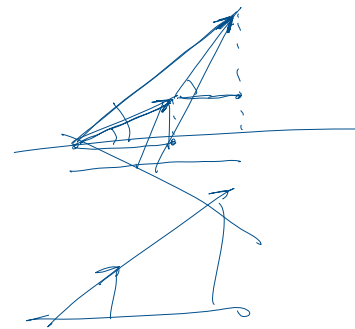
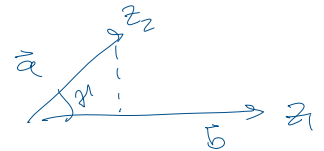
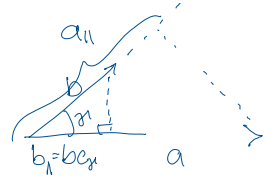


LOS Inner/Outer Products

Wednesday, January 24, 2024 9:54 AM

- The linear structure of LO4 entailed parallel lines, but NOT explicit angle/length
 - Projections needed two complementary spaces: along and into
 - Even \hat{x} , \hat{y} , \hat{z} had no dependence on unit length or perpendicularity (just implied).
 - The dot product provides notions for all of these (unit vectors, orthogonal projections)
- Like vectors themselves, the **dot product** has geometric significance and analytic power
 - Connection between these two aspects is again the linear structure of the dot product
 - Definition: $\vec{a} \cdot \vec{b} = ab \cos \gamma = a_{\parallel} b = ab_{\parallel} = a_x b_x + a_y b_y + a_z b_z = \mathbf{a}^T \mathbf{b} = \delta_{ij} a_i b_j$ $c_{\gamma} \equiv \cos \gamma$
 and $\vec{a} = \hat{a} a = \hat{x} a_x + \hat{y} a_y + \hat{z} a_z = (\hat{x} \hat{y} \hat{z}) \begin{pmatrix} a_x \\ a_y \\ a_z \end{pmatrix} = \hat{\mathbf{e}} \mathbf{a} = \hat{e}_i a_i$; similarly for \vec{b} ; and $\gamma = \angle(ab)$.
 - Length of \vec{a} is $a = \sqrt{\vec{a} \cdot \vec{a}}$ and direction is the unit vector $\hat{a} = \vec{a}/a$, since $\cos 0^\circ = 1 \rightarrow \vec{a} \cdot \vec{a} = a^2$.
 - Geometry: dot product is **product of parallel projected lengths** $\vec{a} \cdot \vec{b} = a_{\parallel} b = ab_{\parallel}$.
 - Either $a_{\parallel} = a \cos \gamma$ (\vec{a} onto \vec{b}) or $b_{\parallel} = b \cos \gamma$ (\vec{b} onto \vec{a}), but not both onto some \vec{n} (see below).
 - The projected product $a_{\parallel} b$ is not the projected length a_{\parallel} unless $\vec{b} = \hat{b}$ is a unit vector.
 - projection**: $P_{\parallel} \vec{v} \equiv \vec{v}_{\parallel} = \hat{n} \hat{n} \cdot \vec{v}$ thus $P_{\parallel} = \hat{n} \hat{n}$ is an **operator** on vectors $P_{\parallel}: \vec{v} \mapsto \vec{v}_{\parallel}$.
 - i. Definition of projection: $P^2 = P$ because $P_{\parallel}^2 = \hat{n} \hat{n} \cdot (\hat{n} \hat{n} \cdot \vec{v}) = \hat{n} (\hat{n} \cdot \hat{n}) \hat{n} \cdot \vec{v} = \hat{n} \hat{n} \cdot \vec{v} = P_{\parallel}$
 - ii. Orthogonal projection: $P^T = P$ because $P_{\parallel}^T = (\hat{n} \hat{n}^T)^T = \hat{n} \hat{n} \cdot = P_{\parallel}$ (see below)
 - The defining properties of the dot product relate the geometric \rightarrow algebraic calculation:
 - scalar** valued: (product of lengths), polar \cdot polar, axial \cdot axial mirror-invariant
 - symmetric** (commutative): $\vec{a} \cdot \vec{b} = ab \cos \gamma = ba \cos(-\gamma) = \vec{b} \cdot \vec{a}$
 - bilinear** (distributive): $(\vec{a}_i \alpha_i) \cdot (\vec{b}_j \beta_j) = \alpha_i (\vec{a}_i \cdot \vec{b}_j) \beta_j = \alpha_i g_{ij} \beta_j$, where g_{ij} is the **metric tensor**.
 This is a direct property of projections, head-to-tail addition, and similar triangles
 - Example - law of cosines: the length² of $\vec{c} = \vec{a} - \vec{b}$ is $c^2 = a^2 - 2\vec{a} \cdot \vec{b} + b^2 = a^2 + b^2 - 2ab \cos \gamma$
 In the complex plane, the vector is $|c|^2 = c^* c = (a - b)^*(a - b) = |a|^2 + |b|^2 - 2\text{Re}(a^* b)$
