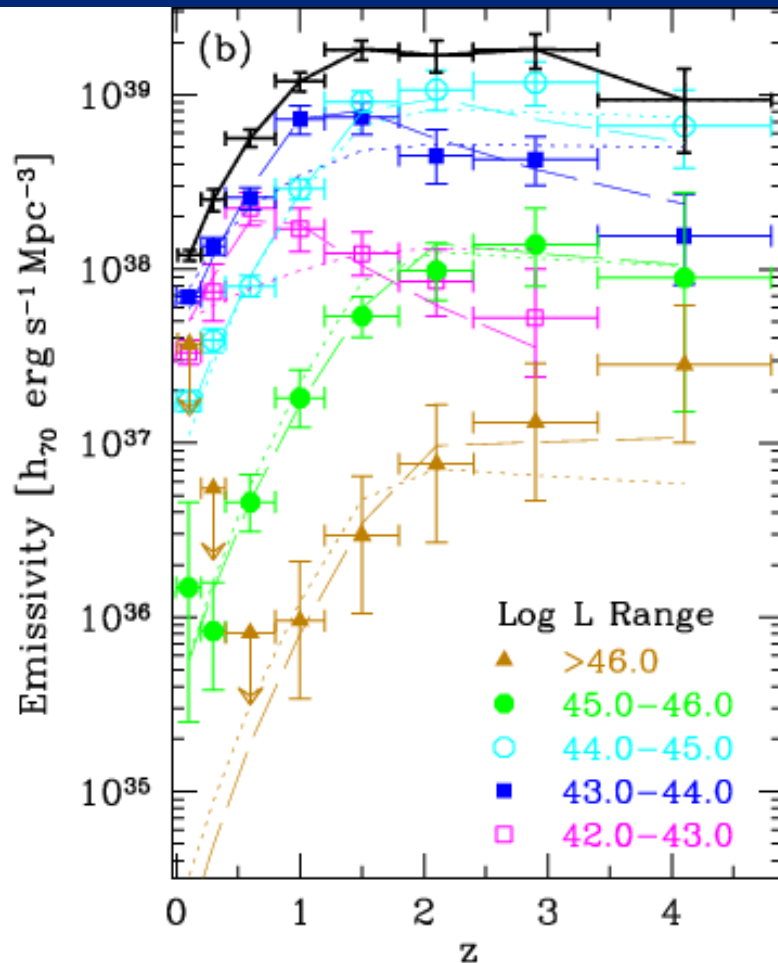
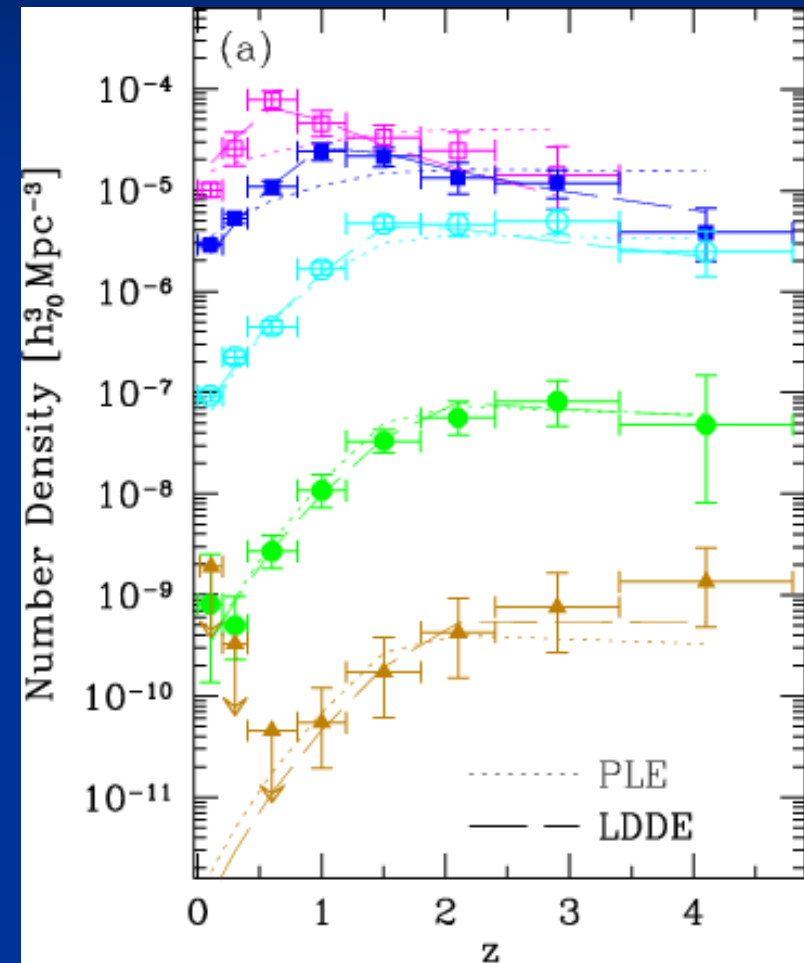
Quasar luminosity function  
(QSO LF) for various redshifts

Quasars become brighter at  
higher redshifts

# AGN ACROSS THE HUBBLE TIME

Number density

Luminosity density



Objects with lower luminosity peak at lower redshift, similar to what is observed for star formation rate (SFR) in galaxies

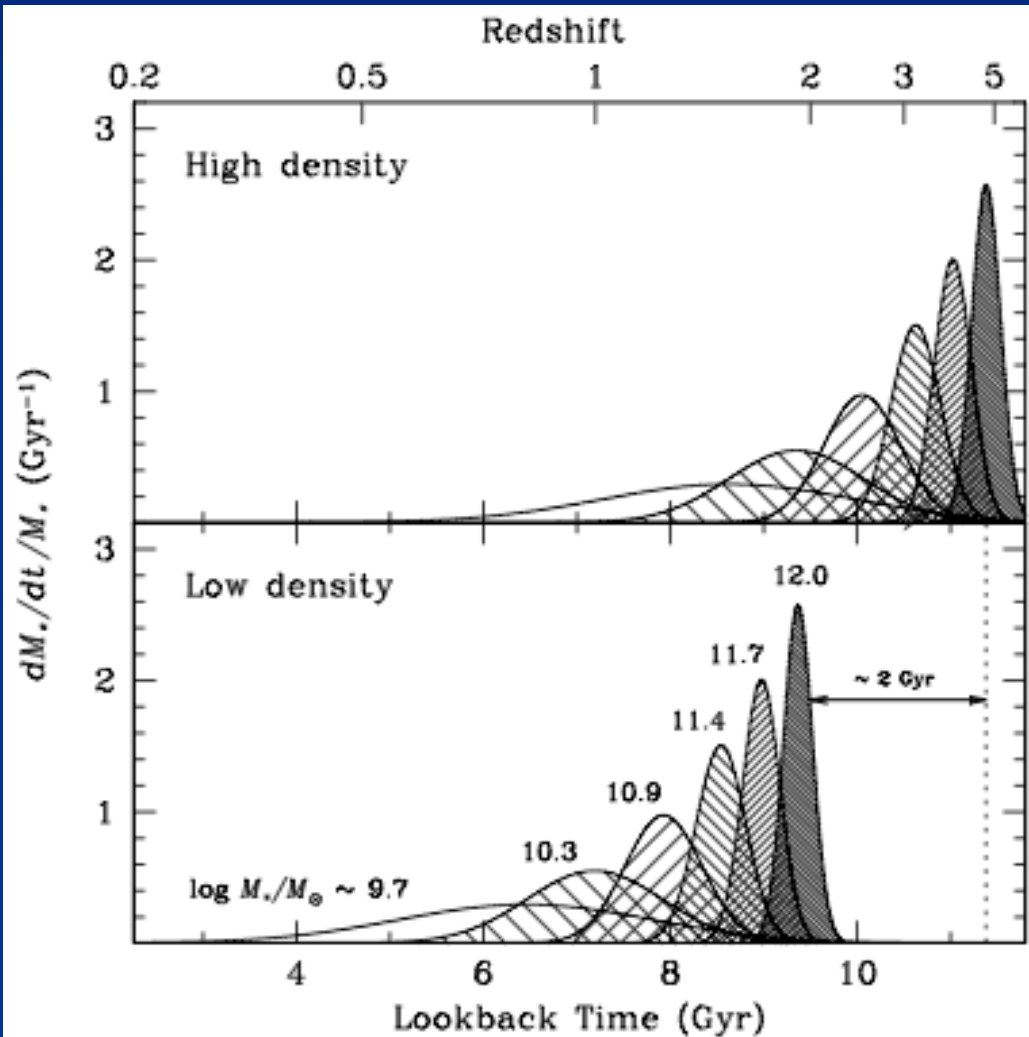
→ cosmic downsizing

QSOs peak at  $z \approx 2-3$

Rest of AGN at  $z \approx 0.5-1$

# AGN ACROSS THE HUBBLE TIME

The number density of the most luminous AGN peaks earlier in cosmic time than for less luminous objects  $\rightarrow$  large BHs are formed earlier than their low-mass counterparts?

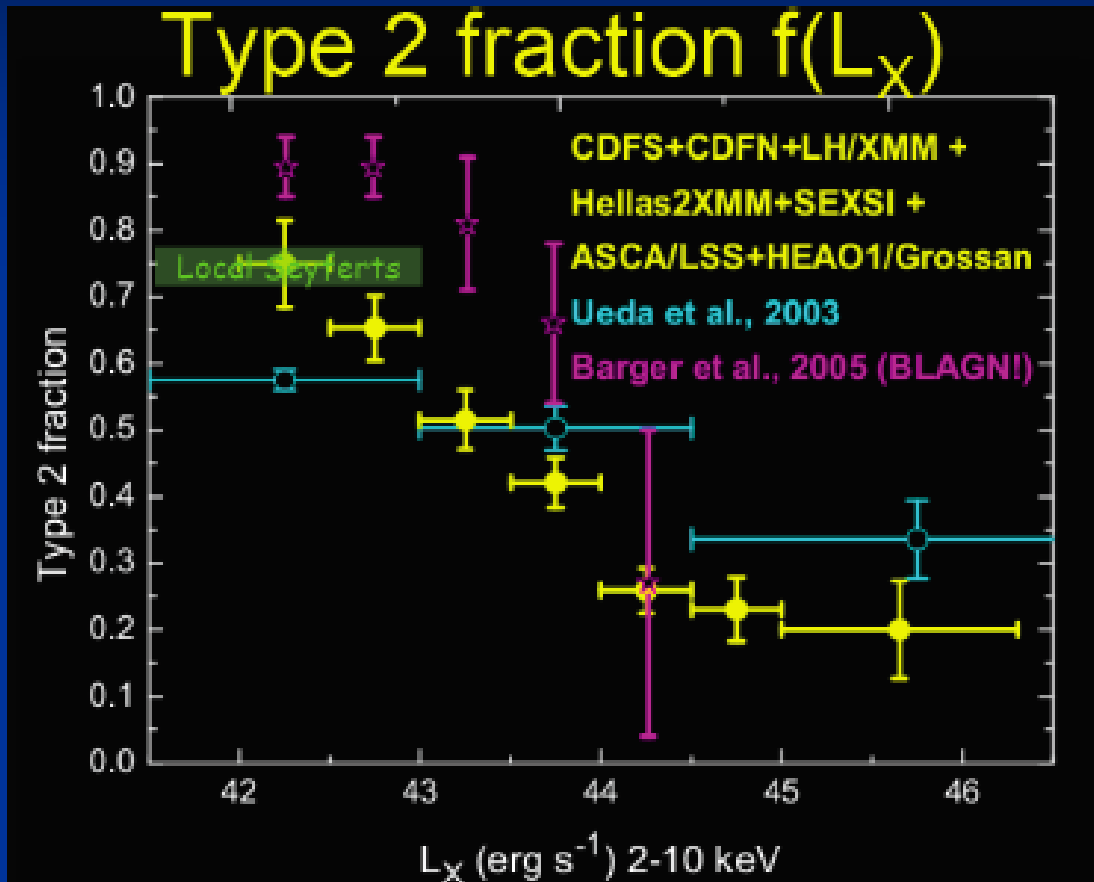


Similar behavior for galaxies:  
massive galaxies tend to form stars earlier and  
faster than less massive galaxies

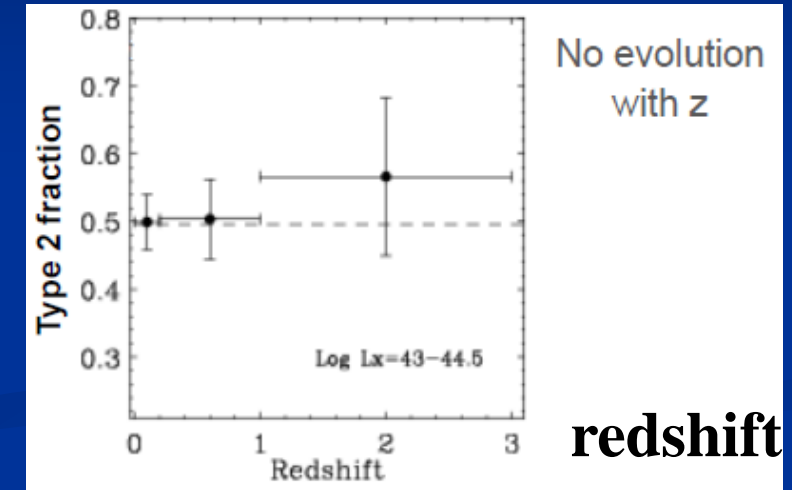
downsizing  
(Cowie et al. 1996)

# AGN ACROSS THE HUBBLE TIME

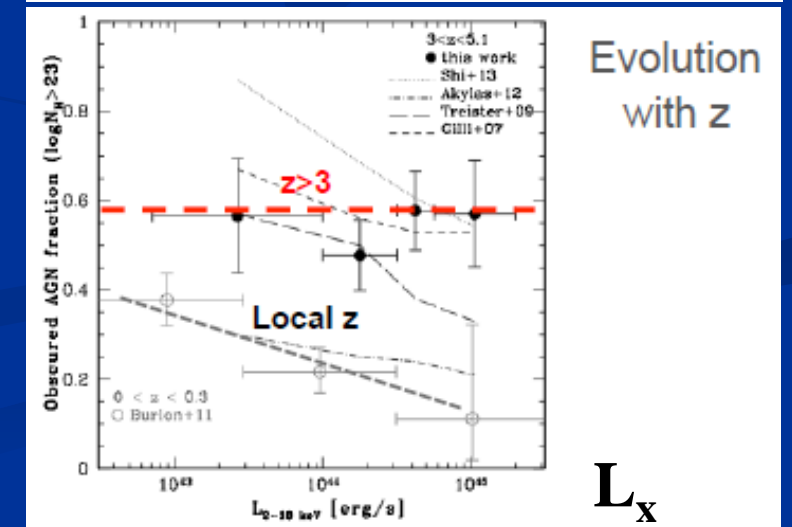
❖ Dependence of obscured AGN fraction on X-ray luminosity and redshift



