

L40 Euler's equations

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* Review of Rotational Mechanics:

- Rotational dynamics: $\dot{\vec{L}} = \vec{N}$

$$\boxed{\vec{L} = \vec{r} \times \vec{p}}$$

$$\dot{\vec{L}} = \frac{d}{dt}(\vec{r} \times \vec{p}) = \vec{r} \times \vec{p}' + \vec{r} \times \vec{p} = \boxed{\vec{r} \times \vec{F} \equiv \vec{N}} \text{ or } \vec{r} \text{ or } \vec{r}'$$

- Inertia: $\vec{L} = \mathbf{I} \vec{\omega}$ $\mathbf{I} = m(r^2 \mathbf{I} - \vec{r} \otimes \vec{r}) = m(r^T r \mathbf{I} - \vec{r} \vec{r}^T)$

$$\vec{L} = \vec{r} \times \vec{p} = \vec{r} \times m(\vec{r} \times \vec{\omega}) = -m\vec{r} \times (\vec{r} \times \vec{\omega}) = m(r^2 - \vec{r} \cdot \vec{r}) \vec{\omega} \equiv \mathbf{I} \vec{\omega}$$

$$\mathbf{I} = m(r^T r \mathbf{I} - \vec{r} \vec{r}^T) = m \begin{pmatrix} x & y & z \\ y & z & x \\ z & x & y \end{pmatrix} \mathbf{I} - \begin{pmatrix} x & y & z \\ y & z & x \\ z & x & y \end{pmatrix} = m \begin{pmatrix} y^2+z^2 & -xy & -xz \\ -yx & z^2+x^2 & -yz \\ -zx & -zy & x^2+y^2 \end{pmatrix}$$

- Since "moment of inertia tensor" $\mathbf{I}^T = \mathbf{I}$, it is diagonalizable
- its eigenvectors/values are called "principle axes/momenta" for which $L_i = I_i \omega_i$

- Rotational Energy:

$$\text{lin: } \vec{p} = m\vec{v} \quad \vec{F} = \vec{p}' = m\vec{a} \quad \int dW = \vec{F} \cdot d\vec{x} = \frac{1}{2} m\vec{v}^2 = \frac{1}{2} \vec{p} \cdot \vec{v} = \frac{\vec{p}^2}{2m}$$

$$\text{rot: } \vec{L} = \mathbf{I} \vec{\omega} \quad \vec{N} = \dot{\vec{L}} = \mathbf{I} \vec{\alpha} \quad \int dW = \vec{N} \cdot d\vec{\theta} = \frac{1}{2} \vec{\omega} \cdot \mathbf{I} \cdot \vec{\omega} = \frac{1}{2} \vec{\ell} \cdot \vec{\omega} = \frac{1}{2} \vec{\ell} \cdot \mathbf{I} \cdot \vec{\ell}$$

- Summary: all formulas take on the same form, exchanging:

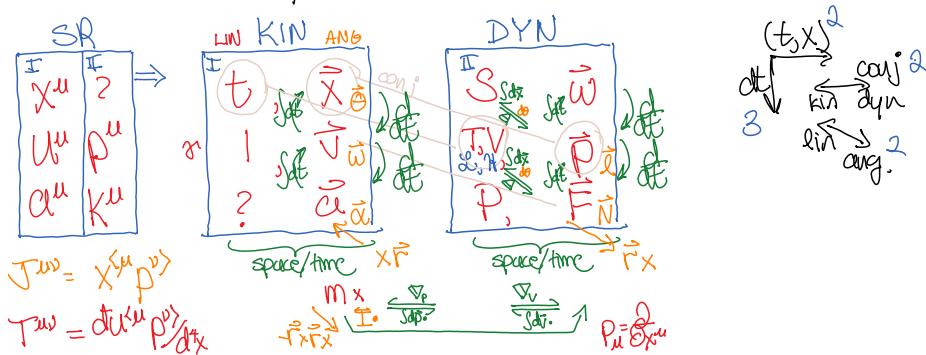
lin:	t	\vec{r}	\vec{v}	\vec{a}	m	\vec{F}	\vec{p}	T, V
	\parallel		$\cancel{1} \times \vec{r}$		$\sqrt{r \times r}$		$\cancel{\sqrt{r \times r}}$	\parallel
rot:	t	$\cancel{\vec{\theta}}$	$\vec{\omega}$	$\vec{\alpha}$	I	\vec{N}	$\vec{\ell}$	T, V

note: do not $\vec{\theta}$!