 - An **orthonormal basis** $\hat{\mathbf{e}} = (\hat{x}, \hat{y}, \hat{z})$, where $\hat{x} \cdot \hat{x} = 1$, $\hat{x} \cdot \hat{y} = 0$, etc., forms a unit cube.
 - Components are easy to calculate in this basis: $v_x = \hat{x} \cdot \vec{v}$, $v_y = \hat{y} \cdot \vec{v}$, $v_z = \hat{z} \cdot \vec{v}$.
 This is trivial in $\hat{\mathbf{e}} = (\hat{x}, \hat{y}, \hat{z})$, but powerful for a rotated orthonormal basis.
 - Dot product also simplifies: $\vec{v} \cdot \vec{w} = (\hat{x} v_x + \hat{y} v_y) \cdot (\hat{x} w_x + \hat{y} w_y) = v_x w_x + v_y w_y = \mathbf{v}^T \mathbf{w} = v_i w_i$
 Thus in an orthonormal basis, $\vec{v} \cdot \vec{w}$ is a sum of projected products. If \vec{v} has only one component, ie. $\vec{v} = \hat{x} v$, this reduces to the geometric interpretation $\vec{v} \cdot \vec{w} = v w_x$.
 - Matrix notation - new concepts we need to treat the dot product.
 - Transpose** M^T flips a matrix M along its main diagonal: $\begin{pmatrix} a & b \\ c & d \end{pmatrix}^T = \begin{pmatrix} a & c \\ b & d \end{pmatrix}$.
 - It is almost invariably connected with the dot product as in the contraction $\mathbf{v}^T \mathbf{w}$.
 - Symmetry: M is **symmetric** if $M^T = M$ and **antisymmetric** if $M^T = -M$.
 - Hermitian conjugate** $M^\dagger = M^{*T}$ also conjugates complex vectors: $\begin{pmatrix} a & b \\ c & d \end{pmatrix}^\dagger = \begin{pmatrix} a^* & c^* \\ b^* & d^* \end{pmatrix}$.
 - The transpose ignores scalars $(a)^T = (a)$ and complex conjugate applies to each element.
 - The inverse M^{-1} and transposes M^T, M^\dagger all share a key property: $(AB)^T = B^T A^T$.
 - They are also all **involutions** (self-inverse): $(A^T)^T = A$.
 - Inner product**: $\mathbf{v}^T \mathbf{w} = \begin{pmatrix} v_x & v_y \end{pmatrix} \begin{pmatrix} w_x \\ w_y \end{pmatrix}$ is a contraction. Thus $(\mathbf{v}^T \mathbf{w})^T = (\mathbf{w})^T (\mathbf{v}^T)^T = \mathbf{w}^T \mathbf{v} = \mathbf{v}^T \mathbf{w}$.
 - Outer product**: $\mathbf{v} \mathbf{w}^T = \begin{pmatrix} v_x \\ v_y \end{pmatrix} \begin{pmatrix} w_x & w_y \end{pmatrix} = \begin{pmatrix} v_x w_x & v_x w_y \\ v_y w_x & v_y w_y \end{pmatrix}$, components of projection (above).
 - Identity matrix**: $\mathbb{I} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$ has three(!) interpretations related to the dot product:
 - Multiplicative identity operator: $\mathbb{I} \mathbf{v} = \mathbf{v}$, $\mathbf{v}^T \mathbb{I} = \mathbf{v}^T$, $\mathbf{v}^T \mathbb{I} \mathbf{w} = \mathbf{v}^T \mathbf{w}$.
 - Canonical basis transformation: $\hat{\mathbf{e}} = (\hat{e}_1 \hat{e}_2) = \mathbb{I}$ where $\hat{e}_1 = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$, $\hat{e}_2 = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$.
 - Orthonormal metric: $\hat{\mathbf{e}}^T \cdot \hat{\mathbf{e}} = (\hat{e}_1 \hat{e}_2)^T \cdot (\hat{e}_1 \hat{e}_2) = \begin{pmatrix} \hat{e}_1 \\ \hat{e}_2 \end{pmatrix} \cdot (\hat{e}_1 \hat{e}_2) = \begin{pmatrix} \hat{e}_1 \cdot \hat{e}_1 & \hat{e}_1 \cdot \hat{e}_2 \\ \hat{e}_2 \cdot \hat{e}_1 & \hat{e}_2 \cdot \hat{e}_2 \end{pmatrix} = \mathbb{I}$.