Evolution of type 2 AGN with  $z$ : not clear yet



redshift

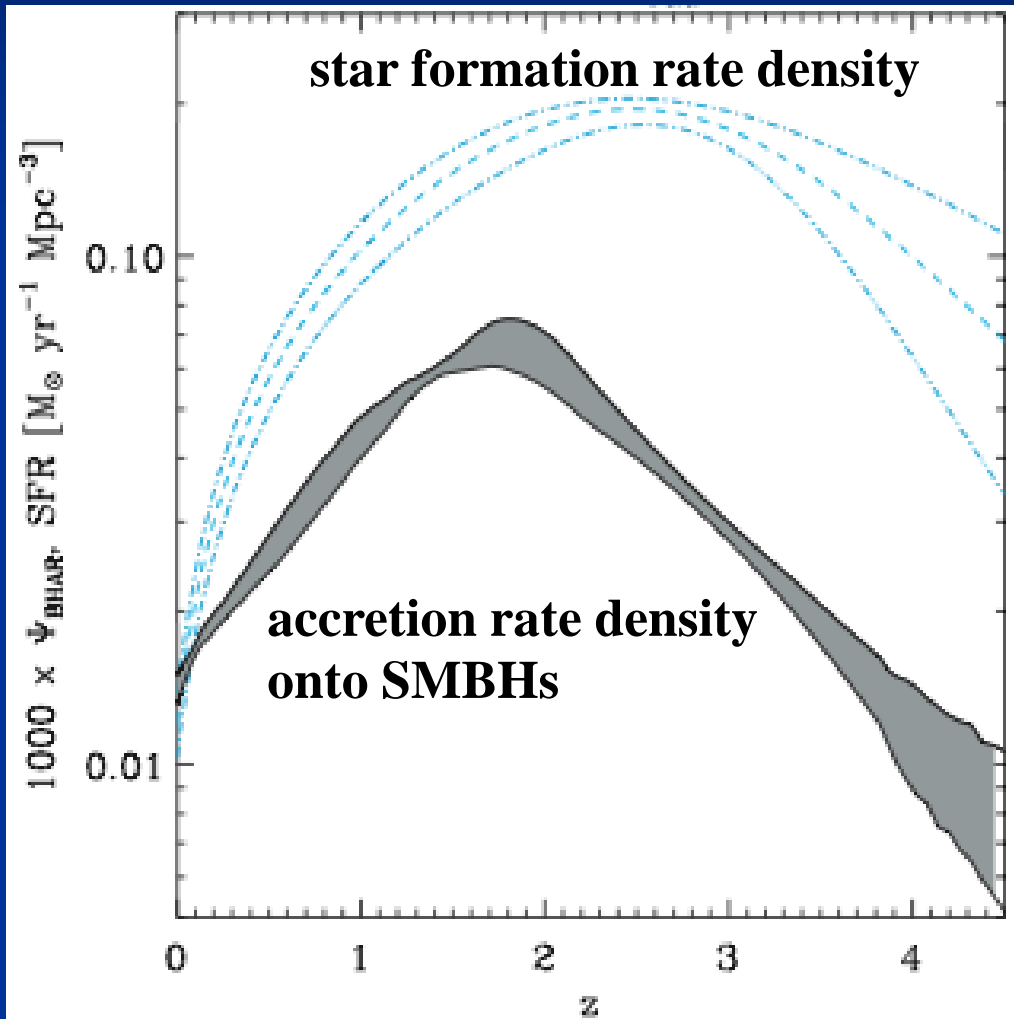


$L_x$

**broad consensus:** obscured fraction of AGN declines toward high intrinsic luminosities  
 → larger opening angle for toroidal obscuration?

# AGN ACROSS THE HUBBLE TIME

## ❖ Host galaxy and SMBH growth



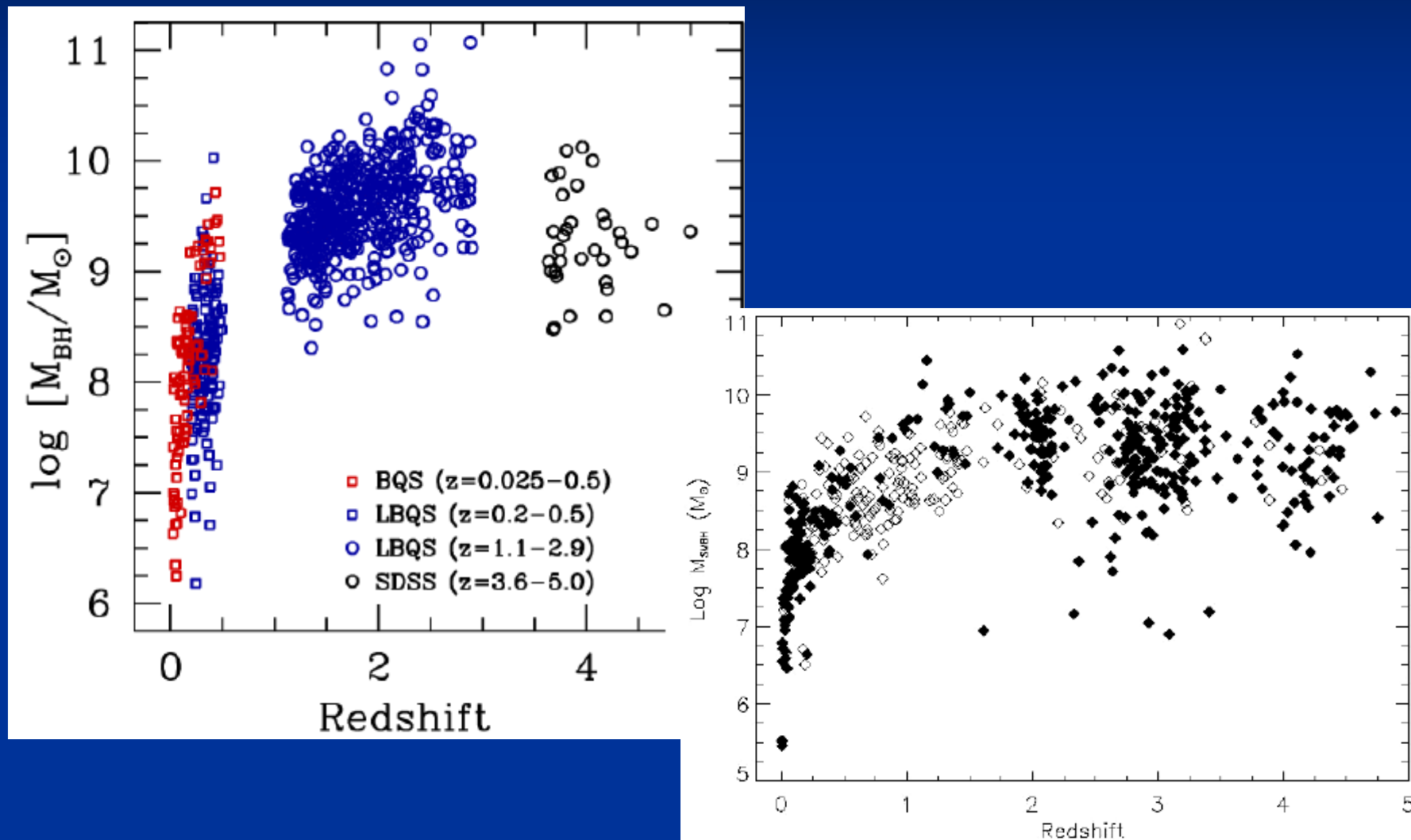
Merloni & Heinz (2008)

Star formation rate density in the universe  
and accretion rate density onto SMBHs



# AGN ACROSS THE HUBBLE TIME

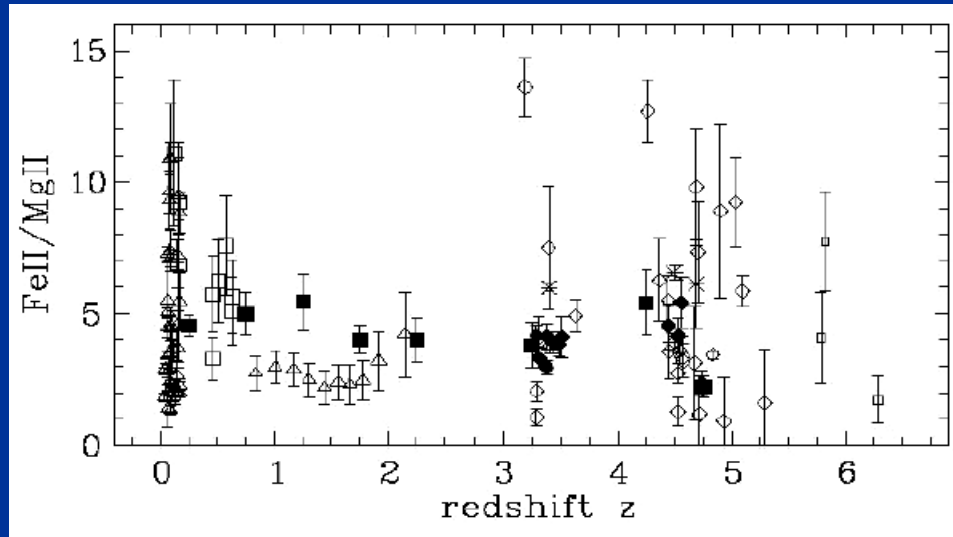
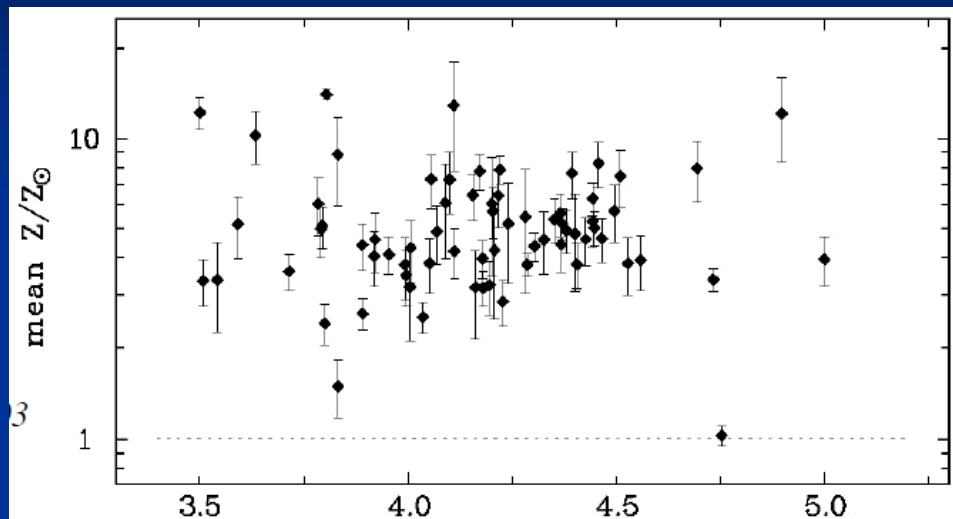
## ❖ Masses of SMBHs in distant quasars



SMBHs with masses of up to  $10^{10} M_{\odot}$  built up quickly, already by  $z \sim 5-6$

# AGN ACROSS THE HUBBLE TIME

## ❖ High-z quasar metallicity



Hamman et al. 2003; Dietrich et al. 2003

Spectra of quasars at all times show very similar metal abundances  $\leftarrow$  most heavy metals come from SNe  
Are all quasars in sites of intense SF and SNe?  
Could quasars have triggered this?  
We start to look earlier than the age of Type I SN  
 $\rightarrow$  should see Fe decline!

High-z quasars: abundance patterns (enhanced  $\text{Fe}/\alpha$ ) similar to Ellipticals

$\rightarrow$  enrichment by Type 1 SNe for  $z < 6$ ,  
with onset of star formation at  $z > 10$

**High-z quasars are very metal rich**

# AGN ACROSS THE HUBBLE TIME

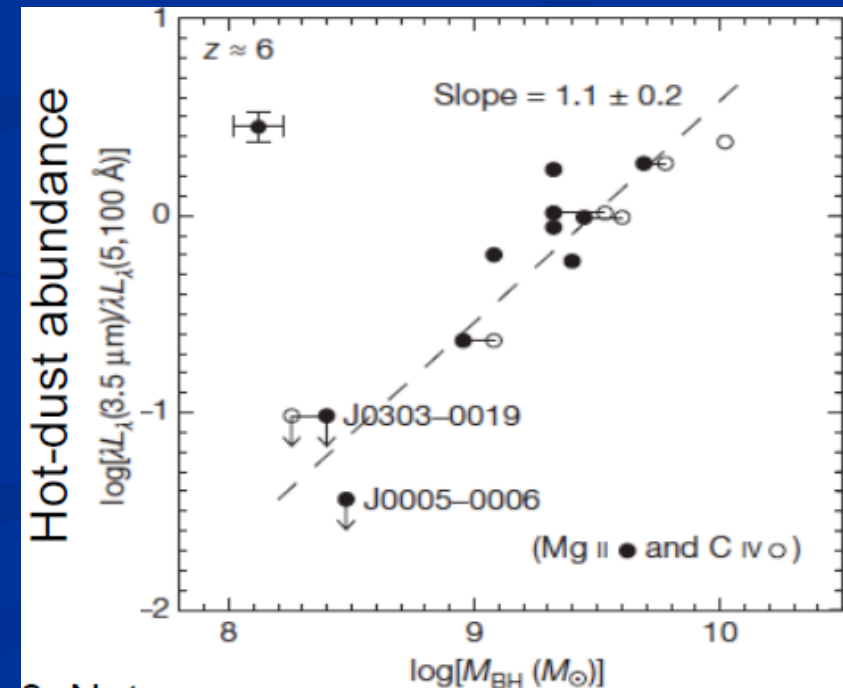
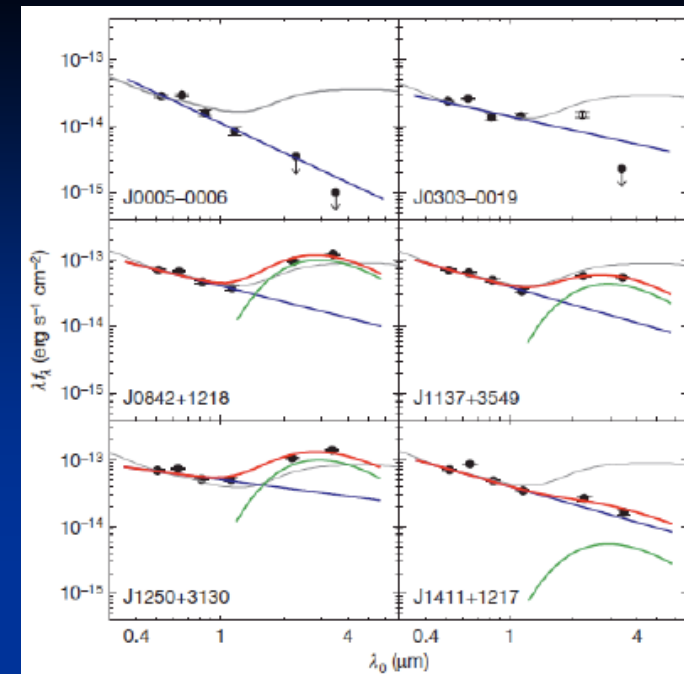
## ❖ High-z ( $\sim 6$ ) quasar dust

Is dust ubiquitous in  $z \sim 6$  quasars?