In particular, time and energy are the same;

spatial coordinates: $\vec{dr} = \vec{d\theta} \times \vec{r}$
 momentum, force : $\vec{l} = \vec{r} \times \vec{p}$ (backwards!)
 inertia: $\vec{\tau} = -m \cdot \vec{r} \times \vec{F}_x$

• Kinematics/Dynamics Map



* Euler's equations:

$\vec{N} = \dot{\vec{l}}$ (space) but $\vec{l} = \vec{I} \vec{\omega}$ and \vec{I} is only constant in the body frame:

$\vec{N} = \dot{\vec{l}} + \vec{\omega} \times \vec{l}$ (body) Euler's equations.

$\vec{I} \ddot{\vec{\omega}} = -\vec{\omega} \times \vec{I} \vec{\omega}$ in the absence of external forces.

$$\begin{aligned}\vec{I}_1 \dot{\omega}_1 &= (\vec{I}_2 - \vec{I}_3) \omega_2 \omega_3 && \text{(Euler equations)} \\ \vec{I}_2 \dot{\omega}_2 &= (\vec{I}_3 - \vec{I}_1) \omega_3 \omega_1 && \text{Newton's laws} \\ \vec{I}_3 \dot{\omega}_3 &= (\vec{I}_1 - \vec{I}_2) \omega_1 \omega_2 && \text{for rigid body}\end{aligned}$$

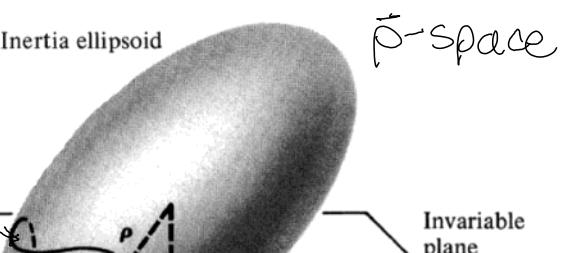
* Free Precession (from Goldstein, Poole, Safko, sec 5.6, p 200)

+ Poincaré construction

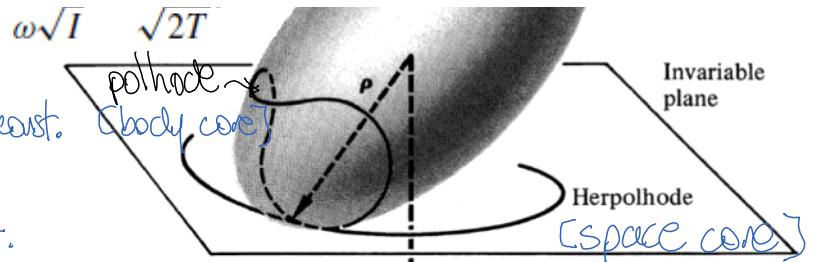
$$F(\rho) = \rho \cdot \mathbf{I} \cdot \rho = \rho_i^2 I_i, = 1 \quad \text{on the ellipsoid: } \rho = \frac{\omega}{\omega \sqrt{I}} = \frac{\omega}{\sqrt{2T}}$$

normal to ellipsoid: $2\mathbf{I} \cdot \boldsymbol{\omega}$

$$\sqrt{2}$$



normal to ellipsoid:
 $\nabla_{\rho} F = 2\mathbf{l} \cdot \boldsymbol{\rho} = \frac{2\mathbf{l} \cdot \boldsymbol{\omega}}{\sqrt{2T}} = \sqrt{\frac{2}{T}} \mathbf{L}$, const.
 projection along \mathbf{l} :
 $\frac{\boldsymbol{\rho} \cdot \mathbf{L}}{L} = \frac{\boldsymbol{\omega} \cdot \mathbf{L}}{L\sqrt{2T}} = \frac{\sqrt{2T}}{L}$ const.



the polhode rolls without slipping on the herpolhode L lying in the invariable plane.

+ Binet construction $I_3 \leq I_2 \leq I_1$,

$$T = \frac{L_x^2}{2I_1} + \frac{L_y^2}{2I_2} + \frac{L_z^2}{2I_3} \quad (\text{conserved})$$

$$\frac{L_x^2 + L_y^2 + L_z^2}{L^2} = 1 \quad (\text{conserved})$$

Energy ellipsoid
 Angular momentum sphere.

