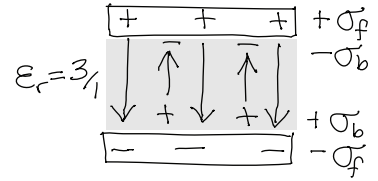


Section 4.2 - Polarization Fields

* review: dipole moment, polarization, forces, dielectrics

$$\begin{aligned} \vec{p} &= \int d\vec{q} \vec{r}' = q\vec{d} & \vec{N} &= \vec{p} \times \vec{E} & \sim \text{does } (\partial q)\vec{d} &= q(\partial \vec{d})? \\ d\vec{p} &= \vec{P} d\tau & U &= -\vec{p} \cdot \vec{E} & \sim \text{what about } d_1 &= -d_2? \\ \vec{F} &= q\vec{E} & \vec{F} &= (\vec{p} \cdot \nabla) \vec{E} & \sim \text{force on a quadrupole?} \end{aligned}$$

E, D, P all point down!



$$\epsilon_0 E = D - P = 3 - 2 = 1$$

* electric potential from polarization: bound charge

$$\begin{aligned} V &= \frac{1}{4\pi\epsilon_0} \int_{V'} \frac{d\vec{p}' \cdot \hat{r}}{r^2} = \frac{1}{4\pi\epsilon_0} \int_{V'} \vec{P}' d\tau' \cdot \nabla' \frac{1}{r} \\ &= \frac{1}{4\pi\epsilon_0} \left[\oint_{\partial V} \frac{\vec{P}' \cdot \hat{n}'}{r} da' + \int_{V'} \frac{-\nabla' \cdot \vec{P}'}{r} d\tau' \right] \\ &= \frac{1}{4\pi\epsilon_0} \left[\oint_{\partial V} \frac{\sigma'_b}{r} da' + \int_{V'} \frac{\rho'_b}{r} d\tau' \right] \end{aligned}$$

$$\begin{aligned} \nabla \frac{1}{r} &= \frac{1}{r^2} \nabla r = -\frac{\hat{r}}{r^2} \\ \nabla \frac{1}{r} &= \frac{1}{r^2} \nabla r = -\frac{\hat{r}}{r^2} \\ \nabla' \frac{1}{r} &= \frac{d}{d\vec{r}'} \frac{1}{r} = -\nabla \frac{1}{r} \\ \nabla \cdot \vec{P} &= \nabla \cdot \vec{P} + \vec{P} \cdot \nabla \frac{1}{r} \end{aligned}$$

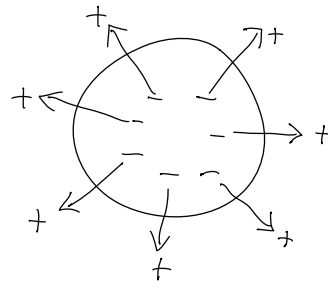
bound charge $\sigma_b = \vec{P} \cdot \hat{n}$ $\rho_b = -\nabla \cdot \vec{P}$ ~ uncancelled charge from overlapping dipoles in the polarized (dielectric) medium

* physical interpretation of bound charge
~ polarization forms "dipole chains"

$$\sigma_b = \frac{q_b}{A} = \frac{q\vec{d} \cdot \hat{n}}{Ad} = \vec{P} \cdot \hat{n}$$

~ divergence finds lone charge at end of each chain

$$\int_V \rho_b d\tau = -\oint_{\partial V} \vec{P} \cdot d\vec{a} = -\int_V \nabla \cdot \vec{P} d\tau \quad \rho_b = -\nabla \cdot \vec{P}$$



~ what fluxes have we considered so far?

~ what are the similarities and differences between \vec{E} and \vec{P} ?

~ why are they so similar?

* example 4.2 - bound charge and fields of a sphere with constant polarization \vec{P}

$$\begin{aligned} \vec{P} &= P\hat{z} \\ \rho_b &= -\nabla \cdot \vec{P} = 0 \\ \sigma_b &= \vec{P} \cdot \hat{n} = P \cos \theta \end{aligned}$$

$$\begin{aligned} V_1 - V_2 &= \frac{\sigma_b}{\epsilon_0}: \sum_l a_l \left(\frac{l}{R} \left(\frac{R}{r} \right)^{l+1} + \frac{l+1}{R} \left(\frac{R}{r} \right)^{l+2} \right) P_l(x) = \frac{P}{\epsilon_0} \cos \theta \\ \sum_l a_l \frac{2l+1}{R} P_l(x) &= \frac{P}{\epsilon_0} P_1(x) \quad a_l = \frac{PR}{3\epsilon_0} \delta_{l1} \end{aligned}$$

$$V_1 = \sum_l a_l \left(\frac{R}{r} \right)^l P_l(x)$$

$$V_1 = \frac{P}{3\epsilon_0} r \cos \theta = \frac{Pz}{3\epsilon_0} \quad \vec{E}_1 = -\nabla V_1 = \frac{-\vec{P}}{3\epsilon_0}$$

$$V_2 = \sum_l a_l \left(\frac{R}{r} \right)^{l+1} P_l(x)$$

$$V_2 = \frac{PR}{3\epsilon_0} \frac{R^2}{r^2} \cos \theta = \frac{\vec{P} \cdot \vec{r}}{4\pi\epsilon_0 r^3} \quad \text{where } \vec{P} = \frac{4}{3}\pi R^3 \vec{P}$$

$$V_1 = V_2:$$

$$\sum_l a_l P_l(x) = \sum_l a_l P_l(x)$$

$$\vec{E}_2 = -\nabla V_2 = \frac{\vec{P}}{4\pi\epsilon_0 r^3} (2\cos \theta \hat{r} + \sin \theta \hat{\theta}) = \frac{3\vec{P} \cdot \hat{r} - \vec{P}}{4\pi\epsilon_0 r^3}$$