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COMPLEX NUMBERS

• Imaginary numbers

 $(\chi^2 - 1) = (\chi + 1)(\chi - 1) = 0$ has 2 roots: $\chi = \pm 1$ $(\chi^2 + 1) = 2$ = 0 needs solutions of $\chi^2 = -1$ $\exists x \in \mathbb{R} \Rightarrow x^2 = -1$ ie $x = \pm \sqrt{-1}$ So define the new imaginary number $\dot{U} = \sqrt{-1}$ then $\chi = \pm \dot{U}$ are the solutions, and $\chi^2 + 1 = (\chi + i)(\chi - i)$ now factors into two binomials like χ^{2} - 1 Likewise, $\chi^2 + Q = \bigcirc \implies \chi = \pm \sqrt{Q} \mathring{c}$ so we have a whole imaginary line $iR = \{y_i | y \in R\}$ analogous to the real line Rey f Z=X+iy Real and imaginary numbers and add separately like vectors, so the direct sum $C = \{Z = X + iy \mid X, y \in \mathbb{R} \}$ forms a vector space called the "Complex plane" (both real and imaginary). $z_1 = (\hat{X}_1 + \hat{U}_1) \quad z_1 = X_1 + \hat{U}_1 \quad z_1 + z_2 = (X_1 + X_2) + \hat{U}(Y_1 + Y_2)$ $\tilde{\mathcal{Z}}_2 = (\hat{X} X_2 + \hat{Y} Y_2) \quad \mathcal{Z}_2 = X_2 + i Y_2 \qquad \mathcal{L} \mathcal{Z}_1 = (\mathcal{A} X_1) + i (\mathcal{A} Y_1)$

• Complex algebra:

Complex numbers have the extra structure $\dot{c}^2 = -1$ beyond standard vector arithmetic.

Thus
$$Z_1 \cdot Z_2 = (X_1 + i Q_1)(X_2 + i Q_2) = (X_1 X_2 - Y_1 Y_2) + i (X_1 Y_2 + X_2 Q_1)$$

This is similar to multiplying a vector by a matrix, except that a matrix has 4 degrees of freedom (each element), while the complex number only has two: $\begin{bmatrix}
\gamma \\
\eta
\end{bmatrix} = \begin{pmatrix}
\chi_z \\
\chi_z \\
\chi_z
\end{bmatrix}$

1. Z [identity] i.Z [ccw rotation by 90°]

Thus multiplication by 'i' is similar to the cross product and also generates rotations (HO2#2).

Either of these operations can be scaled by multiplying by a real number. Unlike matrix multiplication, these operations preserve angles and are called "conformal". Thus we can use complex functions to form various orthogonal coordinate systems.

• Conjugate and Polar coordinates:

 $Z^{\star} = (\chi - i \chi)$ sends 'i' to '-i'. $\varphi = (\chi - i \chi)$ sends 'i' The complex conjugate

It is used find the 'magnitude' or absolute value |z| of a complex number z

$$|Z|^2 = Z^* Z = (\chi - iq)(\chi + iq) = \chi^2 + q^2$$

The 'argument' or angle can be found using normal trig:

 $\tan \phi = \frac{y}{x}$ $x = \rho \cos \phi \quad y = \rho \sin \phi$

Euler's identity $e^{i\phi} = \cos\phi + i\sin\phi$ combines these expressions:

$$Z = \rho e^{i\phi} = \rho \cos\phi + i \rho \sin\phi = \chi + i \gamma$$

The product (and thus exponentials) are simpler in polar coordinates (also double-angle formulae):

 $Z_1 \cdot Z_2 = (p_1 e^{i\phi_1})(p_2 e^{i\phi_2}) = (p_1 p_2) e^{i(\phi_1 + \phi_2)}$ Multiplying by $e^{i\phi}$ rotates counterclockwise by the angle ϕ $P_1 e^{i\phi_1}$

and multiplying by \bigcap stretches by the factor \bigcap

The complex conjugate in polar coordinates is: $\rho e^{-i\phi} = \chi - i q$.

• The universality of complex numbers is captured by the Fundamental Theorem of Algebra: Any n-th order polynomial can be factored into a product of n binomials over $n_{i,j}$ and therefore has n complex roots (accounting for multiplicity). If the coefficients are real then the roots or either real or occur in complex-conjugate pairs.

$$P_{n}(z) = Q_{n} Z^{n} + Q_{ny} Z^{n-1} + \dots + Q_{2} Z^{2} + Q_{1} Z^{1} + Q_{0}$$
$$= Q_{n} (Z - Z_{1}) (Z - Z_{2}) - \dots + (Z - Z_{n})$$

Thus complex numbers solve a much larger problem than the original one posed.