2/21  $z > 6$  quasars w/o hot dust: are quasars born in dust-free environment or not enough time to produce dust?

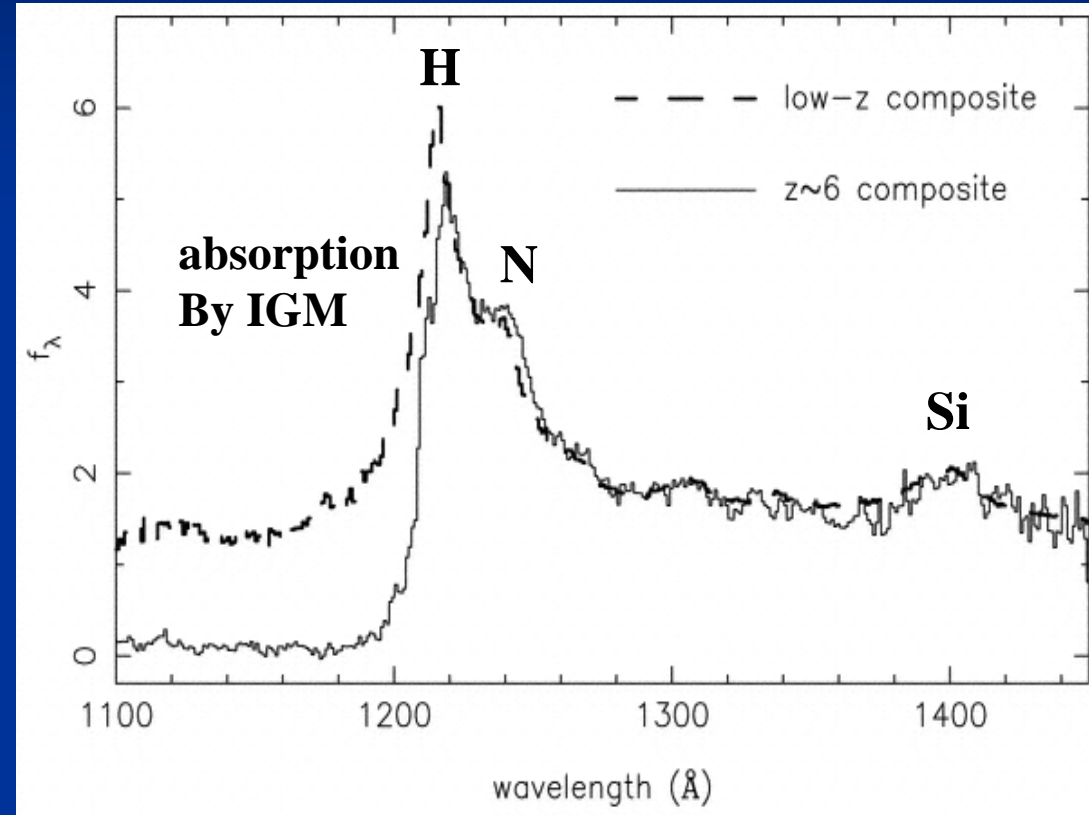
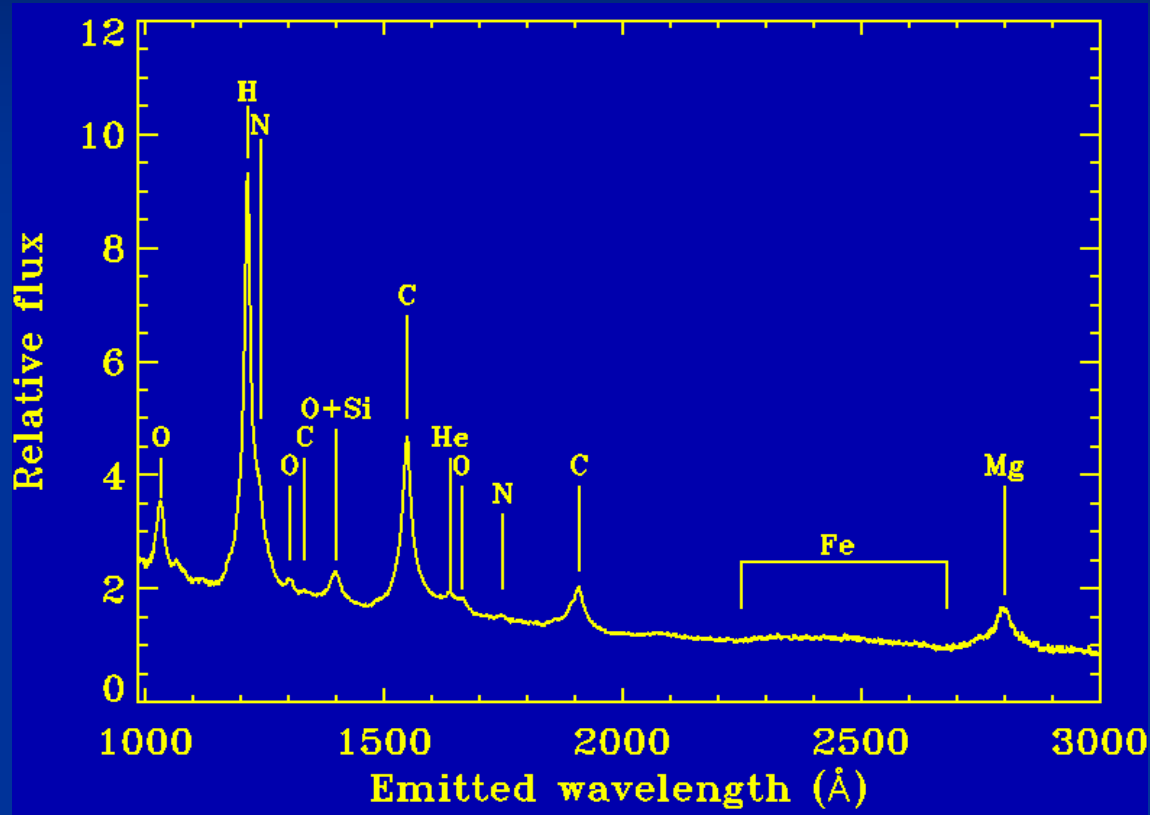
Similar findings for Type 1 AGN in C-COSMOS (Chandra+HST) survey (Hao et al. 2010)

→ from 6% at  $z < 2$  to 20% at  $z \sim 2-3.5$



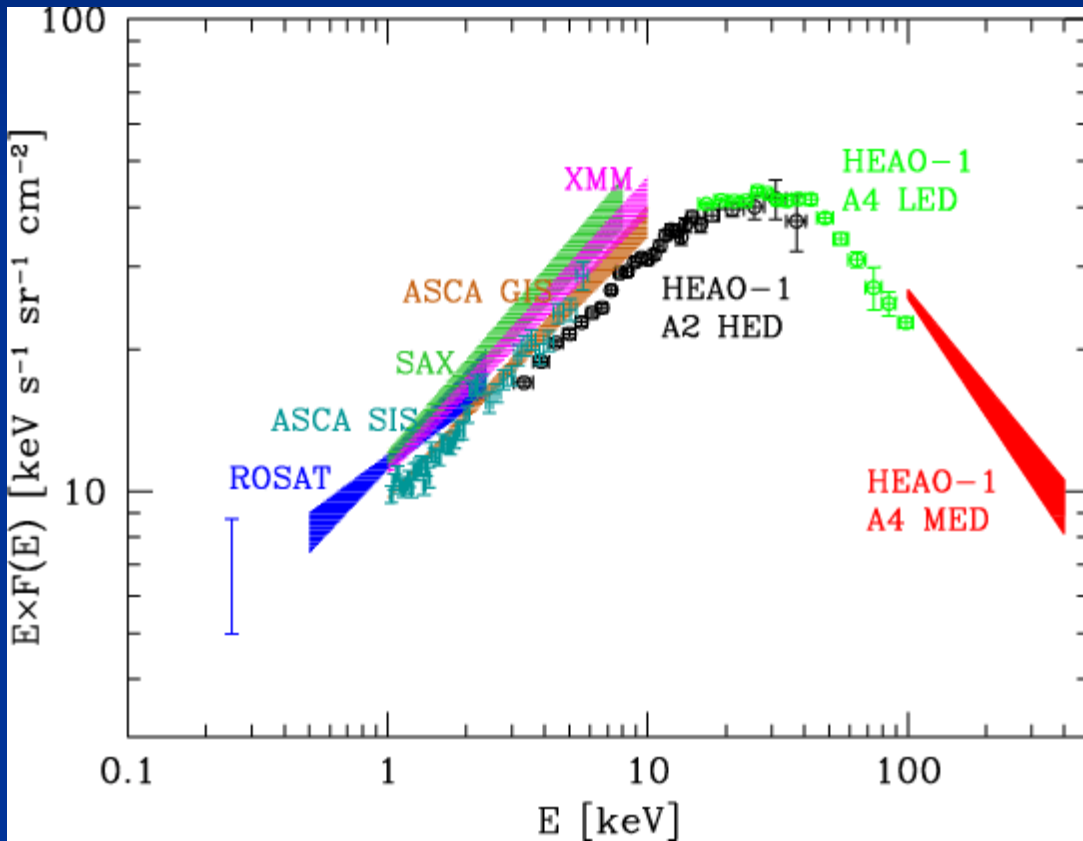
# AGN ACROSS THE HUBBLE TIME

❖ Composite spectra of high-redshift quasars



# AGN ACROSS THE HUBBLE TIME

## ❖ AGN contribution to X-Ray background



First spectral data (1980s) in the 3-60 keV band could be reproduced accurately by thermal emission from an optically thin plasma:

$$F(E) \approx E^{-0.29} e^{-E/41\text{keV}} \quad (\text{bremsstrahlung})$$

So, can a diffuse plasma emission explain the XRB?

**NO!**

Subtracting AGN  $\rightarrow$  XRB spectrum no more compatible with bremsstrahlung emission!

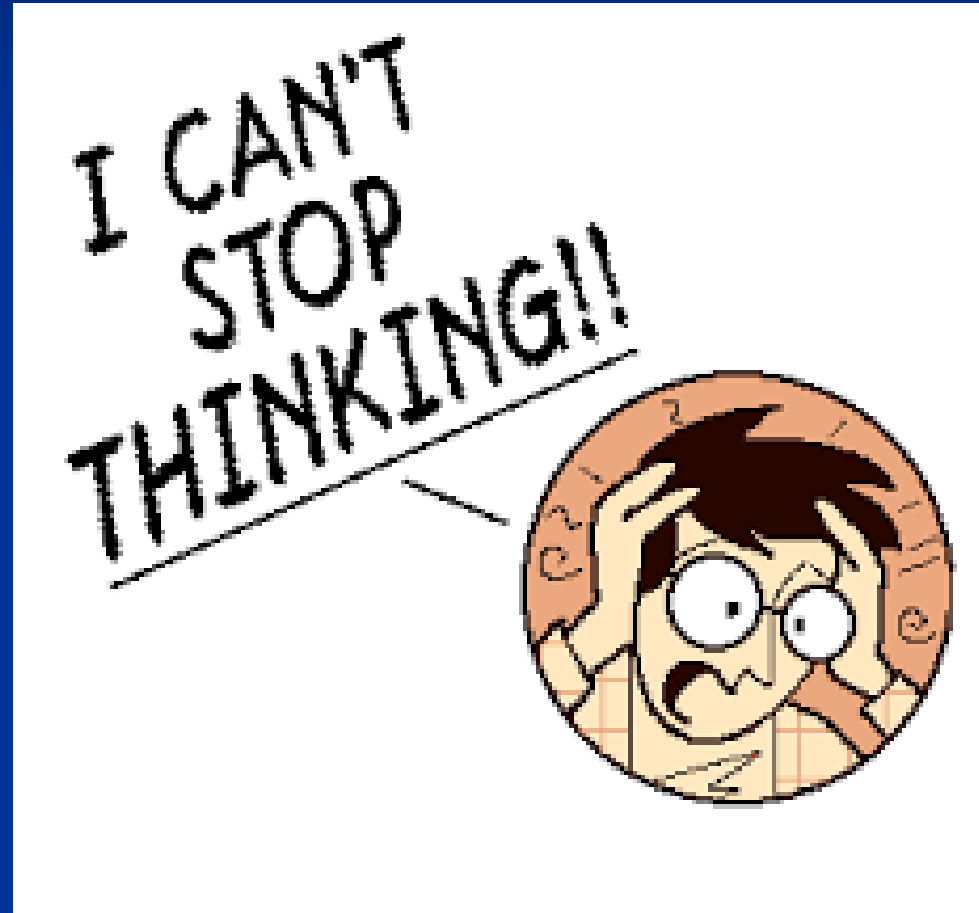
CMB represents a perfect blackbody:

hot gas  $T \sim 40 \text{ keV} \approx 4 \times 10^8 \text{ K}$

would produce distortions by inverse Compton effect

**Emission by unresolved, faint individual sources  $\rightarrow$  AGN**

# WHAT ABOUT AGN AT HIGH REDSHIFTS?



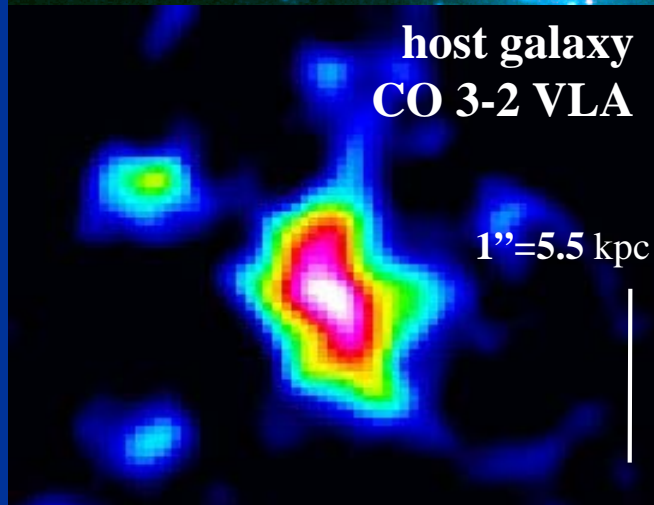
# HIGH-REDSHIFT AGN

## ❖ High-z QSOs

SDSS J1148+5251



$z=6.42 \rightarrow$   
age  $\sim 0.87$  Gyr  
 $L_{\text{bol}} \sim 10^{14} L_{\odot}$



host galaxy  
CO 3-2 VLA

$1'' = 5.5$  kpc

$M_{\bullet} \sim 3 \cdot 10^9 M_{\odot}$  (Willott et al. 03)

$M_{\bullet}$  derived from  $L_E$ , Mg II width, other

Kurk + 2006-8, Fan + 2003-8

Barth et al.

