

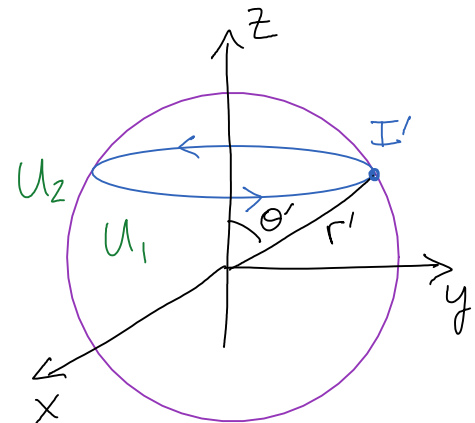
Magnetic Multipoles

Tuesday, April 15, 2014
11:08 AM

* Partial Differential Equation: Maxwell Eq's $\Rightarrow$ Laplace Eq.

$$\begin{aligned} \nabla \times \vec{H} &= \vec{J} \rightarrow 0 \Rightarrow \vec{H} = -\nabla U \\ \nabla \cdot \vec{B} &= 0 \Rightarrow -\nabla \cdot (\mu \nabla U) = 0 \\ B &= \mu H \end{aligned} \quad \text{or} \quad \boxed{\nabla^2 U = 0}$$

$$\begin{aligned} U_1 &= \sum_l (A_l r^l + B_l r^{-l-1}) P_l \\ U_2 &= \sum_l (C_l r^l + D_l r^{-l-1}) P_l \end{aligned} \quad \text{general solution}$$



* Boundary conditions - always derived the same way.

$$\begin{aligned} \nabla &\rightarrow \Delta \hat{n} \\ \frac{d}{d\epsilon} &\rightarrow \frac{d}{da} \end{aligned} \quad \int_{-\epsilon}^{\epsilon} d\vec{n} \left(\hat{n} \frac{\partial}{\partial n} + \dots \right) \times \vec{H} = \int_{-\epsilon}^{\epsilon} d\vec{n} \vec{K} \delta(n) \quad \int_a^b dx = \Delta x$$

$$\hat{n} \times \Delta \vec{H} = \vec{K}$$

$$\begin{aligned} \hat{n} \cdot \Delta \vec{B} &= 0 \\ \hat{n} \times \Delta \vec{H} &= \vec{K} \end{aligned} \quad \begin{aligned} B_{2n} - B_{1n} &= 0 \\ H_{2t} - H_{1t} &= K_s \end{aligned} \quad \boxed{\begin{aligned} -\mu_2 \frac{\partial U_2}{\partial n} + \mu_1 \frac{\partial U_1}{\partial n} &= 0 \\ -\frac{\partial U_2}{\partial \epsilon} + \frac{\partial U_1}{\partial \epsilon} &= K_s \end{aligned}} \quad \text{boundary conditions} \quad -\Delta U = I$$

* Application to azimuthal currents on a sphere:

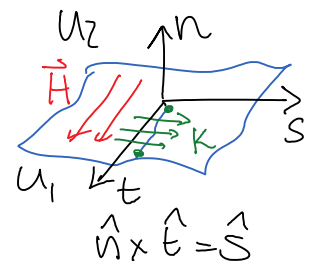
$$1) -\frac{\partial U_2}{\partial r} \Big|_{r'} + \frac{\partial U_1}{\partial r} \Big|_{r'} = 0$$

$$\sum_l [D_l (-l-1) r'^{-l-2} + A_l(l) r'^{l-1}] P_l = 0 \quad \text{basis}$$

$$D_l = -\frac{l}{l+1} r'^{2l+1} A_l$$

$$2) -\frac{\partial U_2}{r \partial \theta} \Big|_{r'} + \frac{\partial U_1}{r \partial \theta} \Big|_{r'} = K_\phi$$

$$\sum_l \left(-\frac{1}{r'} D_l r'^{l-1} + \frac{1}{r'} A_l r'^l \right) P'_l(\cos \theta) (-\sin \theta) = K_{\perp}$$



$$\sum_l -A_l r'^{l-1} \frac{2l+1}{l+1} \underbrace{P'_l(c_\theta)}_{\text{basis fn.}} s_\theta = K_\phi(\theta)$$

