Electromagnetism in a Nutshell

* Maxwell's equations et al.

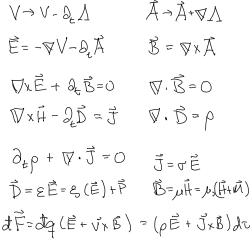
gauge potential field source $\lambda \stackrel{\triangle}{\rightarrow} (V, \hat{A}) \stackrel{\triangle}{\rightarrow} (\hat{E}, \hat{B}) \stackrel{\triangle}{\rightarrow} 0$ Ellu Maxwell eg.'s invariance

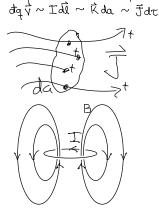
 $U \xrightarrow{d} (\vec{D}, \vec{H}) \xrightarrow{d} (\vec{p}, \vec{J}) \xrightarrow{d} O$ continuity

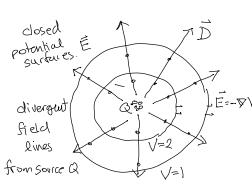
* Flux and Flow

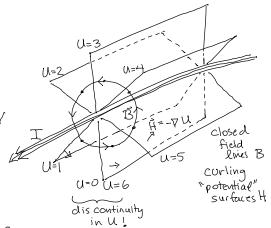
~ conserved currents 9, E, E, đa ~ 2dl ~ oda ~ pdz

~ integral equations \$= Q &+ & == 0 $\underline{\Phi}_{\mathcal{G}} = 0$ $\mathcal{E}_{H} - \mathcal{O}_{t} \underline{\Phi}_{D} = \underline{\mathcal{I}}$









* utility of D & B flux lines, E & H equipotential surfaces

~ flux through a surface $S = \Phi_B = \int_S \bar{B} \cdot d\bar{a} = \#$ of lines that poke through a surface S~ flow along a curve/path $P = \mathcal{E}_{\mathbf{E}} = \mathcal{L}_{\mathbf{E}} = \mathcal{L}_{\mathbf{E}} = \mathcal{L}_{\mathbf{E}}$ of surfaces that a path P pokes through

* potentials, from Helmholtz theorem, qV = potential energy $q\vec{A} = potential$ momentum" ~ transverse and longitudinal components $N^2\vec{V} = \vec{n} \cdot \vec{n} \cdot \vec{V} - \vec{n}_{\star} \vec{n}_{\star} \vec{V}$ $\nabla = \hat{N} \cdot \hat{\vec{Q}}_{n} + \nabla_{t}$

$$\vec{E} = -\vec{\nabla} \left(\vec{\nabla}^2 \vec{\nabla} \cdot \vec{E} \right) + \nabla \times \left(-\vec{\nabla}^2 \vec{\nabla} \times \vec{E} \right)$$

$$= -\vec{\nabla} \vec{\nabla} \vec{\nabla} \cdot \vec{E} + \nabla \times \left(-\vec{\nabla}^2 \vec{\nabla} \times \vec{E} \right)$$

$$\vec{B} = -\vec{\nabla} (\vec{\nabla}^2 \vec{\nabla} \cdot \vec{E}) + \nabla \times \left(-\vec{\nabla}^2 \vec{\nabla} \times \vec{E} \right)$$

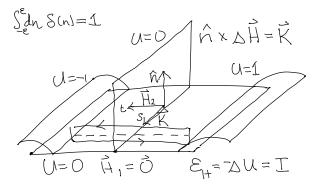
$$\vec{B} = -\vec{\nabla}(\vec{\nabla}^2 \vec{\nabla} \cdot \vec{B}) + \nabla \times (-\vec{\nabla}^2 \vec{\nabla} \times \vec{B})$$

$$\vec{B} = \nabla \times \vec{A}$$

* boundary conditions - integrate Maxwell's equations over the surface

 $\vec{J} \rightarrow \vec{k}$ $\int_{\epsilon}^{\epsilon} dn \frac{\partial}{\partial n} = \Delta$ $\int_{\epsilon}^{\epsilon} dn S(n) = \Delta$ $\nabla \rightarrow \hat{\mathsf{n}} \Delta$ $\rho \rightarrow \sigma$ Ezt=EIL RXAE=0 $D_{2n} - D_{n} = \sigma$ $\hat{N} \cdot \Delta \vec{D} = \sigma$ $\Phi_{0}=Q -\Delta \varepsilon \frac{\partial V}{\partial n}=\sigma'$ ~ magnetic

