

L10-Propagation of a Wave Packet: Dispersion

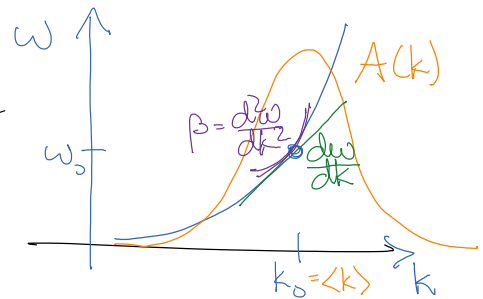
Friday, September 16, 2016 08:00

* How does a wave packet evolve? Add the time dependence:

$e^{ikx} \rightarrow e^{i(kx - \omega t)}$ of each frequency component, and

Taylor expand the dispersion relation

$$\omega(k) \approx \underbrace{\omega_0}_{\omega(k_0)} + \left. \frac{d\omega}{dk} \right|_{k_0} (k - k_0) + \frac{1}{2} \left. \frac{d^2\omega}{dk^2} \right|_{k_0} (k - k_0)^2$$



$$kx - \omega t = k_0 x + (k - k_0)x - \omega(k)t$$

$$\approx k_0(x - \underbrace{\frac{\omega_0}{k_0}}_{v_\phi} t) + (k - k_0)(x - \underbrace{\left. \frac{d\omega}{dk} \right|_{k_0}}_{v_g} t) - \frac{1}{2} (k - k_0)^2 \underbrace{\left. \frac{d^2\omega}{dk^2} \right|_{k_0}}_{\beta} t$$

$$\Psi(x, t) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} dk A(k) e^{i(kx - \omega(k)t)} \quad (\text{adding the time dependence})$$

$$\approx e^{i k_0(x - v_\phi t)} \cdot \underbrace{\frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} dq A(k_0 + q) e^{iq(x - v_g t) - \frac{i}{2} q^2 \beta t}}_{\text{envelope (wave packet)}} \quad \text{where } q = (k - k_0)$$

phase translation broadening

The ripples travel at the phase velocity.

The "particle" is represented by the envelope, a function of x_t , which thus travels at the group velocity.

The dispersion dephases the Fourier components on the fringe of the wave packet, which causes it to spread out.

* Example: Gaussian wave packet: recall the Fourier Transform pair

$$\begin{aligned} \Psi(x) &= \underbrace{\frac{1}{\sqrt{2\pi}\Delta x}}_{N_x = \text{norm}} \underbrace{e^{-\frac{1}{4}(\frac{x-x_0}{\Delta x})^2}}_{\text{envelope}} \underbrace{e^{ik_0 x}}_{\text{phase}} \\ &= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} dk A(k) e^{ikx} \end{aligned} \quad \Leftrightarrow \quad \begin{aligned} &\underbrace{\frac{1}{\sqrt{2\pi}\Delta k}}_{N_k = \text{norm}} \underbrace{e^{-\frac{1}{4}(\frac{k-k_0}{\Delta k})^2}}_{\text{envelope}} \underbrace{e^{-i\Delta k x}}_{\text{phase}} = A(k) \\ &\frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} dx \Psi(x) e^{-ikx} = \end{aligned}$$

$x \leftrightarrow k$

DIFFERENCE!
↑
DIFFERENCE!

Center the momentum: $A(k_0 + q) = N_k e^{-\frac{1}{4} \frac{q^2}{\Delta k^2} - \frac{i q x_0}{\Delta k} - i k_0 x_0}$

gaussian shift global phase

$$\begin{aligned} \Psi(x, t) &= e^{i(\phi - k_0 x_0)} \cdot \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} dq N_k e^{-q^2 \left(\frac{1}{4\Delta k^2} + \frac{i}{2} \beta t \right)} e^{iq(x - \underbrace{(x_0 + v_g t)}_{x_0(t)})} \\ &= e^{i \underbrace{\phi}_{\text{phase}}} e^{-ik_0 x_0} \underbrace{\frac{1}{\sqrt{2\pi}\Delta x(t)}}_{\text{envelope}} e^{-\frac{1}{4} \left(\frac{x - x_0(t)}{\Delta x(t)} \right)^2} \quad \text{where} \end{aligned}$$