* Orthogonality relation for $P'_l(c_\theta) \cdot s_\theta = P'_l(x) \sqrt{1-x^2}$:

$$\text{note: } P_l^m = (-1)^m (1-x^2)^{m/2} \left(\frac{d}{dx} \right)^m P_l(x) \Rightarrow P'_l(x) = -(1-x^2)^{1/2} P'_l(x)$$

$$\int_{-1}^1 P'_l(x) P'_l(x) dx = \frac{2}{2l+1} \frac{(l+m)!}{(l-m)!} \rightarrow \frac{2l(l+1)}{2l+1} \text{ for } m=1 [P'_l(x)]$$

$$\text{thus } \langle P'_l | P'_l \rangle \equiv \int_{-1}^1 P'_l(x) P'_l(x) (1-x^2) dx = \langle P'_l | P'_l \rangle = \frac{2l(l+1)}{2l+1} \delta_{ll'}$$

$$\begin{aligned} \sum_l -A_l \frac{2l+1}{l+1} r'^{l-1} \int_0^\pi P'_l(c_\theta) P'_l(c_\theta) s_\theta^2 d\theta &= \int_0^\pi P'_l(c_\theta) K_\phi(\theta) s_\theta^2 d\theta \\ &= -A_{l'} \frac{2l'+1}{l'+1} r'^{l'-1} \cdot \frac{2l'(l'+1)}{2l'+1} \end{aligned}$$

$$A_l = -\frac{r'^{l-1}}{2l} \int_0^\pi P'_l(c_\theta) K_\phi(\theta) s_\theta^2 d\theta$$

$$D_l = +\frac{r'^{l+2}}{2(l+1)} \int_0^\pi P'_l(c_\theta) K_\phi(\theta) s_\theta^2 d\theta$$

* Reduction to a current loop

$$\text{Current loop: } K_\phi(\theta) = \frac{I'}{r'} \delta(\theta - \theta')$$

$$\int_0^\pi P'_l(c_\theta) K_\phi(\theta) s_\theta^2 d\theta = \frac{I' s_\theta^2}{r'} P'_l(c_{\theta'})$$

$$q_l \equiv A_l = -\frac{I' s_\theta^2}{2l r'^l} P'_l(c_\theta) \quad P_l \equiv 4\pi D_l = \frac{2\pi I' s_\theta^2 r'^{l+1}}{l+1} P'_l(c_\theta)$$

$$U_{\text{int}} = \sum_l q_l r^l P_l(c_\theta)$$

$$U_{\text{ext}} = \sum_l \frac{P_l}{4\pi r^{l+1}} P_l(c_\theta)$$

$$\frac{1}{r^{2l+1}} \quad \text{ex: } \frac{1}{r^3} \quad \text{external multipole}$$

$$\frac{1}{r^{2l+1}} \quad \text{ex: } \frac{1}{r^3} \quad \text{internal multipole}$$

Note: $q_l^E = \frac{P_l^E}{4\pi r^{2l+1}}$
[electric]

$q_l^M = \frac{-(l+1) P_l^M}{l 4\pi r^{2l+1}}$
[magnetic]

* external $[q_l]$
vs. internal $[p_l]$
moments.

