

# L25-Uncertainty Principle: Time

Friday, November 20, 2015 07:28

\* time dependence of expected value:

$$\begin{aligned} \frac{d}{dt} \langle Q \rangle &= \frac{d}{dt} \langle \psi | \hat{Q} | \psi \rangle \\ &= \left\langle \frac{\partial \psi}{\partial t} | \hat{Q} | \psi \right\rangle + \langle \psi | \frac{\partial \hat{Q}}{\partial t} | \psi \rangle + \langle \psi | \hat{Q} | \frac{\partial \psi}{\partial t} \rangle \\ &= \left\langle \frac{1}{i\hbar} \hat{H} \psi | \hat{Q} | \psi \right\rangle + \langle \psi | \frac{\partial \hat{Q}}{\partial t} | \psi \rangle + \langle \psi | \hat{Q} | \frac{1}{i\hbar} \hat{H} \psi \rangle \end{aligned}$$

$$\left\langle \frac{dQ}{dt} \right\rangle = \langle \psi | \frac{1}{i\hbar} [\hat{H}, \hat{Q}] + \frac{\partial \hat{Q}}{\partial t} | \psi \rangle$$

even true without  $\langle \psi | \dots | \psi \rangle$   
as an operator equation

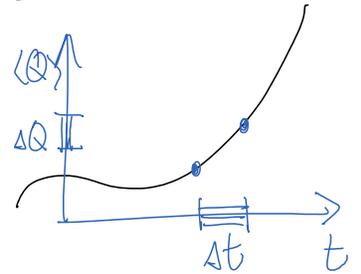
- if  $[\hat{H}, \hat{Q}] = 0$  then  $\hat{Q}$  is conserved ( $\langle Q \rangle$  constant)

\* how much time does it take  $\langle Q \rangle$  to change?

$$\sigma_H \sigma_Q \geq \left| \frac{1}{2i} \langle [\hat{H}, \hat{Q}] \rangle \right| = \left| \frac{1}{2i} \frac{\hbar}{i} \frac{d\langle Q \rangle}{dt} \right| = \frac{\hbar}{2} \left| \frac{d\langle Q \rangle}{dt} \right|$$

let  $\Delta E = \sigma_H$ ,  $\Delta t \equiv \frac{\sigma_Q}{\left| \frac{d\langle Q \rangle}{dt} \right|}$  time for  $\langle Q \rangle$  to change  $\sigma_Q$

then  $\Delta E \Delta t \geq \frac{\hbar}{2}$ .



- a pure energy state never changes.  $\sigma_H = 0 \Rightarrow \frac{\Delta Q}{\Delta t} = 0$

- sudden changes require infinite energy range  $\Delta t \rightarrow 0 \Rightarrow \sigma_H \rightarrow \infty$

\* Breit-Wigner resonances

[http://quantummechanics.ucsd.edu/ph130a/130\\_notes/node428.html](http://quantummechanics.ucsd.edu/ph130a/130_notes/node428.html)

- damped, undriven classical harmonic oscillator

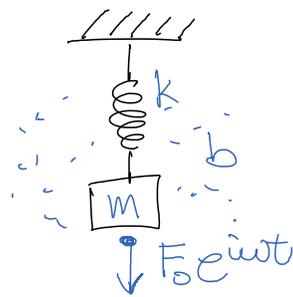
$$kx + b\dot{x} + m\ddot{x} = 0$$

let  $x = e^{\lambda t}$

$$\frac{1}{1}$$

$$kx + b\dot{x} + m\ddot{x} = 0 \quad \text{let } x = e^{\lambda t}$$

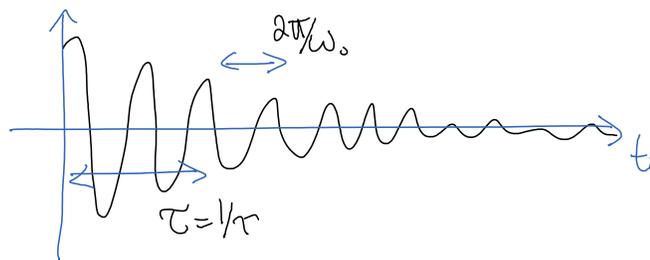
$$k + b\lambda + m\lambda^2 = 0 \quad \lambda = \frac{-b}{2m} \pm \sqrt{\left(\frac{b}{2m}\right)^2 - \frac{k}{m}}$$