Bertoldi et al. (2009/10)

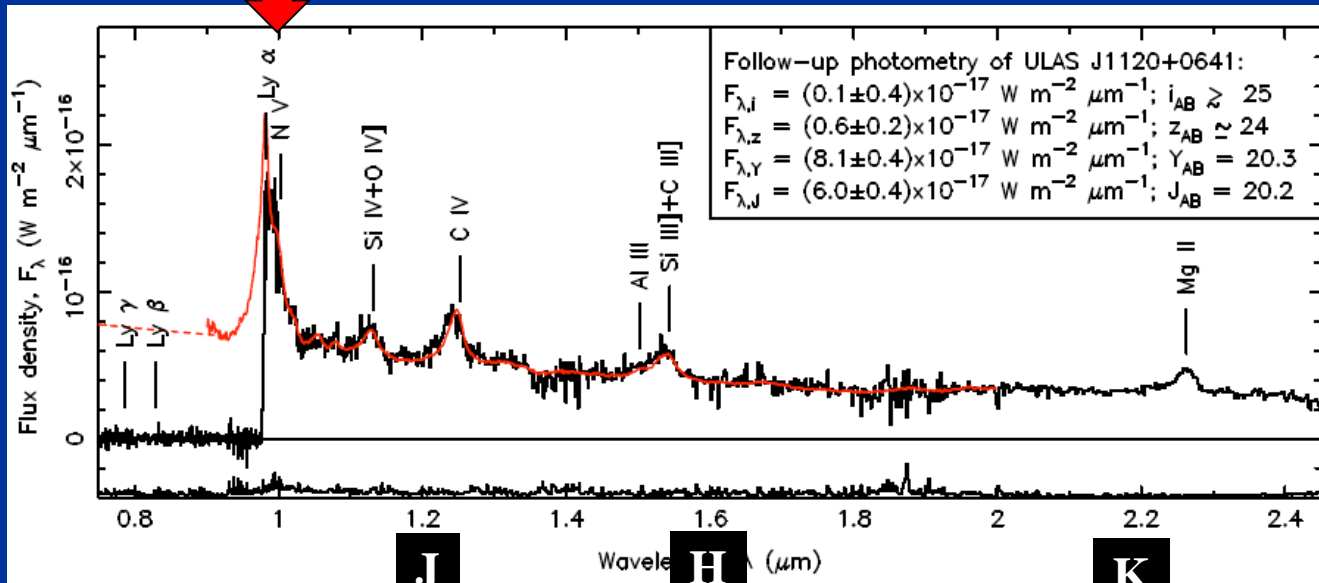
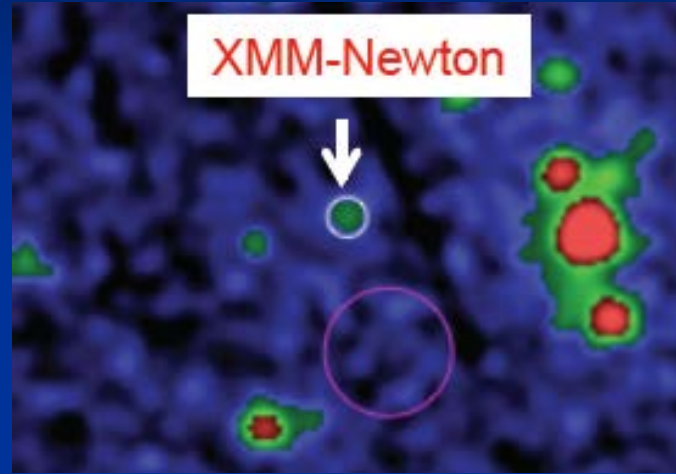
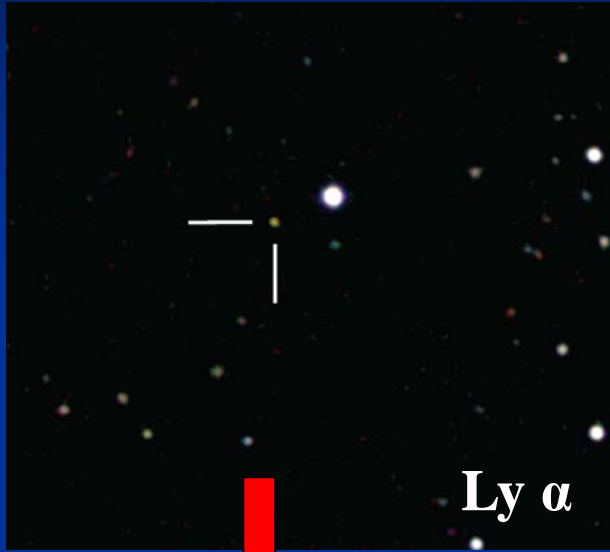
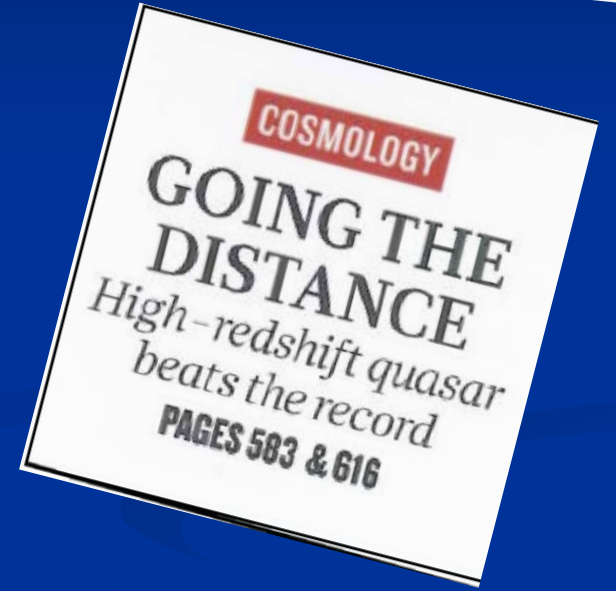
Fan et al.

# HIGH-Z QSOS

❖ Mortlock + (2011):  $z \sim 7.085$

ULAS J1120+0641

$M_{\bullet} \sim 2 \times 10^9 M_{\odot}$

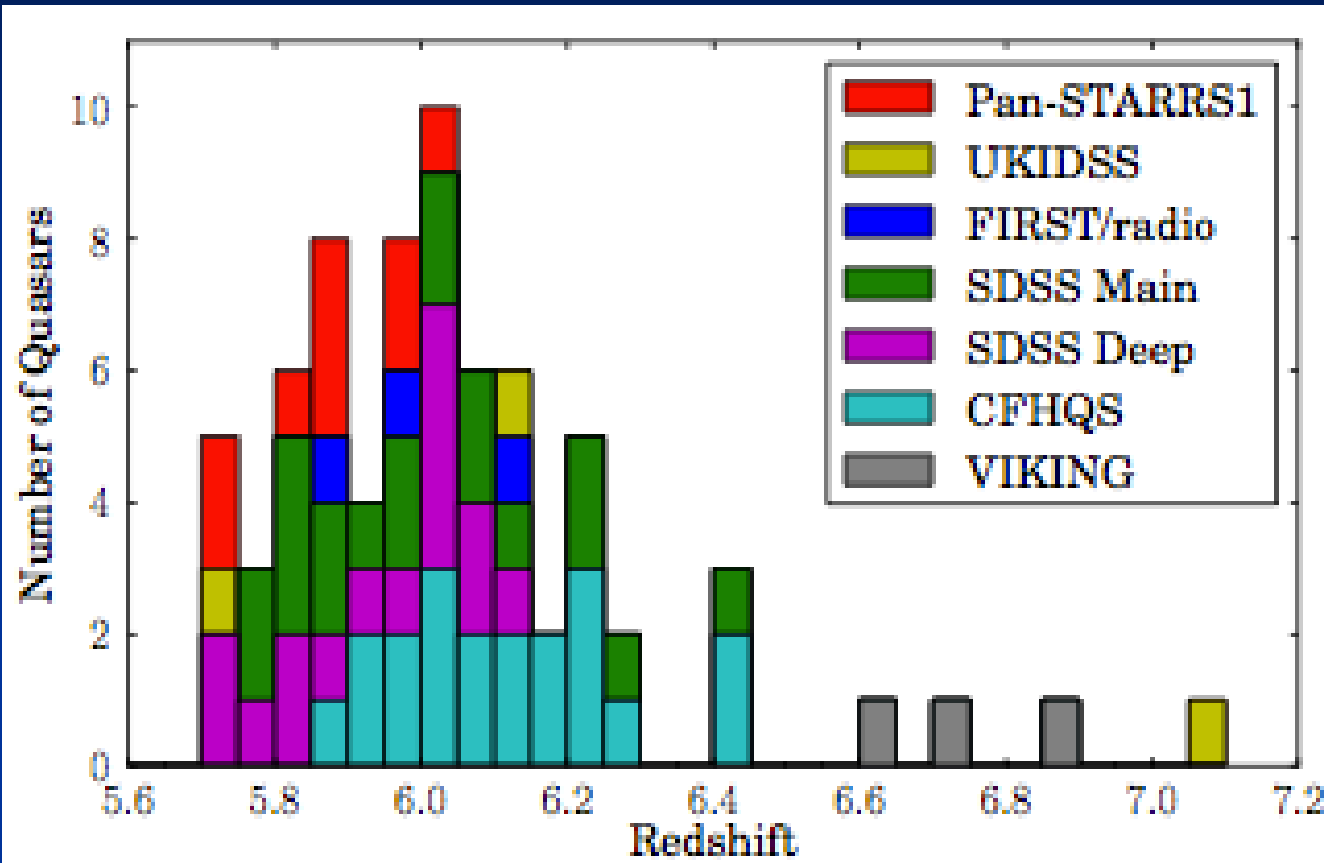


$L_{2-10\text{keV}} \sim 6.7 \times 10^{44} \text{ erg/s}$   
 $\rightarrow$  QSO accretes at the Eddington rate

Moretti et al. 2014



# WHERE DO WE STAND?



About 80 QSOs at  $z > 5.7$

→ SDSS, CFHQS, Pan-STARRS1

(Fan et al. 2000–06; Jiang et al. 2008,09;  
Willott et al. 2007,09,10; Banados et al.14)

including UKIDSS/VISTA

(Mortlock et al. 2011; Venemans et al. 2013)

Less than 25% with X-ray detections

SDSS traces the most luminous QSOs  
 $\log L_x \sim 45$ ,  $\log L_{bol} \sim 46.5$ ,  $M_{1450} = [-24, -28]$

**Faint end of the LF still to be achieved!**

# WHERE DO WE STAND?

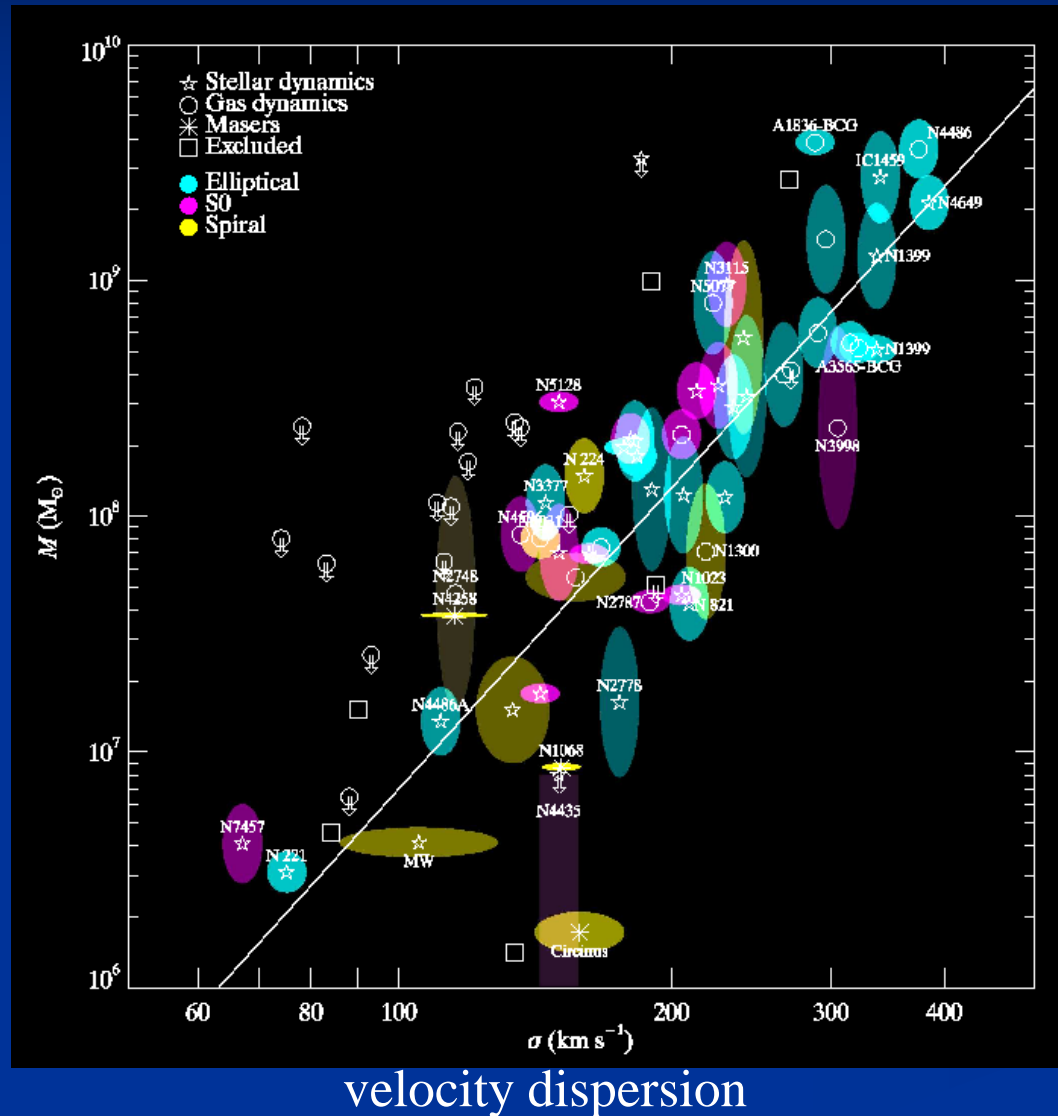
❖ Number of high- $z$  AGN detected so far:

|         | SDSS | X-ray sel. |
|---------|------|------------|
| $z > 3$ | 8000 | 100        |
| $z > 4$ | 1500 | 15         |
| $z > 5$ | 150  | 3-4        |
| $z > 6$ | 10   | 0          |