* Magnetic dipole moment:

$$m = P_l = \frac{2\pi}{2} I' s^2 = I' A'$$

$$P_l' = x \Rightarrow P_l' = 1$$

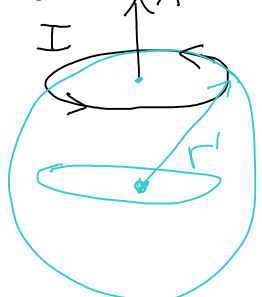
Electric

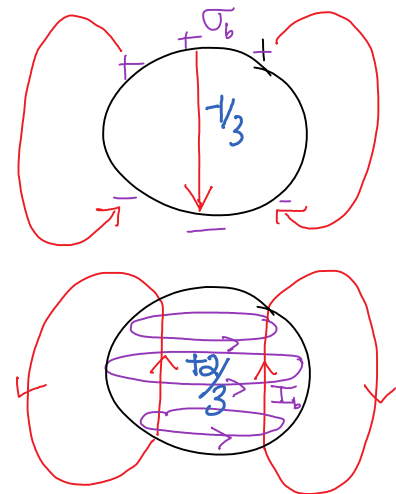
$$\vec{d} = \frac{q}{\epsilon_0} \vec{r}$$

$$\vec{p} = q \vec{d}$$

vs. Magnetic

$$\vec{A} = \pi s^2 \hat{z}$$

$$\vec{m} = I \vec{A}$$




* Dipole fields:

$$\vec{E} = \frac{1}{4\pi\epsilon_0} \frac{1}{r^3} [3\vec{p} \cdot \hat{r} \hat{r} - \vec{p}] - \left(\frac{1}{3}\right) \vec{p} \delta^3(\vec{r})$$

$$\vec{B} = \frac{\mu_0}{4\pi} \frac{1}{r^3} [3\vec{m} \cdot \hat{r} \hat{r} - \vec{m}] + \left(\frac{2}{3}\right) \vec{m} \delta^3(\vec{r})$$

* Legendre polynomials & derivatives:

$$P_0 = 1 \quad P_2 = \frac{1}{2}(3c^2 - 1) \quad P_4 = \frac{1}{8}(35c^4 - 30c^2 + 3)$$

$$P_1 = c \quad P_3 = \frac{1}{2}(5c^3 - 3c) \quad P_5 = \frac{1}{8}(63c^5 - 70c^3 + 15c)$$

$$P_0' = 0 \quad P_2' = 3c \quad P_4' = \frac{5}{2}(7c^3 - 3c)$$

$$P_1' = 1 \quad P_3' = \frac{3}{2}(5c^2 - 1) \quad P_5' = \frac{5}{8}(63c^4 - 52c^2 + 3)$$

* Helmholtz coil:

2 symmetric coils @ $\theta' = \theta_1$
 $\Rightarrow$ even multipoles $P'_l(c_0)$ vanish: $l=2, 4, 6, \dots$
 $l=1$: dipole moment

- we have one degree of freedom to cancel out a higher moment: $l=3$

$$q_3 = \frac{I_1 s_1^2}{2 \cdot 3 r^{13}} P'_3(c_1) = 0 \quad \text{if} \quad P'_3 = \frac{3}{2}(5c_1^2 - 1) = 0 \quad \boxed{c_1 = \frac{1}{5}}$$

* Maxwell coil:

add a centre ring with current I_0 $c_0=0$ r'
to the symmetric pair with I_1 ; c_1 also r'
- use I_0 to fix the magnitude, then we have two parameters I_1, c_1 to null q_3 and q_5

$$q_3 = \frac{-1}{6r^{13}} [I_0 \cancel{s_0^2} P'_3(c_0) + 2I_1 s_1^2 P'_3(c_1)] = 0$$

$$\frac{-1}{6r^{13}} \cdot \frac{3}{2} [I_0 (5c_0^2 - 1) + 2I_1 s_1^2 (5c_1^2 - 1)] = 0$$

$$-I_0 + 2I_1(1 - c_1^2)(5c_1^2 - 1) = 0 \quad [1]$$

$$q_5 = \frac{-1}{10r^{15}} [I_0 \cancel{s_0^2} P'_5(c_0) + 2I_1 s_1^2 P'_5(c_1)] = 0$$