$$x = (A_1 e^{i\omega t} + A_2 e^{-i\omega t}) e^{-\Gamma/2 t} \quad \text{where}$$

$$\omega_0^2 = \Gamma^2 + \frac{k}{m} \quad \text{resonant frequency}$$

$$\Gamma = \frac{b}{m} = 1/\tau \quad \text{lifetime}$$

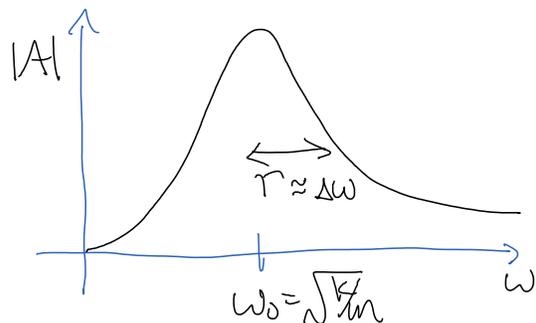


- drive it with an external frequency  $F_{ext} = F_0 e^{i\omega t}$

$$F_{ext} - kx - b\dot{x} = m\ddot{x}$$

$$F_0 e^{i\omega t} = (-m\omega^2 + i b\omega + k) A e^{i\omega t}$$

$$A = \frac{F_0}{(k - m\omega^2) + i b\omega} \quad |A|^2 = \frac{F_0^2}{(k - m\omega^2)^2 + b^2\omega^2}$$



- transition of energy  $E$  with lifetime  $\tau = 1/\Gamma$

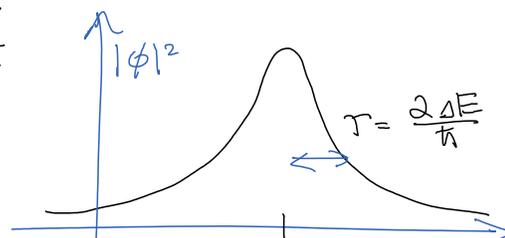
$$\Psi(x,t) = \psi(x) e^{-iEt/\hbar} e^{-\Gamma t/2} \quad \text{so that } P = \int dx |\Psi|^2 \sim e^{-\Gamma t}$$

$$\phi(x,\omega) = \mathcal{F}\{\Psi(x,t)\} = \int dt e^{i\omega t} [\psi(x) e^{-iEt/\hbar} e^{-\Gamma t/2}]$$

$$= \int_0^\infty dt e^{i(\omega - \omega_0 + i\Gamma/2)t} = \frac{i}{(\omega - \omega_0) + i\Gamma/2} \quad E = \hbar\omega_0$$

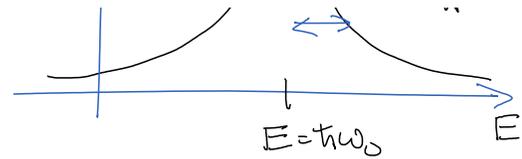
$$\text{thus } I(\omega) = \|\phi(x,\omega)\|^2 = \frac{1}{(\omega - \omega_0)^2 + \frac{\Gamma^2}{4}}$$

"Breit-Wigner" line shape. Fullwidth  $\Gamma$



uncertainty principle wave packet. minimum

note:  $\Delta E \cdot \Delta t = \hbar \frac{\Gamma}{2} \cdot \tau = \hbar \frac{\Gamma}{2}$



- the longer a transition takes, the more oscillations go into the wave packet, and the better defined the frequency is:  $\Delta t \Delta \omega \gg \frac{1}{2}$  (just like  $\Delta x \Delta k \gg \frac{1}{2}$ )

