

Square Well Potentials

* review of Wave mechanics

1) $\hat{H}\Psi = \hat{E}\Psi$ where $\hat{H} = \frac{-\hbar^2}{2m} \frac{\partial^2}{\partial x^2} + V(x)$ "hamiltonian"

2) $\int_{-\infty}^{\infty} |\Psi|^2 dx = 1$ "probability density"
 $\hat{E} = \frac{i\hbar \partial}{\partial t}$ TDSE or \underline{E} TISE "energy"

3) $\Psi(x), \frac{\partial \Psi}{\partial x}$ are a) continuous b) finite c) single-valued d) $\rightarrow 0$ as $x \rightarrow \pm\infty$
 "normalizable" "well-behaved"

* infinite square well

a) at $x < 0$ or $x > L$:

$$\frac{-\hbar^2}{2m} \frac{\partial^2 \Psi}{\partial x^2} + \infty \cdot \Psi(x) = E \cdot \Psi(x)$$

$\Psi(x) = 0!$

b) between $0 < x < L$

$$\frac{-\hbar^2}{2m} \frac{\partial^2 \Psi}{\partial x^2} = E \Psi(x)$$

$$\Psi(x) = e^{ikx} = \cos(kx) + i \sin(kx)$$

$$\Rightarrow A \cos(kx) + B \sin(kx)$$

c) boundary conditions

$$0 = \Psi(0) = A \quad (\text{LHS})$$

$$0 = \Psi(L) = B \sin(kL) \quad k_n L = n\pi$$

d) normalization:

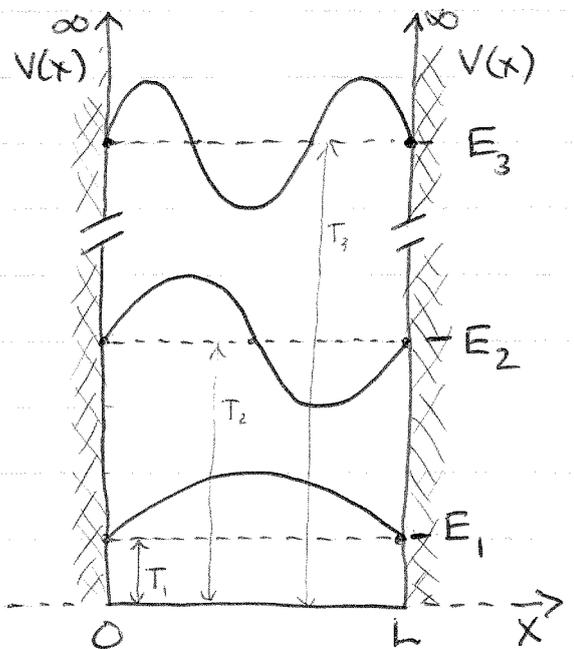
$$\int_{-\infty}^{\infty} |\Psi|^2 dx = \int_0^L B^2 \sin^2(k_n x) dx = B^2 \cdot \frac{L}{2} = 1$$

e) solution:

$$\boxed{\Psi_n(x) = \sqrt{\frac{2}{L}} \sin\left(\frac{n\pi x}{L}\right)}$$

time-dependence: $\Psi_n(x,t) = \underbrace{\sqrt{\frac{2}{L}} \sin\left(\frac{n\pi x}{L}\right)}_{\Psi(x)} e^{\underbrace{-i\omega_n t}}_{\phi(t)} = \frac{1}{2i} \sqrt{\frac{2}{L}} \left\{ e^{i(k_n x - \omega_n t)} - e^{-i(k_n x + \omega_n t)} \right\}$

standing "skipping rope" wave
 = counter-propagating travelling waves.



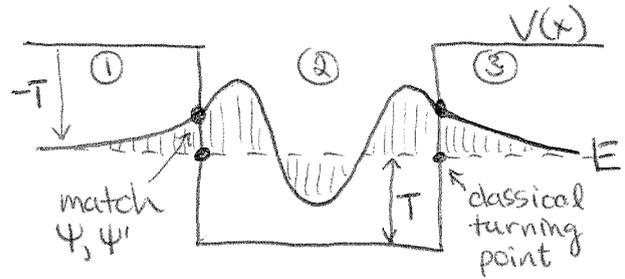
$$E_n = \frac{\hbar^2 k_n^2}{2m} = \frac{\hbar^2}{2m} \left(\frac{n\pi}{L}\right)^2 = \hbar \omega_n$$

$$p_n = \hbar k_n = \frac{\hbar n\pi}{L}$$

* finite square well

$$\frac{-\hbar^2}{2m} \frac{\partial^2 \psi}{\partial x^2} = [E - V(x)] \psi$$

$\frac{\hbar^2 k^2}{2m} = \frac{p^2}{2m}$ T



- find the solution for ψ in each region of constant V : ① ② ③

②: $\psi = e^{ikx}$ $k = \sqrt{2mT}/\hbar = \sqrt{2m(E-V)}/\hbar$

oscillatory, spatial frequency $k \sim \sqrt{T}$

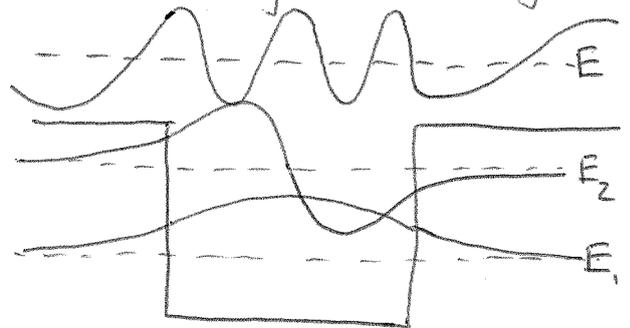
①, ③: $E - V(x) < 0$! classically forbidden.