$$\frac{-1}{10r^{15}} \frac{5}{8} [I_0 (63c_0^4 - 52c_0^2 + 3) + 2I_1 s_1^2 (63c_1^4 - 52c_1^2 + 3)] = 0$$

$$3I_0 + 2I_1(1 - c_1^2)(63c_1^4 - 52c_1^2 + 3) = 0 \quad [2]$$

$$3 \cdot [1] + [2]: 0I_0 + 2I_1(1 - c_1^2)[3(5c_1^2 - 1) + (63c_1^4 - 52c_1^2 + 3)] = 0$$

$$3(5c_1^2 + 21c_1^4 - 14c_1^2) = 9c_1^2(7c_1^2 - 3) = 0$$

$$C_1^2 = 3/7 \quad S_1^2 = 4/7 \quad I_0/I_1 = 2S_1^2(5C_1^2 - 1) = 8/7^2$$

check [2]: $3 \cdot \frac{8}{7^2} I_1 + 2 I_1 \left(\frac{4}{7}\right) \left(63 \frac{3^2}{7^2} - 52 \frac{3}{7} + 3\right)$

$$= \frac{3 \cdot 8 I_1}{7^2} [8 + (3 \cdot 3^2 - 14 \cdot 3 + 7)] = 0 \quad \checkmark$$

Solution: put a coil pair at angle $\cos \theta = \sqrt{3/7}$ and one in the center with $64/47 \times$ more current, all on the surface of a sphere.

* Gradient coil:

2 coils with opposite current

By symmetry odd multipoles vanish.

q_2 is the gradient, adjust $C_1 = \cos(\theta_1)$ to null q_4

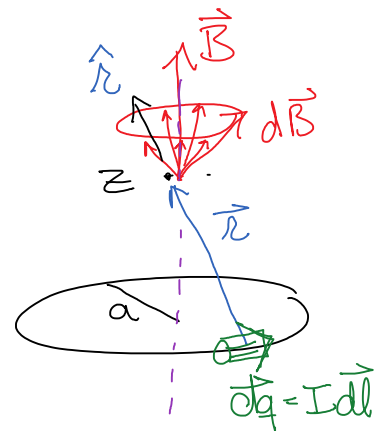
$$q_4 = \frac{2 I_1 S_1^2}{8 r^4} P_4'(C_1) = 0 \quad 7C_1^2 - 3 = 0 \quad C_1^2 = 3/7$$

This is the same configuration as the Maxwell coil without the center coil.

* External multipole:

$$\begin{aligned} \vec{H} &= \int \frac{d\vec{q} \times \vec{r}}{4\pi r^3} = \int_0^{2\pi} \frac{I(z\hat{s} + a\hat{z})a d\phi}{4\pi(z^2 + a^2)^{3/2}} \\ &= \frac{2\pi I a^2 \hat{z}}{4\pi(z^2 + a^2)^{3/2}} = \frac{\vec{m}}{2\pi r^3} \end{aligned}$$

where $\vec{m} = I \vec{A} = I \cdot \pi a^2 \hat{z}$



$$\begin{aligned} d\vec{l} \times \vec{r} &= a d\phi \hat{\phi} \times (z\hat{z} - a\hat{s}) \\ &= (z\hat{s} + a\hat{z}) a d\phi \end{aligned}$$

$$U_{int} = \sum_l q_l r^l P_l(\cos\theta)$$

$$U_{int} = q_1 z = \frac{mz}{2\pi r^3}$$

$$\vec{H} = -\nabla U = \frac{\vec{m}}{2\pi r^3} \quad \checkmark$$

$$q_l = \frac{-I' S_{\theta'}^2}{2l r^l} P_l'(\cos\theta)$$

$$q_1 = \frac{-I' \pi (r' s_{\theta})^2}{2\pi r^3} = \frac{-I' A}{2\pi r^3} = \frac{-m}{2\pi r^3} = \frac{-p_1}{2\pi r^3}$$

not 4π ? because of $-\frac{(l+1)}{2} = -2$