

University of Kentucky, Physics 404G  
Homework #11, Rev. A, [bonus]

1. Physical interpretation of the **vector potential  $\mathbf{A}$** . The quantity  $q\mathbf{A}$  has physical significance as potential momentum, just as  $qV$  is potential energy. The purpose of this problem is to explore the implications of this connection. For more background information see [Am. J. Phys. 64, 1368 \(1996\)](#).

a) Show that  $\partial\rho/\partial t + \nabla \cdot \mathbf{J} = 0$  (conservation of charge). In 4-vector notation, this is written  $\partial_\mu J^\mu = 0$ , where  $\partial = (\partial_t, \nabla)$ , and  $J = (\rho, \mathbf{J})$ .

b) For a continuous distribution  $\rho(\mathbf{r}, t)$  of charges in a velocity field  $\mathbf{v}(\mathbf{r}, t)$  so that  $\mathbf{J} = \rho\mathbf{v}$ , show that the *convective derivative* of  $\rho$  (the time derivative of  $\rho(\mathbf{r}(t), t)$  along the path  $\mathbf{r}(t)$  with velocity  $\mathbf{v} = d\mathbf{r}/dt$ ) is  $d\rho/dt = \partial\rho/\partial t + \mathbf{v} \cdot \nabla\rho$ , where  $\partial\rho/\partial t$  is the time derivative of  $\rho$  at a fixed position  $\mathbf{r}$  and  $\nabla\rho$  is the gradient of  $\rho$  at a fixed time. Use part a) to show that  $d\rho/dt + \rho\nabla \cdot \mathbf{v} = 0$ . What is the physical interpretation of  $\nabla \cdot \mathbf{v}$  for a velocity field?

c) Show that using the velocity-dependent potential  $U = q(V - \mathbf{v} \cdot \mathbf{A})$  in Lagrange's equation produces the Lorentz force:

$$\mathbf{F} \equiv \frac{d}{dt} \frac{\partial U}{\partial \mathbf{v}} - \frac{\partial U}{\partial \mathbf{r}} = q(\mathbf{E} + \mathbf{v} \times \mathbf{B}), \quad (1)$$

where  $\partial/\partial \mathbf{r} = \nabla = \hat{x}\partial_x + \hat{y}\partial_y + \hat{z}\partial_z$  and  $\partial/\partial \mathbf{v} = \nabla_{\mathbf{v}} = \hat{x}\partial_{v_x} + \hat{y}\partial_{v_y} + \hat{z}\partial_{v_z}$  are gradients with respect to the independent variables  $\mathbf{r}$  and  $\mathbf{v}$ , respectively. Use  $\mathbf{E} \equiv -\nabla V - \partial\mathbf{A}/\partial t$ , which includes time dependence. This generalizes the standard formula  $\mathbf{F} = -\partial U/\partial \mathbf{r} = -\nabla U$ . Thus the electromagnetic Lagrangian is  $L = T - U = \frac{1}{2}mv^2 - q(V - \mathbf{v} \cdot \mathbf{A})$ .

d) Calculate the *canonical momentum*  $\boldsymbol{\pi} \equiv \frac{\partial L}{\partial \mathbf{v}}$ . This generalized momentum includes contributions from both the particle and the magnetic field.

e) Calculate the *Hamiltonian*  $H \equiv \boldsymbol{\pi} \cdot \mathbf{v} - L$ . Express  $H$  in terms of  $\boldsymbol{\pi}$  and  $\mathbf{r}$ , not  $\mathbf{v}$ . This is the total energy in the system. Thus the Hamiltonian and canonical momentum  $(H, \boldsymbol{\pi})$  are the generalized energy and momentum of a system. Each includes a kinetic and potential term.

f) Show that  $\frac{dH}{dt} = \frac{d}{dt}(T + qV) = \frac{\partial}{\partial t}q(V - \mathbf{v} \cdot \mathbf{A}) = \frac{\partial U}{\partial t}$ , where  $\frac{\partial}{\partial t}$  acts on the potentials, but not  $\mathbf{r}$  or  $\mathbf{v}$ . Thus the total energy  $H$  is conserved if the generalized potential is time-independent.

g) Show that  $\frac{d\boldsymbol{\pi}}{dt} = \frac{d}{dt}(\mathbf{p} + q\mathbf{A}) = -\frac{\partial}{\partial \mathbf{r}}q(V - \mathbf{v} \cdot \mathbf{A}) = -\frac{\partial U}{\partial \mathbf{r}}$ , where again  $\frac{\partial}{\partial \mathbf{r}}$  acts on the potentials, but not  $\mathbf{v}$ . Thus  $\boldsymbol{\pi}$  is conserved if the generalized potential is position-independent.