$\psi = e^{-\kappa x}$ $k = \sqrt{2m(E-V)}/\hbar = \sqrt{-2m(V-E)}/\hbar$
positive
 $= i \sqrt{2m(V-E)}/\hbar \equiv i\kappa$

exponential decay in classically forbidden region
 "evanescent wave"

- final wave function:

- match up $\psi(x), \psi'(x)$ on the boundary.
- $\psi(x) \rightarrow 0$ as $x \rightarrow \pm \infty$
- still get a quantization of "n" oscillations in the well.
- only a finite number of bound states.



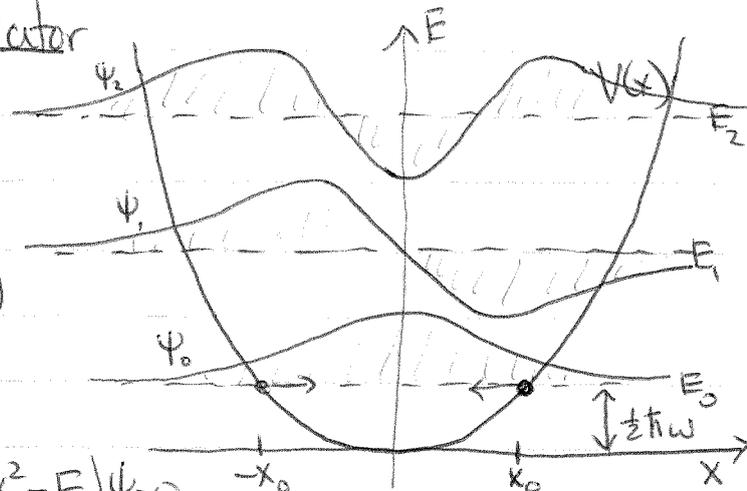
Simple Harmonic Oscillator

classical: $ma = F = -kx$

$$m\ddot{x} + kx = 0$$

$$-m\omega^2 + k = 0 \quad (x = e^{i\omega t})$$

$$\omega^2 = k/m$$



quantum: $-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} \psi + (\frac{1}{2} m\omega^2 x^2 - E)\psi = 0$

$$V = \frac{1}{2} kx^2 = \frac{1}{2} m\omega^2 x^2$$

ground state: $E_0 = \frac{1}{2} \hbar \omega$ at $x_0 = \sqrt{\frac{\hbar}{m\omega}}$

so let $\epsilon = E/E_0$ and $\xi = x/x_0$

$$\frac{\partial^2 \psi}{\partial \xi^2} + (\epsilon - \xi^2) \psi = 0$$

now, let $\psi = h(\xi) e^{-\frac{1}{2}\xi^2}$

$$\text{then } \frac{\partial^2 \psi}{\partial \xi^2} = \frac{\partial}{\partial \xi} \left[h' \cdot e^{-\frac{1}{2}\xi^2} + h \cdot (-\xi) \cdot e^{-\frac{1}{2}\xi^2} \right]$$

$$= h'' \cdot e^{-\frac{1}{2}\xi^2} + h'(-\xi) e^{-\frac{1}{2}\xi^2} + h'(-\xi) e^{-\frac{1}{2}\xi^2} - h e^{\frac{1}{2}\xi^2} + h \xi^2 e^{-\frac{1}{2}\xi^2}$$

$$= (h'' - 2\xi h' + (\xi^2 - 1)h) e^{-\frac{1}{2}\xi^2}$$

so $\boxed{h'' - 2\xi h' + (\epsilon - 1)h = 0}$

$$\epsilon = 2n + 1$$

differential equation for Hermite polynomials.

n	ϵ	$h(\xi)$	$\psi(x)$	E_n
0	1	1	$A_0 e^{-m\omega x^2/2\hbar}$	$\frac{1}{2} \hbar \omega$
1	3	2ξ	$A_1 \sqrt{\frac{m\omega}{\hbar}} x e^{-m\omega x^2/2\hbar}$	$\frac{3}{2} \hbar \omega$
2	5	$4\xi^2 - 2$	$A_2 \left(1 - \frac{2m\omega x^2}{\hbar}\right) e^{-m\omega x^2/2\hbar}$	$\frac{5}{2} \hbar \omega$
3	7	$8\xi^3 - 12\xi$	$A_3 \left(3x - 2\frac{m\omega}{\hbar} x^3\right) \sqrt{\frac{m\omega}{\hbar}} e^{-m\omega x^2/2\hbar}$	$\frac{7}{2} \hbar \omega$
4	9	$16\xi^4 - 48\xi^2 + 12$	$A_4 \dots$	$\frac{9}{2} \hbar \omega$