We can use these properties to derive the inner product form of the dot product:

$$\vec{v} \cdot \vec{w} = (\hat{e}v) \cdot (\hat{e}w) = (\hat{e}v)^T \cdot (\hat{e}w) = v^T (\hat{e}^T \cdot \hat{e}) w = v^T \mathbb{I} w = v^T w = v_x w_x + v_y w_y \quad \text{or} \quad \cdot \rightarrow^T$$

The same formalism applies to the metric of a non-orthonormal basis:

$$\vec{v} \cdot \vec{w} = (\hat{e}v) \cdot (\hat{e}w) = (\hat{e}v)^T \cdot (\hat{e}w) = v^T (\hat{e}^T \cdot \hat{e}) w = v^T \mathbb{g} w = (v_x \quad v_y) \begin{pmatrix} g_{xx} & g_{xy} \\ g_{yx} & g_{yy} \end{pmatrix} \begin{pmatrix} w_x \\ w_y \end{pmatrix}, \quad \cdot \rightarrow^T \mathbb{g}$$

where the metric $\mathbb{g} = \hat{e}^T \cdot \hat{e} = \begin{pmatrix} \vec{e}_1 \cdot \vec{e}_1 & \vec{e}_1 \cdot \vec{e}_2 \\ \vec{e}_2 \cdot \vec{e}_1 & \vec{e}_2 \cdot \vec{e}_2 \end{pmatrix} = \begin{pmatrix} g_{xx} & g_{xy} \\ g_{yx} & g_{yy} \end{pmatrix}$ characterizes the dot product.

- Example: in statistics the multivariable extension of $z = \frac{x-\mu}{\sigma}$ is the residual $\vec{\chi} = \vec{x} - \vec{\mu}$.

To generalize σ , we square σ to obtain the symmetric covariance matrix $\Sigma = \begin{pmatrix} \sigma_{11}^2 & \sigma_{12}^2 \\ \sigma_{21}^2 & \sigma_{22}^2 \end{pmatrix}$,

which appears in $\chi^2 = \vec{\chi} \cdot \vec{\chi} = \chi^T W \chi$, as the weight $W = \Sigma^{-1}$ called the [Mahalanobis metric](#).

- Index (implicit summation or Einstein) notation

- Indexed elements of a matrix: $M = \begin{pmatrix} m_{11} & m_{12} \\ m_{21} & m_{22} \end{pmatrix} \sim m_{ij}$ always row,col same as contractions

- Contraction: $\vec{v} \cdot \vec{w} = \sum_i v_i w_i = v_i w_i$ where the summation is implied (since it's always there!)

- Augmentation: $C = AB$ or $\begin{pmatrix} c_{11} & c_{12} \\ c_{21} & c_{22} \end{pmatrix} = \begin{pmatrix} a_{11} & a_{12} \\ a_{21} & a_{22} \end{pmatrix} \begin{pmatrix} b_{11} & b_{12} \\ b_{21} & b_{22} \end{pmatrix}$ or $c_{ik} = \sum_j a_{ij} b_{jk} = a_{ij} b_{jk}$,

which represents 4 equations, one for each i and j , for example $c_{12} = a_{1j} b_{j2} = a_{11} b_{12} + a_{12} b_{22}$, representing a contraction for each row of A and column of B , all augmented together.

The j is called a 'dummy index'. It can be replaced by any symbol and expands to a sum.

Be careful to use a different dummy index for each new summation (only 2 i 's or j 's).

- The **Kronecker delta** symbol $\delta_{ij} = \begin{cases} 1, & i = j \\ 0, & i \neq j \end{cases}$ represents the components of the identity matrix

$\mathbb{I} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} = \begin{pmatrix} \delta_{11} & \delta_{12} \\ \delta_{21} & \delta_{22} \end{pmatrix}$, and has the same property: $\mathbb{I}v = v$ or $\delta_{ij} v_j = v_i$. Thus the rule to simplify equations is to remove δ_{ij} and replace all of one dummy index with the other. If both i and j are repeated, like $\alpha = v_i \delta_{ij} w_j$, you can replace either to get $\alpha = v_i w_i = v_j w_j$.

- Conservation law of symbols: everything must match on the LHS and RHS of an equation:

1) units, 2) tensor rank (scalar, vector, operator, ...), 3) matrix dimension, 4) indices

With these rules, we can repeat the calculation of dot product in terms of components:

$$\vec{v} \cdot \vec{w} = (\hat{e}_i v_i) \cdot (\hat{e}_j w_j) = v_i (\hat{e}_i \cdot \hat{e}_j) w_j = v_i \delta_{ij} w_j = v_i w_i = v_x w_x + v_y w_y.$$

The same formalism applies to the metric of a non-orthonormal basis:

$$\vec{v} \cdot \vec{w} = (\hat{e}_i v_i) \cdot (\hat{e}_j w_j) = v_i (\hat{e}_i \cdot \hat{e}_j) w_j = v_i g_{ij} w_j = v_x g_{xx} w_x + v_x g_{xy} w_y + v_y g_{yx} w_x + v_y g_{yy} w_y.$$

The indices 'run over' either 1,2,3 or x,y,z for the labels of the different components.