From: Silverman et al. 2008; Hasinger 2008; Brusa et al. 2009;  
Civano et al. 2010; Fiore et al. 2012; Vito et al. 2013

# DO SMBHS KNOW ABOUT HOST GALAXIES?

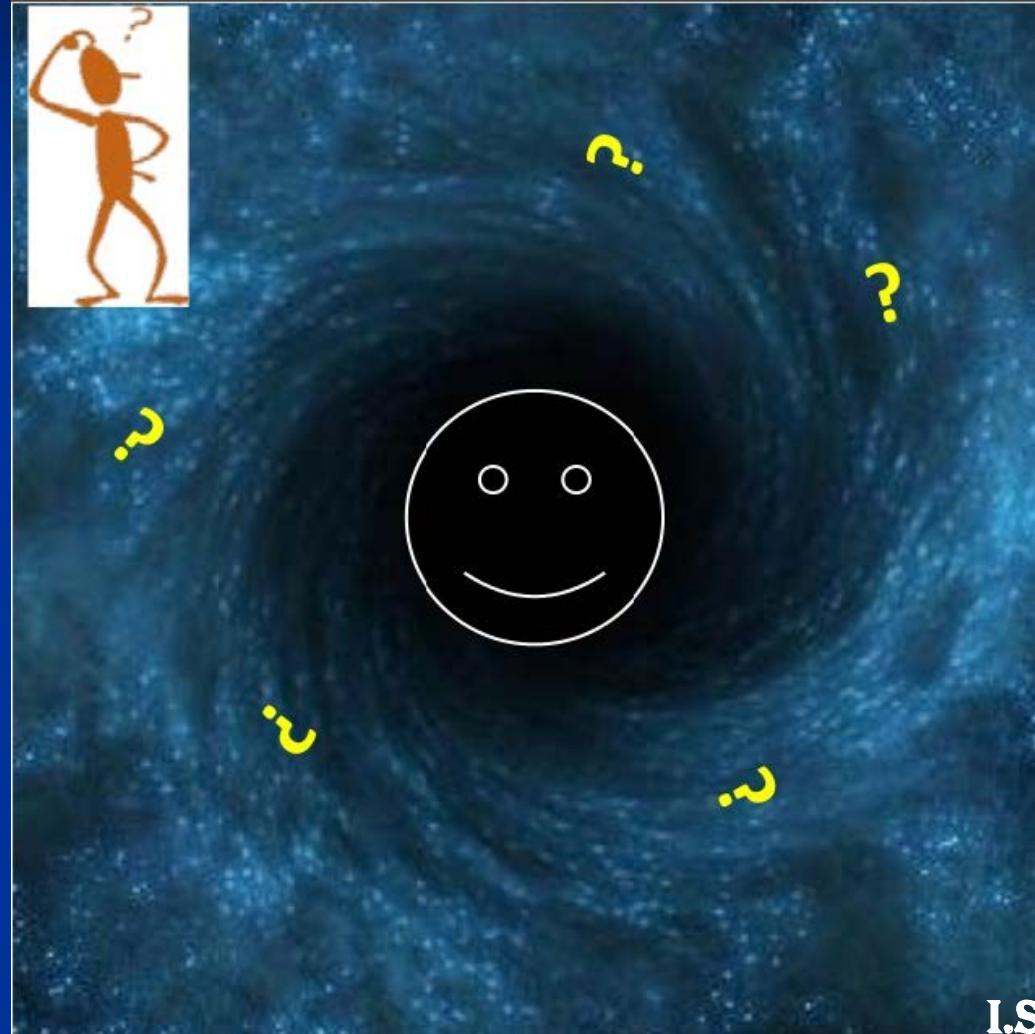
## ❖ SMBHs and host galaxies: scaling relations



Scaling relations between SMBH masses and host galaxy properties (stellar bulge mass, luminosity, velocity dispersion)

Are AGN and galaxies closely tied in their evolution?

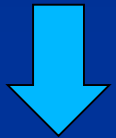
# PATHWAYS TO BLACK HOLE SEEDS



# PATHWAYS TO BLACK HOLE SEEDS

❖ How do SMBHs form?

AGN are the most luminous long-lived objects in the universe  
 SMBHs power AGN



What are the seeds of SMBHs in the early universe or anywhere?

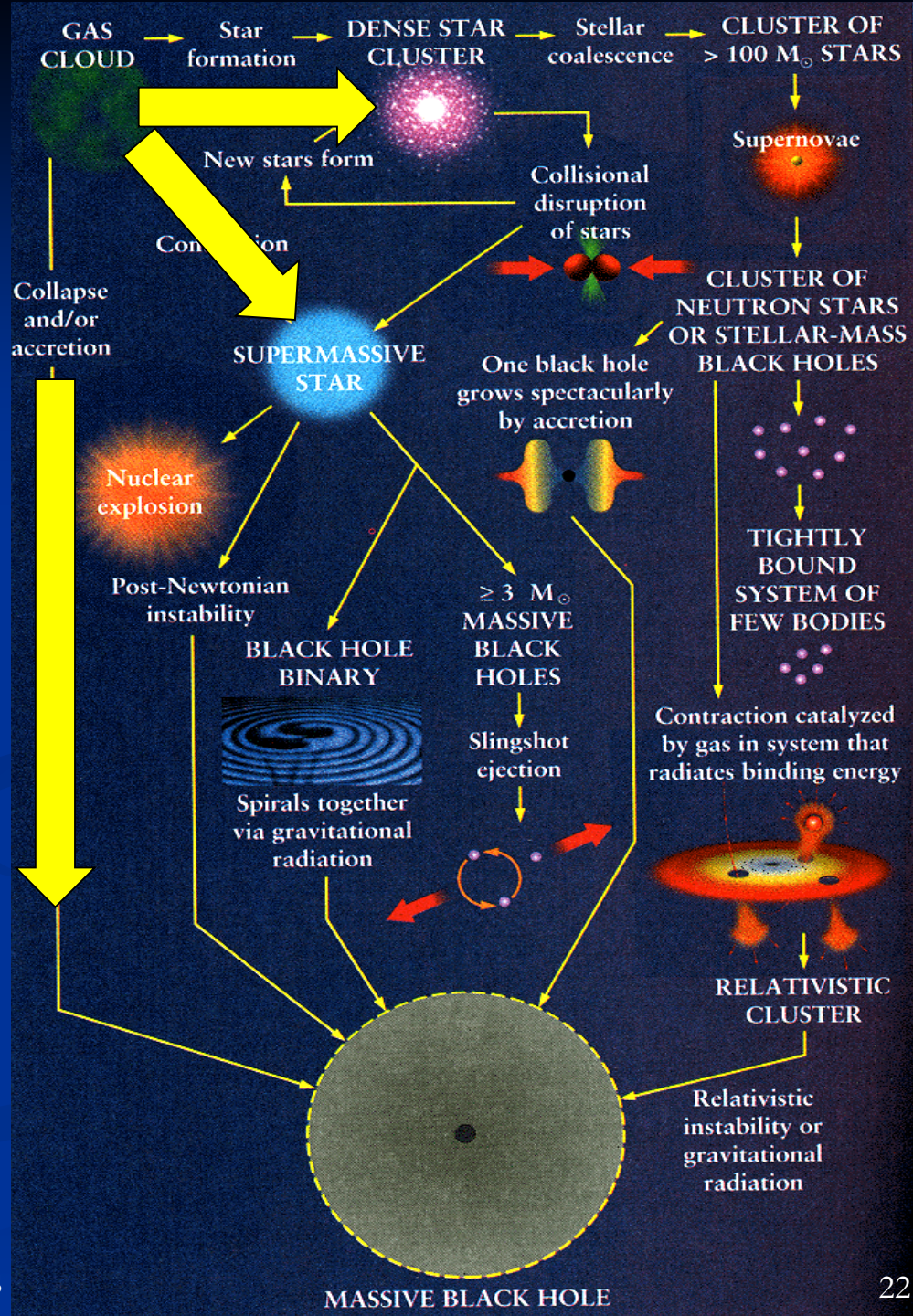
When did they form? What type of seeds: light, intermediate or heavy?

Are SMBH seeds related to the bulge growth?

Impact of first quasars on ISM/IGM structure formation?

SMBH-galaxy co-evolution is unconstrained at  $z > 6$

Heavily obscured accretion mostly unconstrained beyond the local universe



Rees (1978)  
 Begelman & Rees

# PATHWAYS TO BLACK HOLE SEEDS



primordial BH seeds



relativistic star clusters



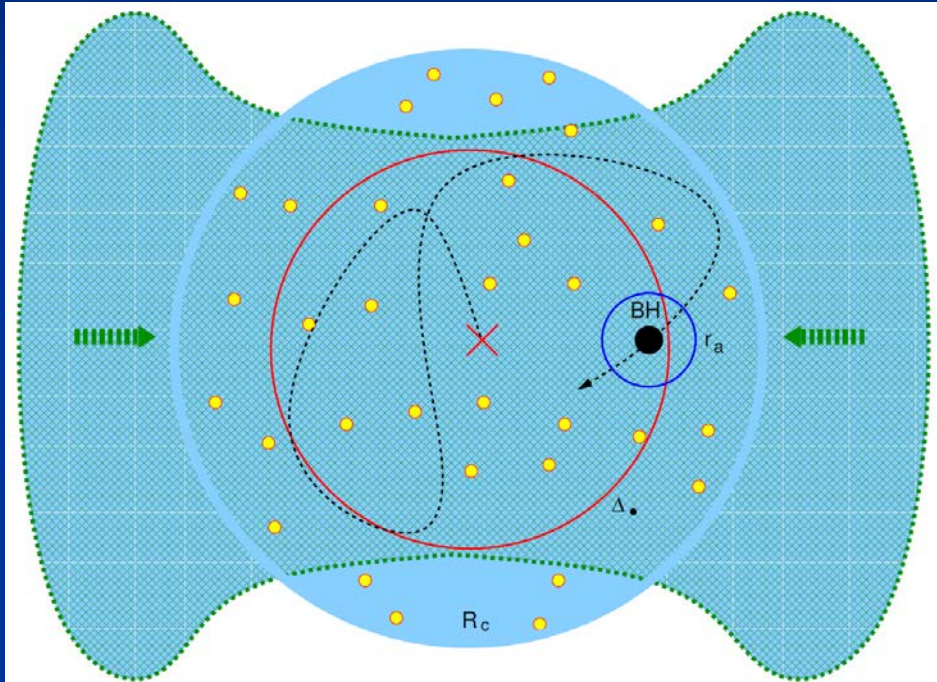
Pop III stars



direct collapse

# PATHWAYS TO BLACK HOLE SEEDS

## ❖ SMBHs from stellar clusters



Alexander & Natarajan (2014)

Supra-exponential growth of a small BH to SMBH... in stellar/gaseous cluster...

where do you find such clusters at high  $z$ ?

# PATHWAYS TO BLACK HOLE SEEDS

## ❖ SMBHs from relativistic stellar clusters



Gravitational instability: when stellar cluster is compact enough

- gravitational redshift from its center  $> 0.5$
- runaway collapse

How can such a cluster form at high  $z$ ?



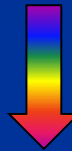
# PATHWAYS TO BLACK HOLE SEEDS

## ❖ SMBHs from Population III remnants

Population III stars forming in minihalos

$$M_* \sim 10^{2-3} M_\odot \quad ?$$

(Bromm & Loeb 2003...)



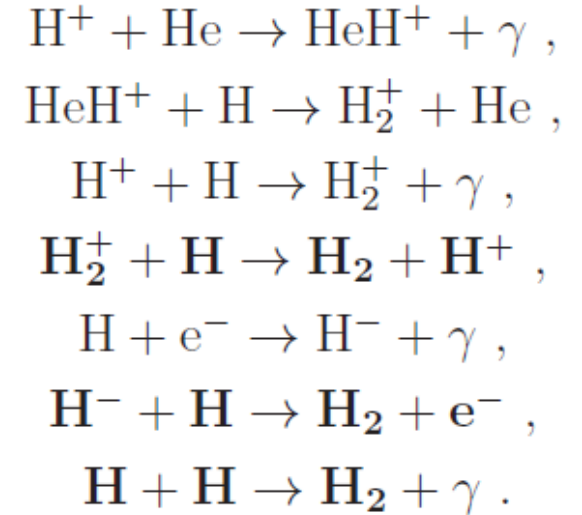
**radiation feedback!**

$$M_* \sim 40 M_\odot$$

Hosokawa+11,14; Wise + 12

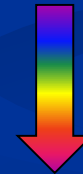


important role  
of H<sub>2</sub>



BH seeds:

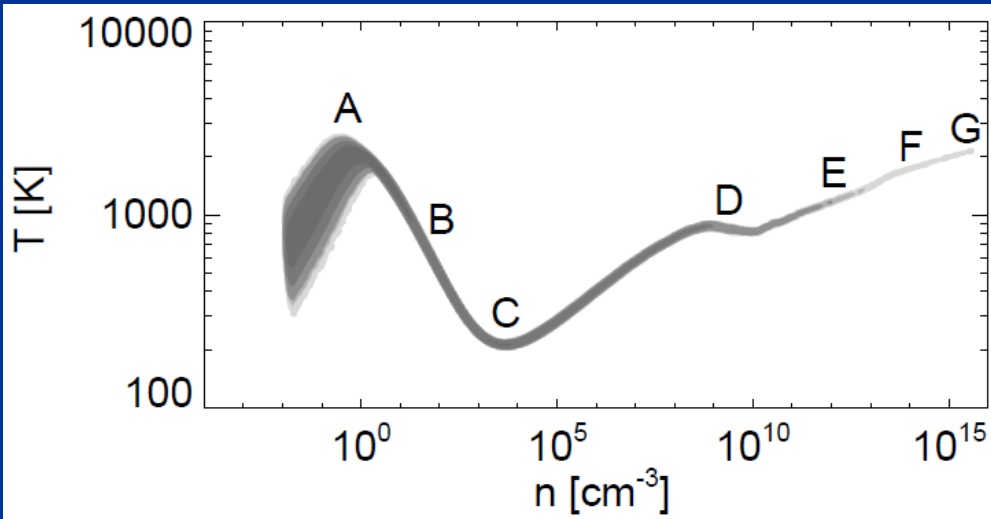
$$M_* \sim 10^2 M_\odot$$



$$M_* \sim 10 M_\odot$$



**BH seeds at z > 20**



Omukai (2001)  
Yoshida (2006)