Quaternions were invented by William Rowan Hamilton (also known for the Hamiltonian) to generalize rotations to 3 dimensions. Seeing that 1 and i represent the x̂ and ŷ, he tried unsuccessfully to add a third imaginary number j to represent ẑ. He later had an epiphany to use a 4-vector a1 + bi + cj + dk with a separate imaginary (i, j, k) for each (x̂, ŷ, ẑ) axis, and carved the famous inscription i² = j² = k² = ijk = -1 into the Brougham (Broom) Bridge, Dublin. The numbers are associative: (ij)k = i(jk), but not commutative. For example, multiplying by -k gives ijk(-k) = ij(-k²) = ij = k, but multiplying by ji on the left gives (ji)ijk = k = -ji, the negative of ij. Thus quaternions have the structure of cross products ij = k for x̂ × ŷ = ẑ and of (negative) dot products i² = j² = k² = -1. A rotation about x̂ in the (y,z)-plane:

$$e^{i\phi}j = (1\cos\phi + i\sin\phi)j = j\cos\phi + k\sin\phi$$

Other tricks were needed to keep i invariant by the rotation. But this math was too much for the 19^{th} century, so <u>Willard Gibbs</u> and <u>Oliver Heaviside</u>, independently distilled the structure of quaternions into the simpler dot and cross products of vector calculus. The use of (i, j, k) as unit vectors is a carry-over from quaternions. Similar contemporary developments were Hermann Grassmann's generalization of cross products to n-dimensions, and William Kingdom Clifford's combination of dot and cross products into a unified "geometric" product, with the same structure as Pauli or Dirac matrices.

ROTATIONS

- Rotations are linear operators (square matrices) that preserve the 'shape' of a set of vectors. They preserve lengths and angles between vectors, thus depends on the metric and are thus called 'orthogonal' transformations. Rotation matrices are also called orthogonal. HW#2 explores the similar structure of the imaginary i and the cross product ×, and this similar structure is use to 'generate' rotations.
- There are two ways to represent rotations:
 - a) Active rotations: physically rotate the vector in a fixed coordinate system for example, to describe a rotating object
 - b) Passive rotations: the physical vectors remains fixed, but the reference frame and basis vectors are rotated, giving the vector different components.

Both types of rotations look the same when transforming components



You can see that the same transformation of components involves rotations in different directions of the vector vs. the axes! To specify this relation, we need to include basis vectors in the passive formulation:

 $\vec{\nabla} = (\hat{X} \hat{Y} \hat{z}) \begin{pmatrix} V_X \\ V_Y \\ V_z \end{pmatrix} = \hat{e} \vee = \hat{e} \cdot \vee^{i} = (\hat{X}' \hat{Y}' \hat{z}') \begin{pmatrix} V_X' \\ V_Y' \\ V_z' \end{pmatrix} = \hat{e} \vee^{i} \vee^{i} = \hat{e} \cdot \vee^{i} \vee^{i} + \hat{e} \cdot \vee^{i} \vee^{i} = \hat{e} \cdot \vee^{i} \vee^{i} + \hat{e} \cdot \vee^{i} + \hat{e}$ if V'=RV, then $\vec{v}=\vec{e}\cdot\vec{R}\vec{v}=\vec{e}\cdot\vec{R}\vec{v}=\vec{e}\cdot\vec{V}$, so $\vec{e}=\vec{e}\cdot\vec{R}$

Notice the difference in primed vs unprimed and in left vs right multiplication of R. These differences are why v^i are called the contravariant components of \vec{v} .

We can remember this by writing way to write the transformations while still conserving "index height":

Let's examine each of these in terms of components in the first column of R:

- is the linear combination 1) $\hat{e} = \hat{e}' R$
- 2) Alternatively actively transform the components of the vector

 $\hat{\chi}' = (\hat{\chi}' \hat{\mathcal{Q}})$

 $\hat{\chi} = R \hat{\chi}' = \begin{pmatrix} C_{\phi} - S_{\phi} \\ S_{\phi} & C_{\phi} \end{pmatrix} \begin{pmatrix} I \\ O \\ C_{\phi} \end{pmatrix} = \begin{pmatrix} C_{\phi} \\ S_{\phi} \end{pmatrix} \text{ in the primed system,} \\ \text{which is the same as above.}$

 $\hat{\chi} = \hat{\chi}' c_{\phi} + \hat{\chi}' s_{h}$

 $\mathbb{R}^{\check{\iota}'}_{\check{J}}$ and $\hat{\mathcal{O}}_{\check{\iota}}, \bigvee^{j}$ so there is only one