apply the above F.T. pair:
 $x_0(t) = x_0 + v_g t$
 $\Delta x^2(t) = \Delta x^2(0) + \frac{1}{2} \beta t$

$$= e^{i k_0 (x - v_g t)} e^{-i k_0 x_0} \frac{1}{\sqrt{2\pi} \Delta x(t)} e^{-\frac{1}{4} \frac{(x - x_0(t))^2}{\Delta x^2(t)}} \quad \text{where} \quad \begin{aligned} x_0(t) &= x_0 + v_g t \\ \Delta x^2(t) &= \Delta x^2(0) + \frac{1}{4} \beta t \end{aligned}$$

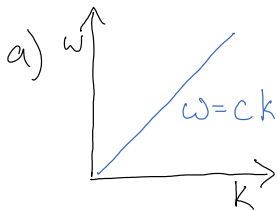
To understand the dispersion (spreading out of the wave packet), calculate the probability density $P(x) = |\Psi(x)|^2$

$$\begin{aligned} |\Psi(x, t)|^2 &= \frac{1}{\sqrt{2\pi} \Delta x(t)} \left| e^{-\frac{1}{4} \frac{(x - x_0)^2}{\Delta x^2(t)}} \right|^2 & |e^{z/2}|^2 &= e^{z/2} \cdot e^{z^*/2} = e^{\frac{z+z^*}{2}} = e^{2 \operatorname{Re}[z]/2} = e^{\operatorname{Re}[z]} \\ &= \frac{1}{\sqrt{2\pi} \Delta x(t)} e^{-\frac{1}{2} \frac{(x - x_0)^2 \operatorname{Re}[\Delta x^2(t)]}{(\Delta x^2(t))^2}} = \frac{1}{\sqrt{2\pi} \Delta x(t)} e^{-\frac{1}{2} \frac{(x - x_0)^2}{\Delta x^2(0) + \frac{1}{4} \beta t \Delta x^2(0)}} \end{aligned}$$

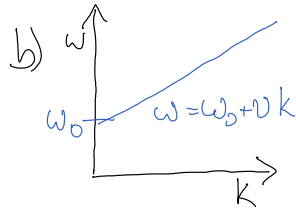
is a Gaussian distribution of width $\Delta x(t) = \Delta x(0) \sqrt{1 + \frac{1}{4} \beta t^2 / \Delta x^4(0)}$, which spreads out in time! (dispersion)

Thus the dispersion constant $\beta = \frac{d^2 \omega}{dk^2}$, i.e. the second derivative of the dispersion relation, is actually what causes dispersion. The bending of light at different angles is due to frequency-dependent impedance $Z(k) = \sqrt{\mu \epsilon} = \mu v = \mu v_k$, which is related to dispersion.

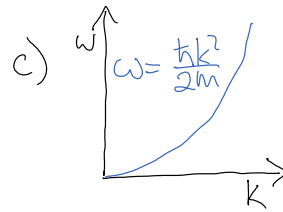
* Types of dispersion relations:



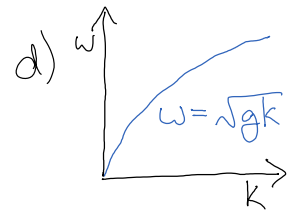
nondispersive medium (E&M vacuum)
constant $v_\phi = v_g = c$



constant v_g
with extra global phase factor $e^{i\omega_0 t}$



free particle:
 $2v_\phi = v_g = \frac{\hbar k}{m} = v_p$
ripples lag behind.



gravity waves:
 $\frac{1}{2} v_\phi = v_g = \frac{1}{2} \sqrt{\frac{g}{k}}$
ripples advance