- A: virial shock in mini-halos
- B: gas/H<sub>2</sub> cool down to 200 K
- C: LTE
- D: 3-body reactions → all H<sub>2</sub>
- E: H<sub>2</sub> cooling inefficient
- F: collisional cooling → radiation
- G: H<sub>2</sub> dissociates at 2,000 K

# PITFALLS WITH POP III BLACK HOLE SEEDS

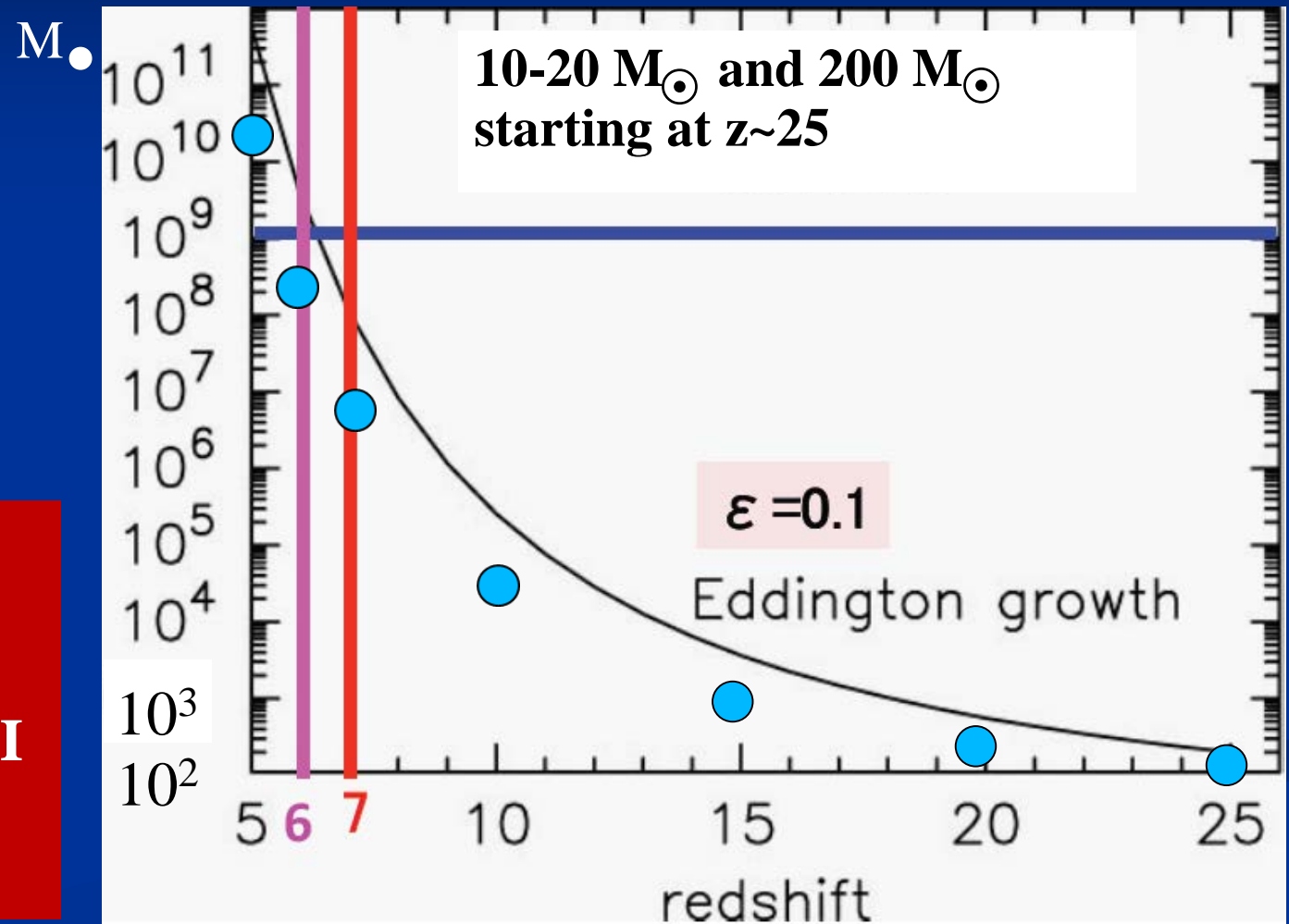
If start with  $\sim 10\text{-}20 M_{\odot} \rightarrow$  we are  $\sim 1.5\text{-}2$  decades off for the SMBH masses at  $z \sim 7$

Needs about

$$\ln(2 \cdot 10^9 M_{\odot} / 10 M_{\odot})$$

$\sim 19$  e-foldings

**too frequent mergers:** BH slingshot and ejection from minihalos  
**BH accretion growth:** limited by PopIII + feedback + accretion cutoff



Maybe Pop III remnants grow to smaller SMBH, e.g.,  $10^{6-7} M_{\odot}$ ?

# PATHWAYS TO BLACK HOLE SEEDS

## ❖ SMBHs from Direct Collapse in dark matter halos

Direct ('monolithic') collapse

If **primordial** gas:

cooling floor  $T_{\text{gas}} \sim 10^4 \text{ K}$

collapse into DM halos of

$$T_{\text{vir}} \geq T_{\text{gas}} \sim 10^4 \text{ K}$$

(Haiman & Loeb 2001;  
Begelman, Volonteri & Rees 2006)

$$M_{\text{vir}} \sim 10^8 M_{\odot}$$

**Planck** CDM Universe: 16% baryons

$$M_{\bullet} \sim 10^{5-7} M_{\odot}$$



If gas **enriched** by Pop III

→ BH seeds

$$M_{\bullet} \sim 10^{4-5} M_{\odot} \quad ?$$

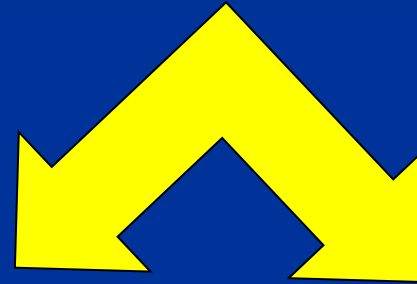
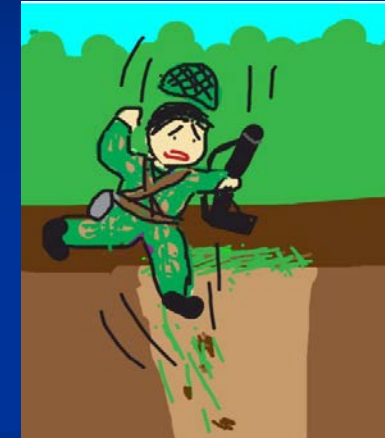
(because gas collapses earlier into smaller halos)

**BH seeds at  $z \sim 10-15$**

# DIRECT COLLAPSE TO SMBH SEEDS: PITFALLS

Growth by accretion/mergers from  $z \sim 20$  to  $z \sim 7 \rightarrow$  to  $10^9 M_{\odot}$

less mergers! BUT...



angular momentum barrier  
 $\rightarrow$  fragmentation  
star formation  
because  $\lambda > 0$

gas collapse  $\rightarrow$  fragmentation  
star formation  
because  $M_{\text{gas}}/M_{\text{J}} > 1$

many stars, small black holes(s)?

WHAT IS THE NATURE OF THE CENTRAL OBJECT?

# SPECIFICS OF GRAVITATIONAL COLLAPSE

Generic accretion rate (with and w/o central BH):

$$\dot{M} \sim \frac{v^3}{G} \sim 2 \times 10^{-4} T_{100}^{3/2} M_{\odot} \text{ yr}^{-1}$$

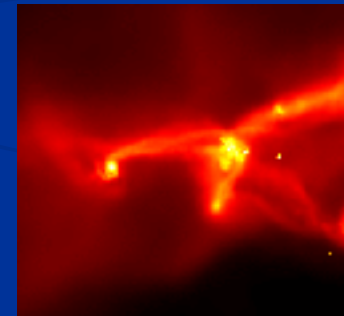
Stellar collapse:  $T \sim 10 \text{ K} - 100 \text{ K}$

$$\dot{M} \sim 10^{-5} - 10^{-4} M_{\odot} \text{ yr}^{-1}$$



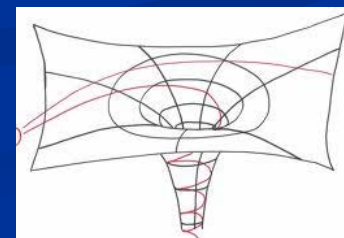
Pop III stellar collapse:  $T \sim 10^2 \text{ K} - 10^3 \text{ K}$

$$\dot{M} \sim 10^{-4} - 10^{-3} M_{\odot} \text{ yr}^{-1}$$



Direct collapse:  $T \geq 10^4 \text{ K}$

$$\dot{M} \geq 0.1 M_{\odot} \text{ yr}^{-1}$$

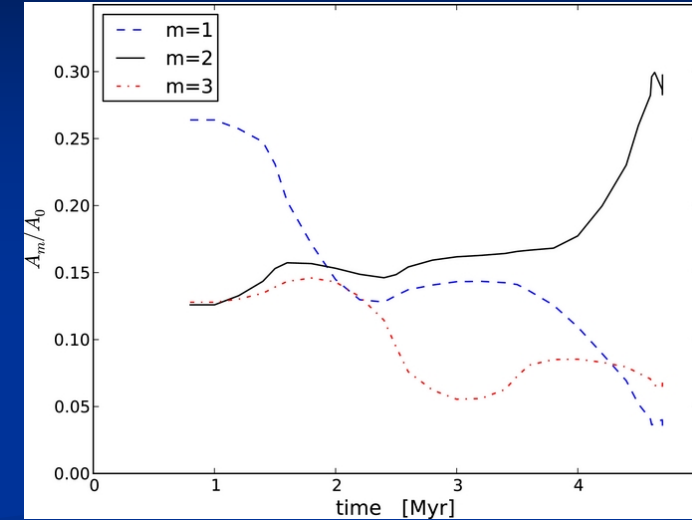


# HOW TO COLLAPSE

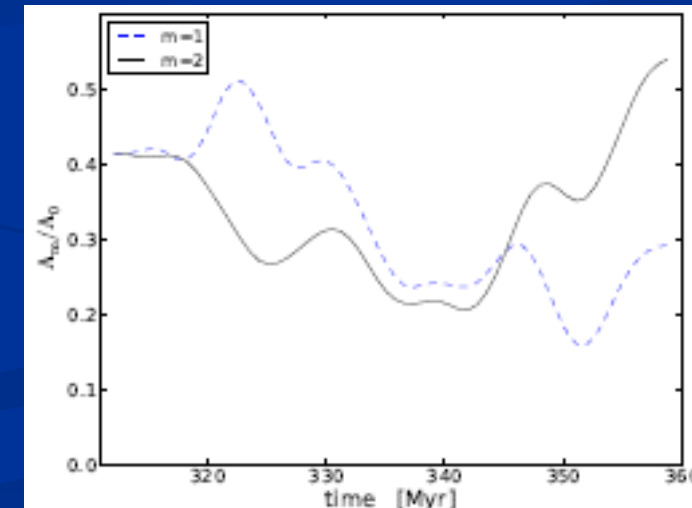
❖ Spherical collapse: forget about this → angular momentum barrier



isolated halo



cosmological halo



even if the original configuration has only  
~ 1% Keplerian support!

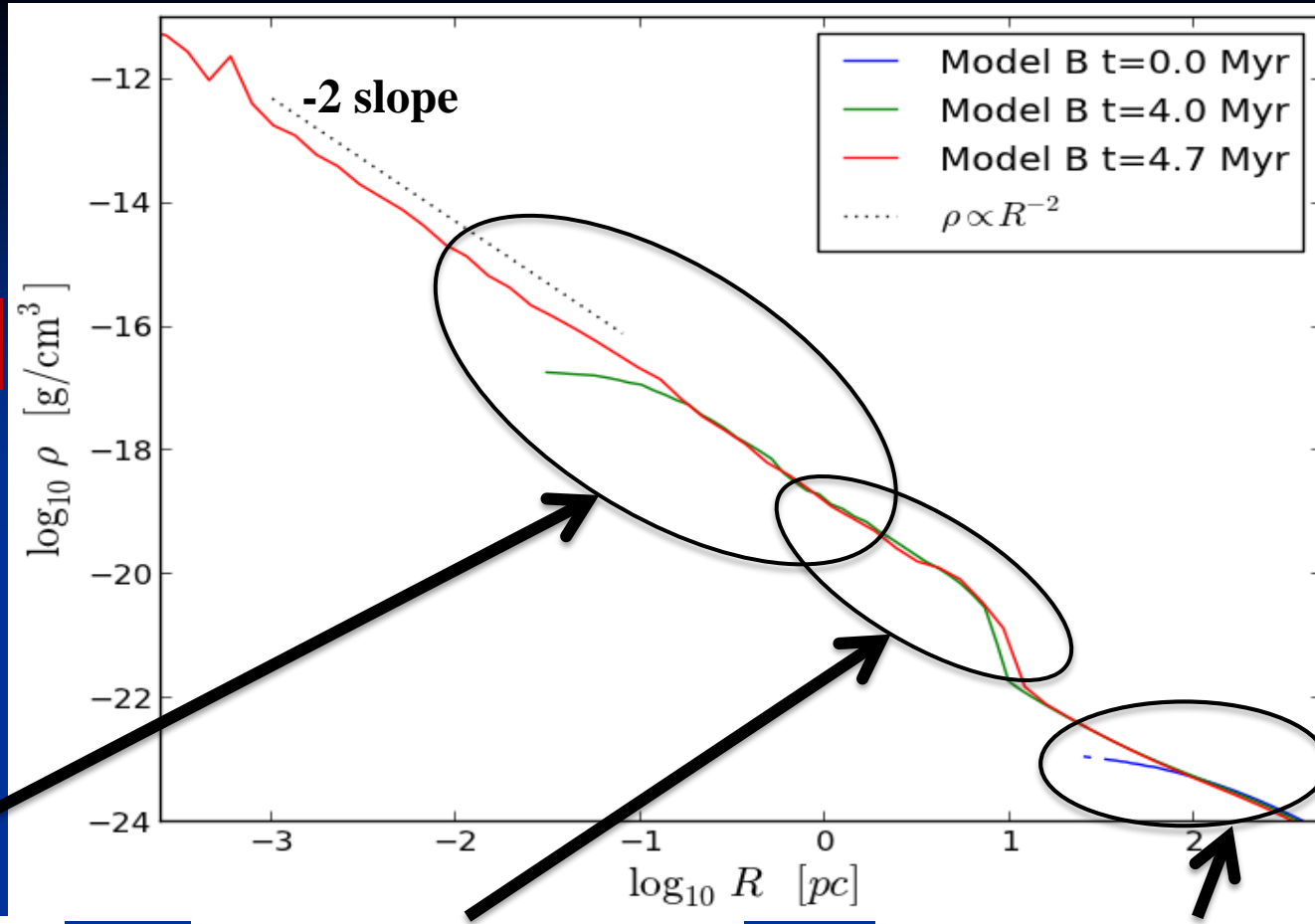
With nonzero spin → angular momentum is not conserved!  
→ gravitational torques