EH=I -AU=I H2E-H1 = Ks RXAH=K $\Phi_{B} = 0$ $\Delta \mu \frac{\partial U}{\partial n} = 0$ B_{2n}=B_{1n} Ñ·△B=O



~ surface current flows along U equipotential

~ U is a SOURCE potential

~ the current I=I2-II flows between any two equipotential lines U=II and U=I2

* electric magnetic dipoles and macroscopic equations (electric and magnetic materials)

$$\vec{A}(\vec{r}) = \frac{\mu_0}{4\pi} \frac{\vec{m} \times \vec{r}}{r^3} \quad \vec{m} = \oint \vec{L} d\vec{a} = \vec{L} \vec{a} \quad \times \Leftrightarrow \quad V_{dip}(\vec{r}) = \frac{1}{4\pi\epsilon_0} \vec{p} \cdot \vec{r} \quad \vec{p} \cdot \vec{r} \Rightarrow \frac{\mu_0 \vec{m}}{4\pi\epsilon_0} \quad \vec{m} = \frac{1}{\epsilon} \int d\vec{r} \vec{m} \cdot \vec{r} \cdot \vec{r$$

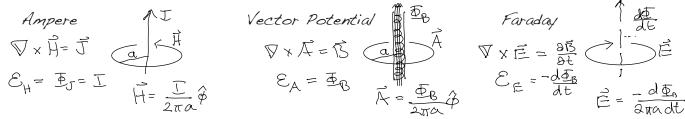
* dynamics of dipoles in fields (compare Electric and Magnetic)

$$\vec{F}_{e} = (\vec{p} \cdot \nabla) \vec{E} = \nabla (\vec{p} \cdot \vec{E}) = -\nabla W \quad (\nabla x \vec{E} = 0) \quad \vec{N}_{e} = \vec{p} \times \vec{E} \quad W = -\int N d\theta = -\vec{p} \cdot \vec{E}$$

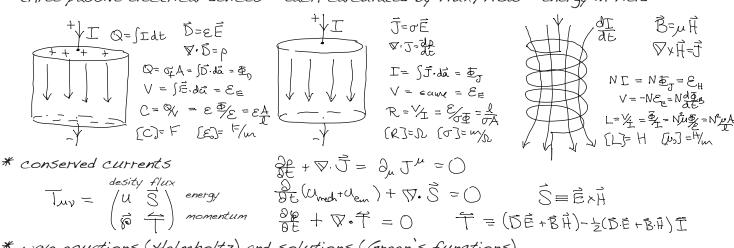
$$\vec{F}_{e} = (\vec{m} \times \nabla) \times \vec{B} = \nabla (\vec{m} \cdot \vec{B}) = -\nabla W \quad (\nabla \cdot \vec{B} = 0) \quad \vec{N}_{e} = \vec{p} \times \vec{E} \quad W = -\int N d\theta = -\vec{m} \cdot \vec{B}$$

* constitutive relations: magnetic susceptibility and permeability
$$\varepsilon_{0}\vec{E} = \vec{D} - \vec{P}$$
 $\vec{D} = \varepsilon_{0}(\vec{E} + \vec{P}) = \varepsilon_{0}(1 + \chi_{e})\vec{E} = \varepsilon_{0}\varepsilon_{r}\vec{E} = \varepsilon\vec{E}$ $\vec{P} - \vec{P} -$

* three Ampere-like laws - each can be solved using Stoke's theorem



* three passive electrical devices - each calculated by flux/flow = energy in field



* application of oblique boundary conditions: Fresnel equations

i)
$$D_{n} \circ D_{2n}$$
: $E_{I} - E_{R} = \beta E_{T}$ $E_{R} = \frac{\lambda - \beta}{\alpha + \beta} E_{I}$ $R = \left(\frac{\lambda - \beta}{\alpha + \beta}\right)^{2}$ $I = \frac{1}{2} \text{ ev } E^{2} \cos \theta$.

ii) $E_{1e} = E_{2e}$: $E_{I} + E_{R} = \lambda E_{T}$ $E_{I} = \frac{\lambda}{\alpha + \beta} E_{I}$ $T = \frac{4 \lambda \beta}{(\alpha + \beta)^{2}}$ $A = \frac{\cos \theta_{2}}{\cos \theta_{1}} \beta = \frac{\mathcal{X}_{2}}{\mathcal{X}_{1}}$

* guides: wave equation for longitudinal component, boundary conditions