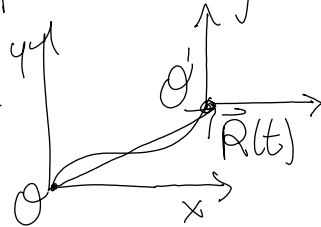


L38 Inertial Forces

Wednesday, December 4, 2019 07:30

* Galilean relativity - the foundation behind special relativity

let $\mathcal{O}xyz$ be an inertial frame where $\vec{F} = m\vec{a}$, and let $\mathcal{O}'x'y'z'$ be a frame with the same unit vectors $\hat{x}, \hat{y}, \hat{z}$ but $\mathcal{O}'$ is translated to $\vec{R}(t)$ w.r. $\mathcal{O}$.



$$\text{then } \vec{r}' = \vec{r} - \vec{R} \quad \vec{v}' = \vec{v} - \dot{\vec{R}} \quad \vec{a}' = \vec{a} - \ddot{\vec{R}}$$

$$\text{or } m\vec{a}' = \vec{F} - m\vec{A} \quad \vec{v} \quad \vec{A}$$

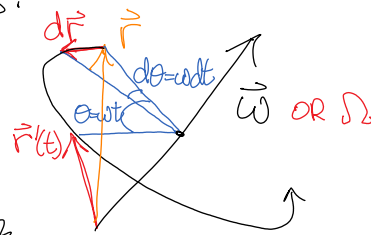
the acceleration of $\mathcal{O}'$ i.e. $\vec{A}$ is called an inertial force $\vec{F}_{\text{inertial}} = -m\vec{A}$ and acts just like a real force

that is the general equivalence principle behind general relativity.

* Rotating frames: we have already seen that $[\vec{\omega} \times]$ generates rotations about the $\vec{\omega}$ axis. Here's a simpler way to see this:

[arc length]

$$d\vec{s} = \vec{\omega} \times \vec{r} \quad \frac{d}{dt} \vec{r} = \vec{\omega} \times \vec{r}$$



this holds for all vectors: $\hat{e}_1, \hat{e}_2, \hat{e}_3$

let $\mathcal{O}'x'y'z'$ be rotating at velocity $\vec{\omega}(t)$ about $\mathcal{O}xyz$, with the same origin

$$\text{then } \vec{v} = \frac{d}{dt} \hat{e}_i r^i = \hat{e}_i \dot{r}^i + \underbrace{\dot{\hat{e}}_i}_{\vec{\omega} \times \hat{e}_i} r^i = \vec{v}' + \vec{\omega} \times \vec{r}'$$

$$\text{or } \left(\frac{d}{dt} \right)_{\mathcal{S}_0} = \left(\frac{d}{dt} \right)_{\mathcal{S}} + \vec{\omega} \times \dots \quad \text{note: } \vec{\omega} = \vec{\omega}' \text{ always.}$$

$$\vec{a} = \left(\frac{d}{dt} \right)_{\mathcal{S}_0} (\vec{v}' + \vec{\omega} \times \vec{r}') = \left(\frac{d}{dt} \right)_{\mathcal{S}} (\vec{v}' + \vec{\omega} \times \vec{r}') + \vec{\omega} \times (\vec{v}' + \vec{\omega} \times \vec{r}')$$

$$= \vec{a}' + \vec{\omega} \times \vec{r}' + (\vec{\omega} \times \dot{\vec{r}}' + \vec{\omega} \times \vec{v}') + \omega \times \omega \times \vec{r}'$$

* General transformation: $Oxyz$ to $O'x'y'z'$ w/ $\hat{e}_1 \hat{e}_2 \hat{e}_3$:

$$\vec{v} = \vec{v}' + \underbrace{\vec{\omega} \times \vec{r}'}_{\text{rotational}} + \underbrace{\vec{V}}_{\text{translation}}$$

$$\vec{a} = \vec{a}' + \underbrace{\vec{\alpha} \times \vec{r}'}_{\text{ang. accel.}} + \underbrace{2\vec{\omega} \times \vec{v}'}_{\text{Coriolis}} + \underbrace{\vec{\omega} \times \vec{\omega} \times \vec{r}'}_{\text{centrifugal}} + \underbrace{\vec{A}}_{\text{translation}}$$

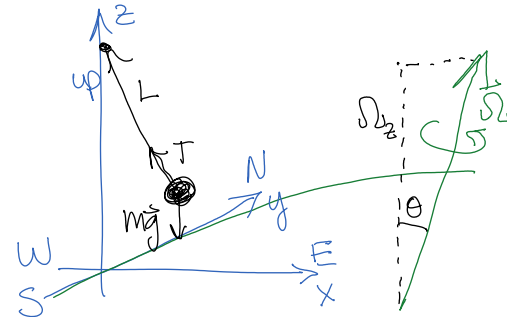
* Inertial forces:

$$m\vec{a}' = \vec{F} - \underbrace{m\vec{a}r\hat{\phi}}_{\text{tangential}} - \underbrace{2m\vec{\omega} \times \vec{v}'}_{\text{Coriolis}} + \underbrace{\frac{m\omega^2}{\rho'}\hat{\rho}}_{\text{centrifugal}} - m\vec{A}$$

* Example: Foucault Pendulum

$$m\vec{a}' = \vec{T} + \underbrace{m\vec{g}_0 - m\vec{\Omega} \times \vec{\Omega} \times \vec{r}}_{m\vec{g}} - 2m\vec{\Omega} \times \vec{v}$$

$T_z \cong T = mg \quad T_x = -\frac{mgx}{L} \quad T_y = -\frac{mgy}{L}$



let $\xi = x + iy \quad \ddot{\xi} = -\omega_0^2 \xi - 2i\Omega_z \dot{\xi}$

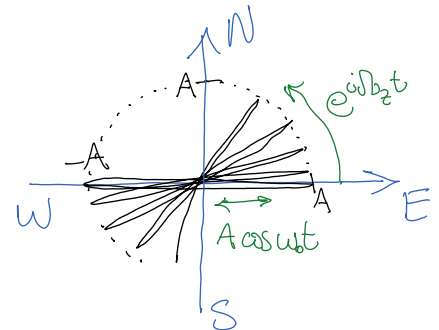
$= e^{i\alpha t} \quad \alpha^2 - 2i\Omega_z \alpha + \omega_0^2 = 0$

$\omega_0^2 = g/L \quad \Omega_z = \Omega \cos\theta$

$\alpha = \Omega_z \pm \sqrt{\Omega_z^2 + \omega_0^2} = \Omega_z \pm \omega_0$

$\xi = e^{i(\Omega_z \pm \omega_0)t} \rightarrow e^{i\Omega_z t} (C_1 e^{i\omega_0 t} + C_2 e^{-i\omega_0 t})$

$= e^{i\omega_0 t} (\underbrace{\tilde{A} \cos \omega t}_{\text{oscillation}} + \underbrace{\tilde{B} \sin \omega t}_{\text{precession}})$



$\tilde{A} = C_1 + C_2 \quad \tilde{B} = i(C_1 - C_2)$

$= A e^{i\phi_0} \quad \text{let } \phi_0 = 0$