Angular momentum transfer is dominated by the low Fourier modes

Choi et al. 2013, 2015

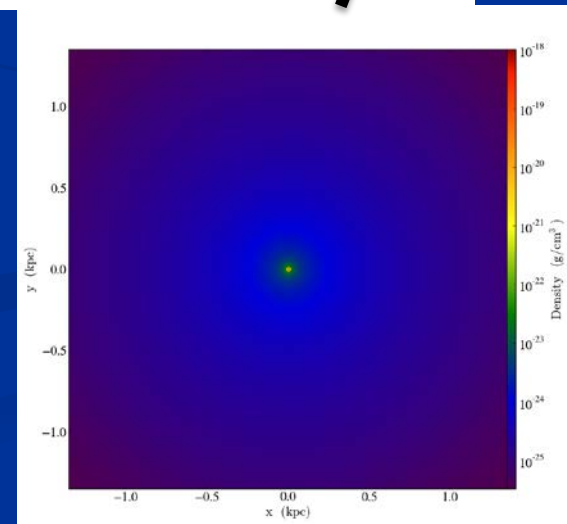
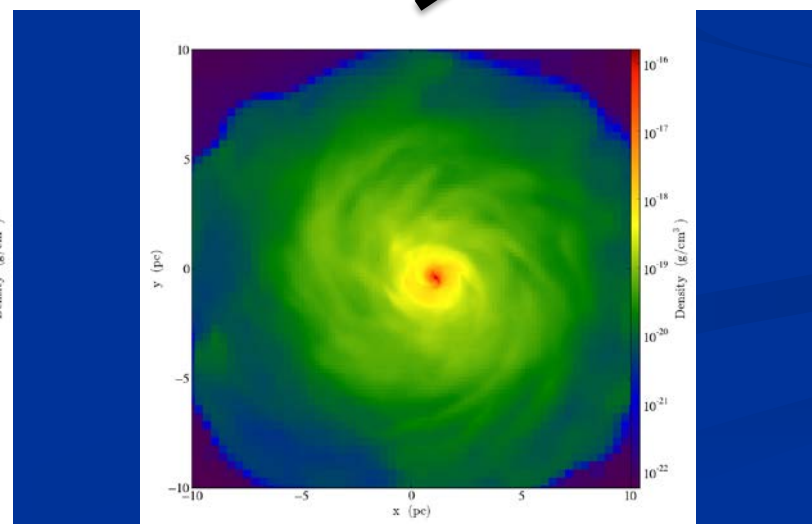
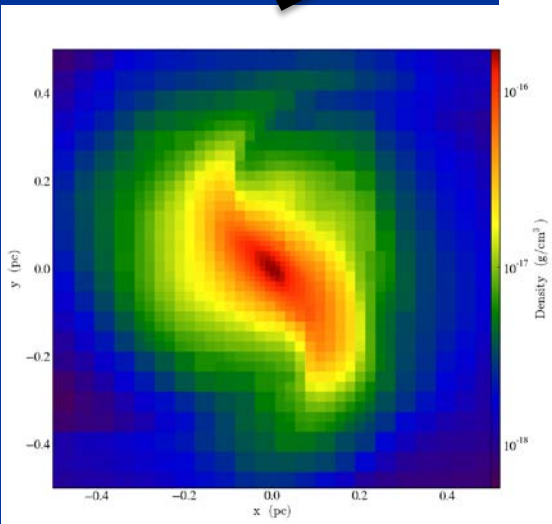
# Overcoming the J-barrier

gas density profile



Self-similar collapse

how to break this symmetry?

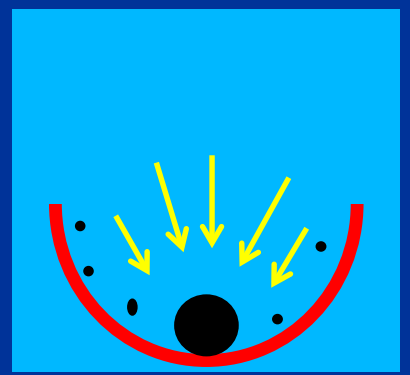


# COSMOLOGICAL SMBH SEEDS: BREAK SYMMETRY OF GRAVITATIONAL COLLAPSE

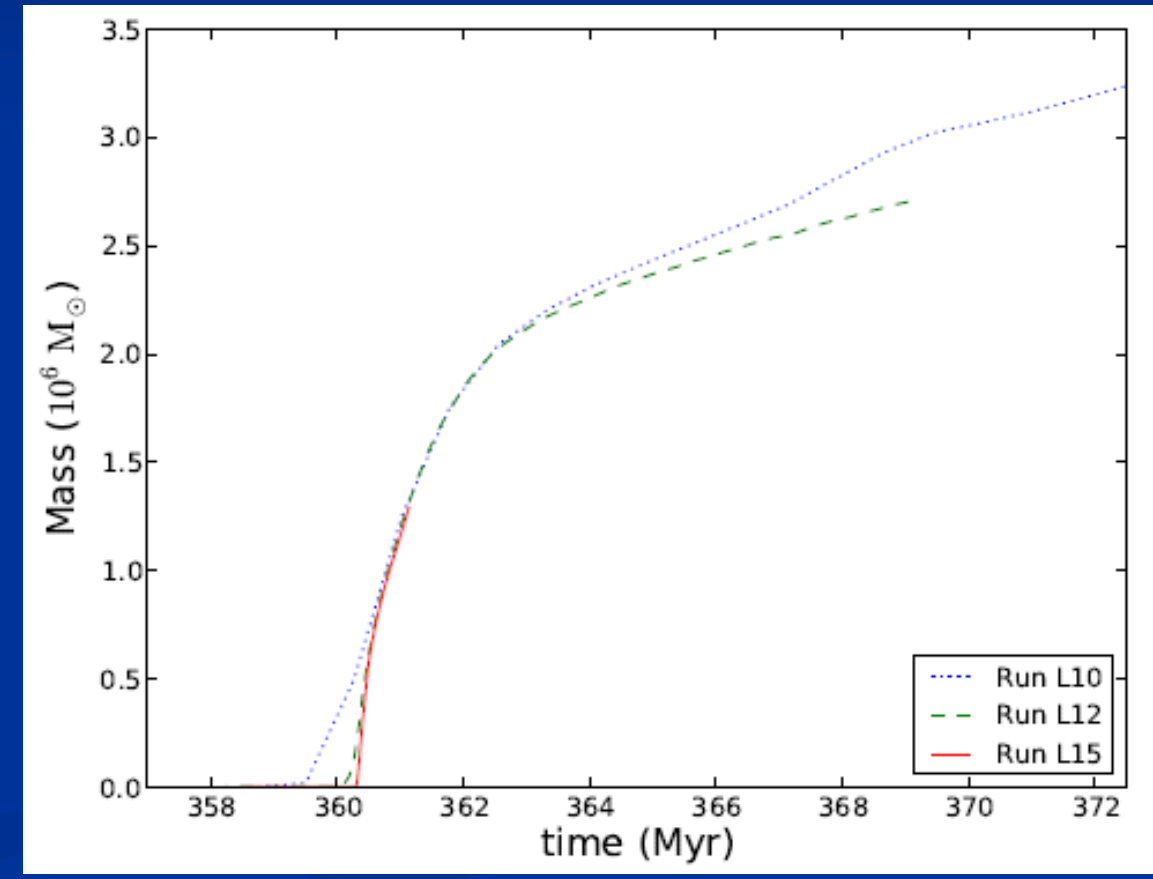
Using **sink particles** to follow up the collapse

Initial masses  $\sim 10 M_{\odot}$

In all models: one sink particle grows by merging and accreting gas and other sink particles