 $\hat{e}_{i'}R^{i'}_{i} = \hat{e}_{i}$ $\vee^{i'} = R^{i'}_{i} \vee^{j}$

 $\hat{\chi} = \hat{\chi} C_{\phi} + \hat{V} S_{\phi}$

Orthogonal transformations are ones that 'preserve the metric' meaning that you get the same value of the dot product of two vectors whether or not you transform them.
 If this is true for the basis vectors, it will be true for the entire linear space.

$$g = \hat{e}^{\mathsf{T}} \cdot \hat{e} = (\hat{e}^{\mathsf{T}} R)^{\mathsf{T}} (\hat{e}^{\mathsf{T}} R) = R^{\mathsf{T}} \hat{e}^{\mathsf{T}} \cdot \hat{e}^{\mathsf{T}} R = R^{\mathsf{T}} g^{\mathsf{T}} R = g^{\mathsf{T}}$$

If the original basis is orthonormal, ie. g = I then the transformed basis is also g' = IThus $g = R^T g R$ so $R^T R = I^* r^{-1}$ or $R^T = R^{-1}$

Note the role of the transpose is to flip components to form dot products of individual columns.

$$\mathbb{R}^{\mathsf{T}} \mathbb{R} = \begin{pmatrix} \hat{\mathbf{x}} & \hat{\mathbf{x}} \\ \hline \mathbf{C}_{\boldsymbol{\phi}} & \mathbf{S}_{\boldsymbol{\phi}} \\ \hline \mathbf{S}_{\boldsymbol{\phi}} & \mathbf{C}_{\boldsymbol{\phi}} \\ \hline \mathbf{S}_{\boldsymbol{\phi}} & \mathbf{S}_{\boldsymbol{\phi}} \\ \hline \mathbf{S}_{\boldsymbol{\phi}} & \mathbf{S}_{\boldsymbol$$

• For orthogonal coordinate systems, the components can be found by taking dot products (the direction cosines)

$$R^{i'}_{j} = \hat{C}_{i'} \circ \hat{C}_{j} = cos \Theta_{i'}$$

$$R^{i}_{i'} R^{i'}_{k} = S^{i}_{k}.$$

• As an example, here is a derivation of the transformation matrix for the <u>Euler angles</u> of an arbitrary rotation in 3-d. You need 3 parameters: $2(\theta, \phi)$ to describe the direction of the new z-axis, and $1(\psi)$ to describe the position of the (x,y)-axes rotated about z. These are used in robotics, and also to describe the rotational motion of a rigid body, the last topic of this course.

An example is the roll (ψ), pitch (θ), and yaw (ϕ)

of an airplaine in a slightly different convention.

We obtain the full transformation matrix by composition three elementary rotations of (ϕ, θ, ψ) .

(wikipedia)

In the standard Euler z-y'-z'' convention, we rotate the coordinate system (ϕ) about the z-axis, then (θ) about the new y-axis, and then (ψ) about the new z-axis. Care is needed to multiply the rotations in the correct order:

a) Passive:
$$\hat{D} \hat{e}' = \hat{e} R_{\hat{z}\phi} \quad \hat{a}) \hat{e}'' = \hat{e}' R_{\hat{y}\theta} \quad \hat{a}) \hat{e}''' = \hat{e}'' R_{\hat{z}\psi}$$

$$\left(\left(\hat{e} \right) R_{\hat{z}\phi} = \hat{e}' \right) R_{\hat{y}\theta} = \hat{e}'' \right) R_{z\psi} = \hat{e}''' \quad ie \quad \hat{e}''' = \hat{e} R_{\hat{z}\phi} R_{\hat{y}\theta} R_{\hat{z}\psi}$$

a) Active: this time we must rotate (ψ) first about the z-axis before the z-axis has changed directions, since all rotations are done in the fixed lab frame. Likewise, we must rotate (θ) second while the y-axis is still in its original position, and finally the (ϕ) rotation about z.

$$\mathbb{R}_{\mathbb{V}} = \mathbb{R}_{\hat{\mathbb{Z}}_{\emptyset}} \left(\mathbb{R}_{\hat{\mathbb{Y}}_{\theta}} \left(\mathbb{R}_{\hat{\mathbb{Z}}_{\psi}} \mathbb{V} \right) \right) = \left(\mathbb{R}_{\hat{\mathbb{Z}}_{\emptyset}} \mathbb{R}_{\hat{\mathbb{Y}}_{\theta}} \mathbb{R}_{\hat{\mathbb{Z}}_{\omega}} \right) \mathbb{V}$$

Either way, the full rotation matrix is the same:

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$$\mathbb{R} = \begin{pmatrix} C_{\varphi} - S_{\varphi} \\ S_{\varphi} & C_{\varphi} \\ S_{$$