

L14-TISE: Separation of variables

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- * Review: Time-Dependent Schrödinger equation (TDSE)
- "operator-ized dispersion relation"

$$\hat{H}\Psi = \hat{E}\Psi \quad \hat{H}(\hat{p}, \hat{x}) = \frac{\hat{p}^2}{2m} + \hat{V}(\hat{x})$$

$$\left(\frac{\hat{p}^2}{2m} + \hat{V}\right)\Psi = \hat{E}\Psi \quad \hat{p} = \hbar\hat{k} = -i\hbar\partial_x \quad \hat{E} = \hbar\omega$$

- we constructed this from eigenfunction plane waves of the standard wave equation when $\hat{V}=0$:

eigenfunction eigenvalue
 $\underbrace{\partial_x}_{\text{operator}} \underbrace{e^{ikx}}_{\langle x|k \rangle} = \underbrace{ik}_{\langle x|k \rangle} \underbrace{e^{ikx}}_{\langle x|k \rangle}$

eigenfunction eigenvalue
 $\underbrace{\partial_t}_{\text{operator}} \underbrace{e^{-i\omega t}}_{\langle t|\omega \rangle} = \underbrace{-i\omega}_{\langle t|\omega \rangle} \underbrace{e^{-i\omega t}}_{\langle t|\omega \rangle}$

"tensor" product
 $\Psi(x,t) = \underbrace{e^{ikx}}_{\langle x|k \rangle} \underbrace{e^{-i\omega t}}_{\langle t|\omega \rangle}$
 of eigenfunctions
 $\langle x,t|k,\omega \rangle$

- the general solution is a linear combination of (tensor) products of eigenfunctions: they form a basis of the solution space.

$$\Psi(x,t) = \int dk \Phi(k) e^{ikx-i\omega t} \quad \sim \int dk \underbrace{\Phi(k)}_{\text{only one } \omega \text{ per } k, \dots} \underbrace{|k\rangle|\omega\rangle}_{\text{dispersion relation:}}$$

- coefficient $\Phi(k)$ from initial conditions

$$\langle k|\Psi_0\rangle = \langle k|\int dk' \Phi(k') |k'\rangle|\omega'\rangle = \int dk' \Phi(k') \cdot 2\pi\delta(k-k') = 2\pi\Phi(k)$$

$$\Phi(k) = \frac{1}{2\pi}\langle k|\Psi_0\rangle = \int dx \frac{1}{2\pi} e^{-ikx} \Psi(x,0) \quad \hat{x} \cdot \vec{\nabla} = \partial_x$$

- * now apply the same technique to other potentials,
- no longer plane waves, what about time-dependence?

- apply separation of variables to obtain new eigenfunctions.

$$\Psi(x,t) = \psi(x) \cdot \varphi(t) \quad \hat{H}|\Psi\rangle = E|\Psi\rangle \quad \hat{E}|\varphi\rangle = E|\varphi\rangle$$

- if $\hat{H}$ is time-independent, $\hat{E}$ eigenfunctions already solved.

$$\langle t | \hat{E} | \varphi \rangle = i\hbar \partial_t \varphi(t) = E \varphi(t) \quad \text{let } E = \hbar\omega$$

$$\frac{d\varphi}{\varphi} = -i\omega dt \quad \varphi(x) = \cancel{\varphi_0} e^{-i\omega t} = \cancel{\varphi_0} e^{-iEt/\hbar}$$

- the physics is in the $\hat{H}$ equation,
the Time Independent Schrödinger Equation (TISE):

$$\hat{H}|\Psi\rangle = E|\Psi\rangle \quad \left(-\frac{\hbar^2}{2m} \nabla^2 + \hat{V} \right) \psi_n(x) = E_n \psi_n(x)$$

- Sturm-Liouville guarantees orthogonality: $\langle \psi_n | \psi_m \rangle = N_n^2 \delta_{nm}$
normalization constant

- general solution: $\Psi(x,t) = \sum_n \underbrace{c_n}_{\text{component}} \underbrace{\psi_n(x)}_{\text{basis function}} e^{-iE_n t/\hbar}$

- initial conditions: $\langle \psi_m | \Psi_0 \rangle = \sum_n c_n \langle \psi_m | \psi_n \rangle = N_m^2 c_m$
(normalized wavefunctions)

* formal solution to TDSE: (fancy representation of same steps!)

$$\hat{H}|\Psi\rangle = \hat{E}|\Psi\rangle = i\hbar \partial_t |\Psi\rangle \quad \hat{H} dt / i\hbar = d|\Psi\rangle / |\Psi\rangle$$

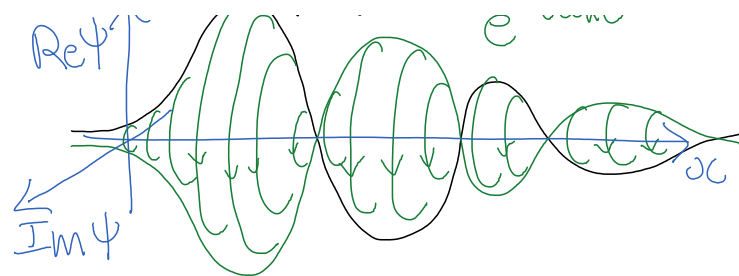
$$|\Psi_t\rangle = \underbrace{e^{\hat{H}t/\hbar}}_{\text{unitary } \hat{U}(t) \text{ time evolution operator}} |\Psi_0\rangle = \underbrace{\sum_n e^{E_n t/\hbar} |\psi_n\rangle}_{\text{spectral representation of } \hat{U}(t)} \underbrace{\langle \psi_n | \Psi_0 \rangle}_{c_n}$$

* properties of $\psi_n(x)$: (Griffiths)

a) stationary states:



- a) stationary states:
"skipping rope" function
like Bohr's orbitals



2-state system:

$$\Psi(x,t) = c_1 \psi_1(x) e^{iE_1 t/\hbar} + c_2 \psi_2(x) e^{-iE_2 t/\hbar} \quad \hbar \Delta\omega = E_2 - E_1$$

$$|\Psi(x,t)|^2 = |c_1|^2 \psi_1^2 + |c_2|^2 \psi_2^2 + 2|c_1^* c_2| \psi_1(x) \psi_2(x) \cos(\Delta\omega t)$$

exactly the Bohr frequency of transition radiation!

- b) Energy eigenstates:
definite total energy
 $\sigma_H^2 = \langle H^2 \rangle - \langle H \rangle^2 = 0$

$$\langle \hat{H} \rangle = \langle \psi_n | \hat{H} | \psi_n \rangle = \langle \psi_n | E_n | \psi_n \rangle = E_n$$

$$\langle f(\hat{H}) \rangle = f(E_n)$$

- c) They form a complete set of states,
a basis of the Hilbert space

$$\sum_n |\psi_n\rangle \langle \psi_n| = I$$