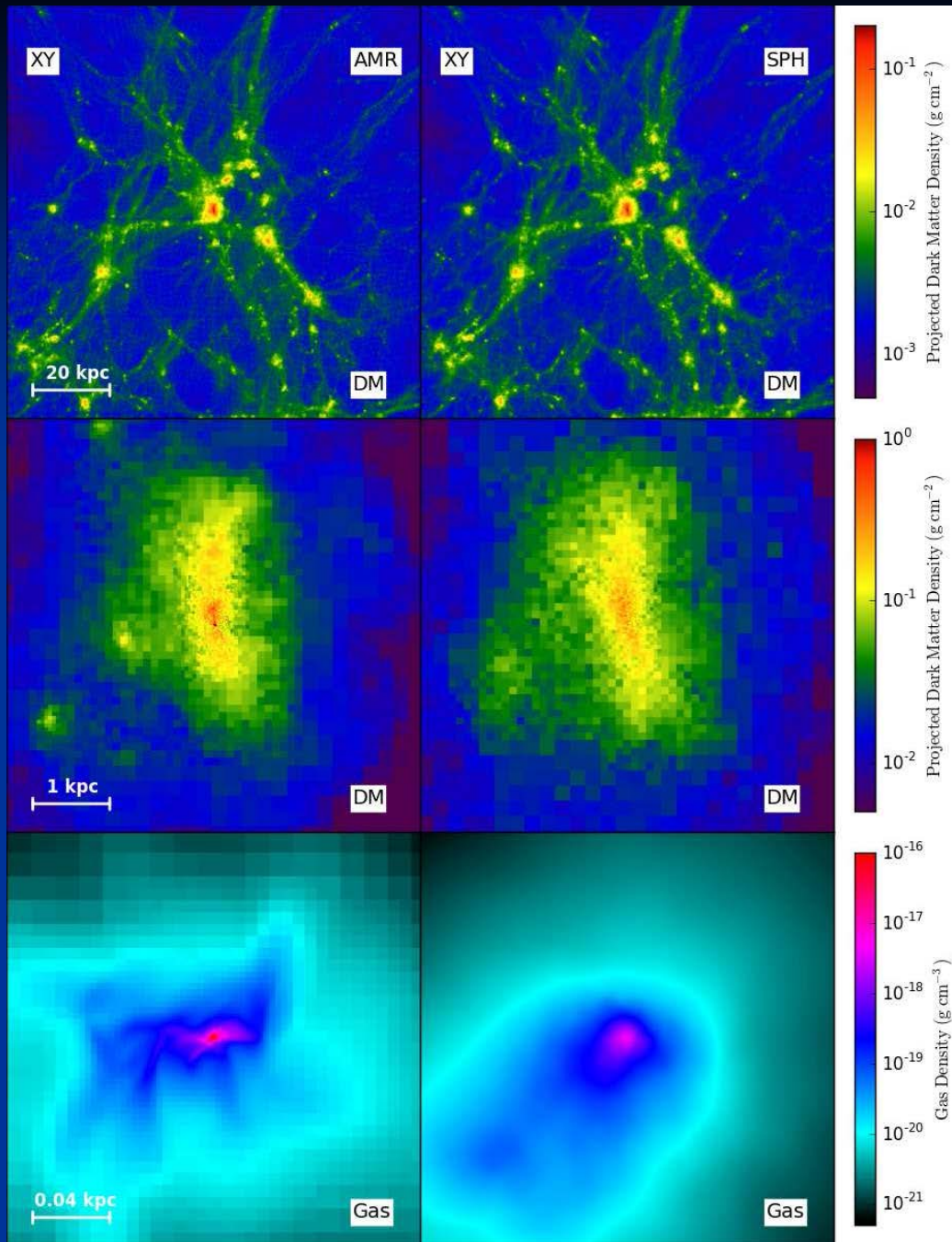


evolution of the central seed



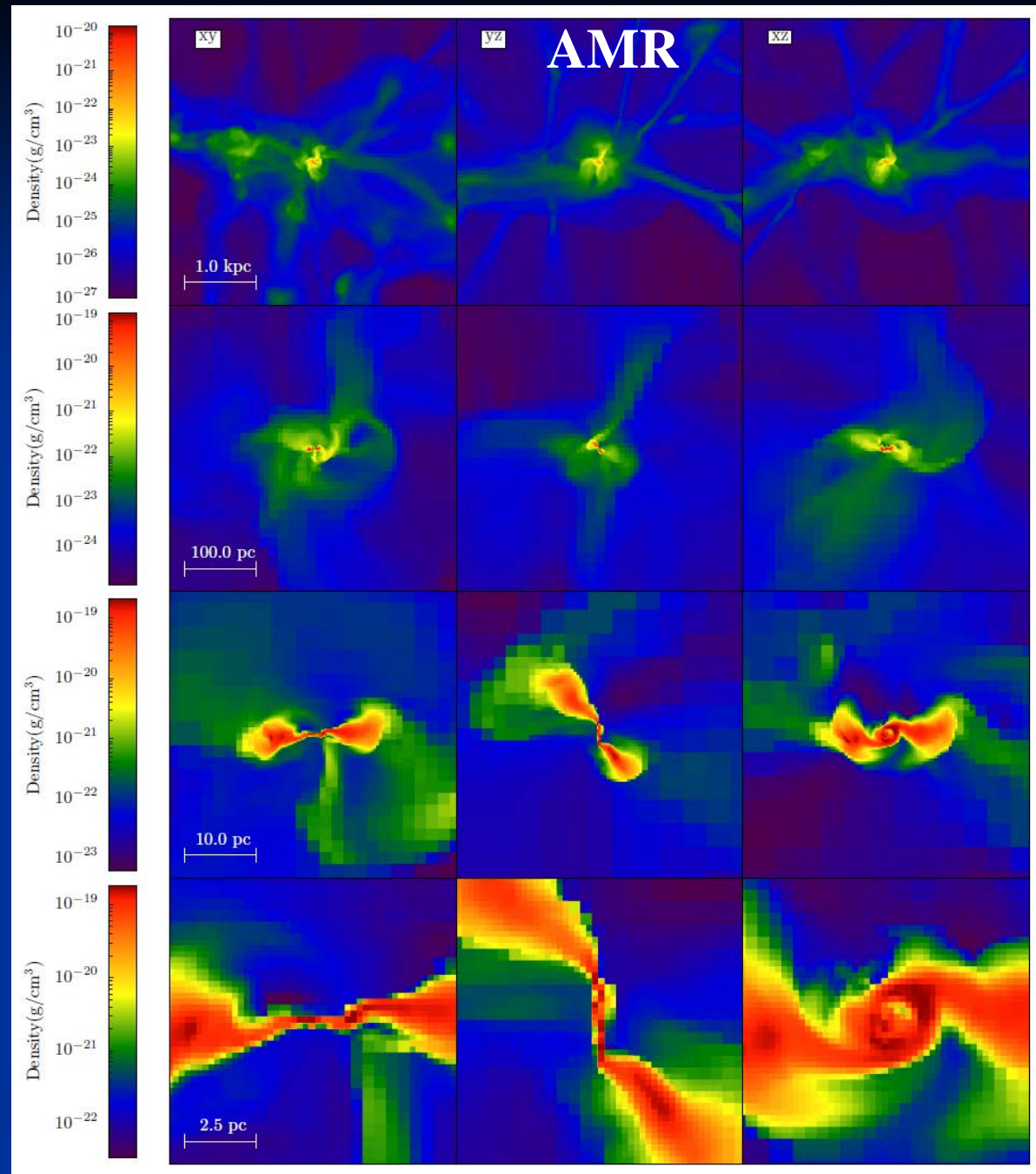


No seeds



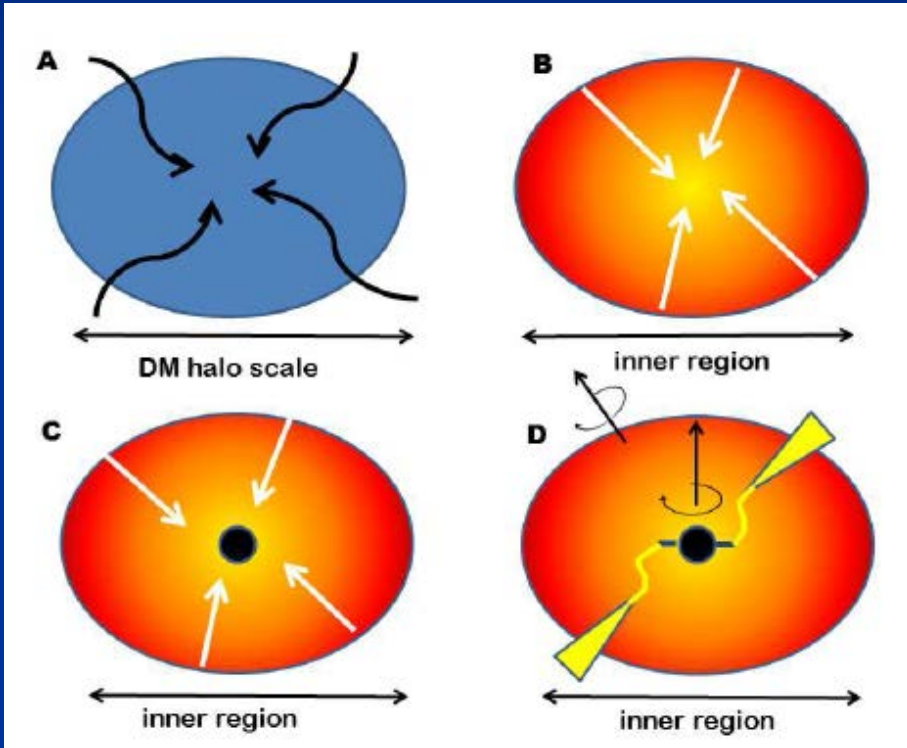
Yang, Nagamine & Shlosman 2015

With seeds

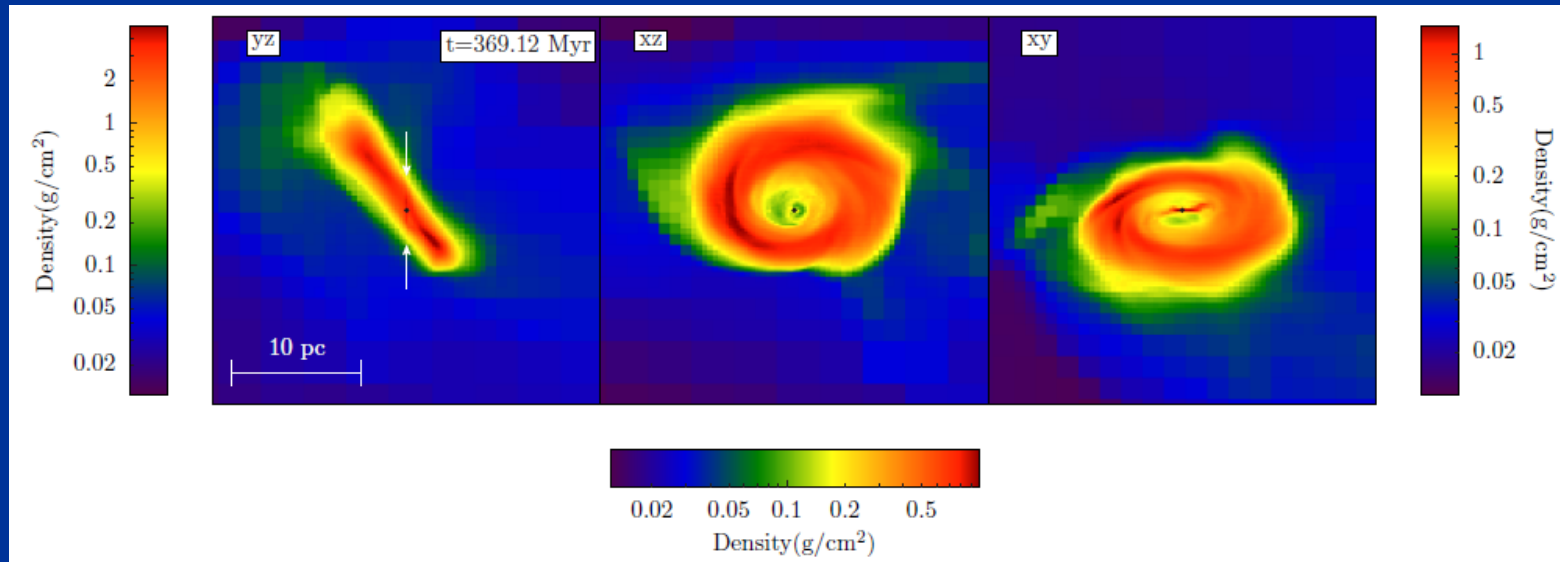


IS, Choi, Begelman & Nagamine 2015

# COSMOLOGICAL SMBH SEEDS: BREAK SYMMETRY OF GRAVITATIONAL COLLAPSE



misaligned disks



Breaking collapse symmetry  $\rightarrow$  disk formation

Variability of angular momentum in accreting gas  $\rightarrow$  misaligned disks

# DEEP INTERIOR FLOW: OPTICALLY-THICK COLLAPSE

Under what conditions will accretion flow continue to smaller  $R$ ?

What type of flow(s) for  $R < R_{\text{trap}}$  ?

Quite a bizarre flow!



Waterfall by: M.C. Escher

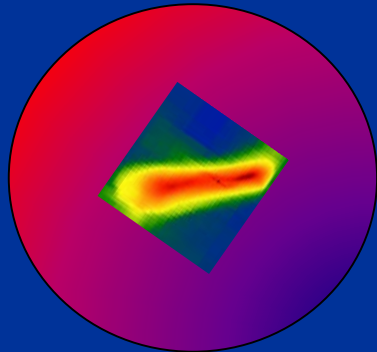
# DEEP INTERIOR FLOW: OPTICALLY-THICK COLLAPSE

Two options here (Begelman & Shlosman 2009)

$$j(R) \sim \beta j_K$$

The flow can stay being rotationally supported (+  $\nabla P_{\text{rad}}$ ) (Choi et al. 2013,2015)

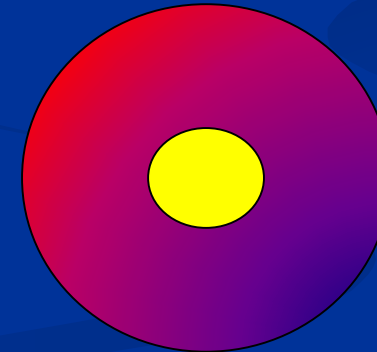
where  $\beta < 1$



Disklike flow  $\rightarrow$  ADAF/CDAF/ADIOS flow?

The flow can lose  $j$  in a few rotations by some unspecified way (Begelman 2012)

where  $\beta \ll 1$

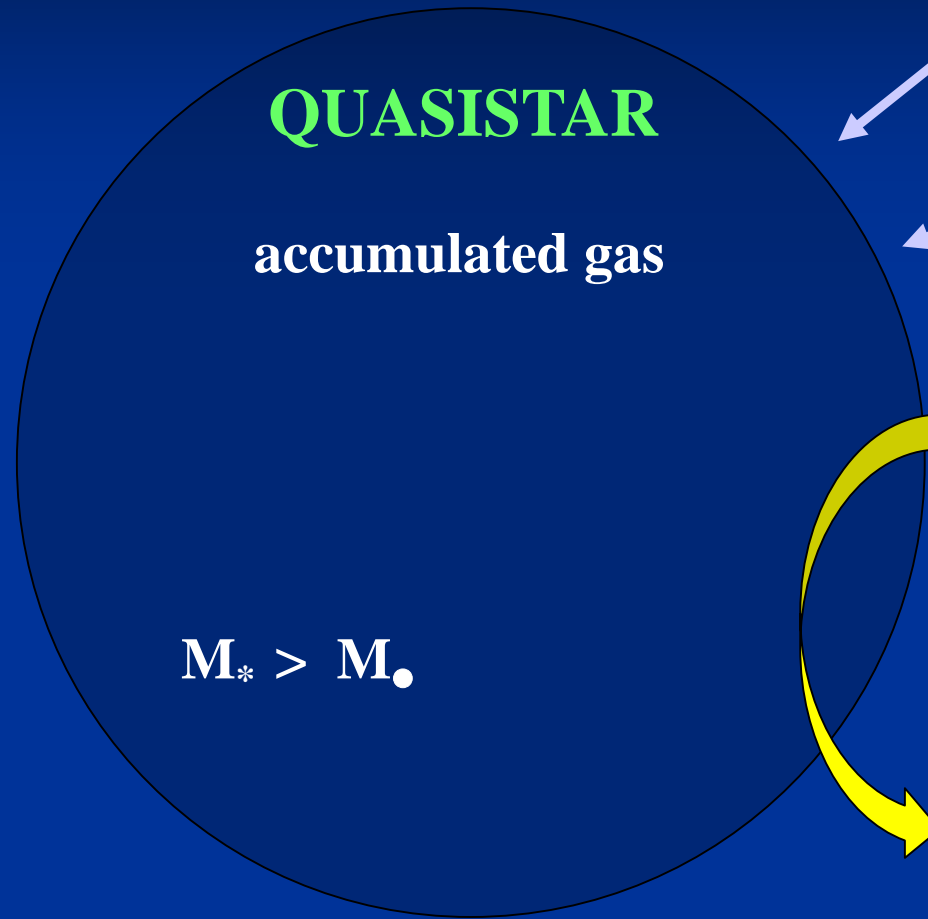


Starlike accretion flow

What does this difference mean?

# A QUASISTAR FORMATION? (Begelman, Volonteri & Rees 2006)


Rapid growth of the BH seed in DM minihalos with  $\sigma \sim 10\sigma_{10} \text{ km s}^{-1}$



$$\dot{M}_* \sim \frac{\alpha \sigma^3}{G}$$

quasistar radius: where accretion becomes super-Eddington for a given opacity  $\rightarrow$  radiation trapping:

thermonuclear reactions:  $t \sim 10^6$  yrs

$$R_* \sim 10^{13} \alpha \sigma_{10}^3 \left( \frac{\kappa}{\kappa_{es}} \right) \text{ cm}$$


**Seed BH  $\rightarrow$  from neutrini cooled core of  $\sim 10 M_{\odot}$   $\rightarrow$  grows at the rate of**

$$\dot{M} \gg \dot{M}_E$$

# TWO ALTERNATIVES IN DIRECT COLLAPSE

Starlike accretion flow: **J goes in smoke**

angular momentum is **un**important  
→ no preferential channel of energy release (Begelman et al. 2006-14)

photon trapping  
→ formation of a single accreting massive object in pressure equilibrium  
→ **quasistar**

thermonuclear reactions  
neutrino-cooled core collapse

**supermassive black hole**  
 $\sim 10 M_{\odot} \rightarrow 10^4 M_{\odot}$

Disklike accretion flow: **J important**

angular momentum dominates  
→ defines preferential channel of energy release  
(Begelman & IS 2009; Choi et al. 2013  
IS & Begelman, in prep)

**no** early photon trapping  
→ formation of a single accreting massive object supported by rotation in **R** and pressure in **z**  
→ **quasi 'torus/disk'**

dynamical instabilities  
**no** thermonuclear reactions  
collapse, photons/neutrino partially trapped/escape

**supermassive black hole**  
 $\sim 10^{5-7} M_{\odot}$

# CONCLUSIONS TO TALK 5

- ❖ **AGN** in general, and quasars in particular, have formed early, and probably compete with Pop III stars in illuminating the universe and ending the Dark Ages
- ❖ **There are real problems** to grow BH seeds from Population III remnants to  $\sim 10^{8-9} M_{\odot}$  in less than 1 Gyr, and intermediate-mass BHs have not been found (so far)
- ❖ **Direct (monolithic) collapse** of gas within  $\sim 10^8 M_{\odot}$  DM halos to BH appears to be feasible and (at least) two alternatives exist  
→ collapse via a supermassive star → **quasi-star**, or via **disklike** → **SMBH flow** → **proto-AGN**. Open issues remain with both approaches. The primary role of angular momentum makes the latter approach a realistic one.
- ❖ **Probably the crucial point** is whether it is possible to avoid the thermonuclear phase → it slows down the collapse by  $\sim 10^6$  yr  
→ angular momentum transfer by magnetic and gravitational torques