

## L25-Operators: Stretches

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### \* Geometry of Stretches & Rotations

$$S v = \lambda v$$

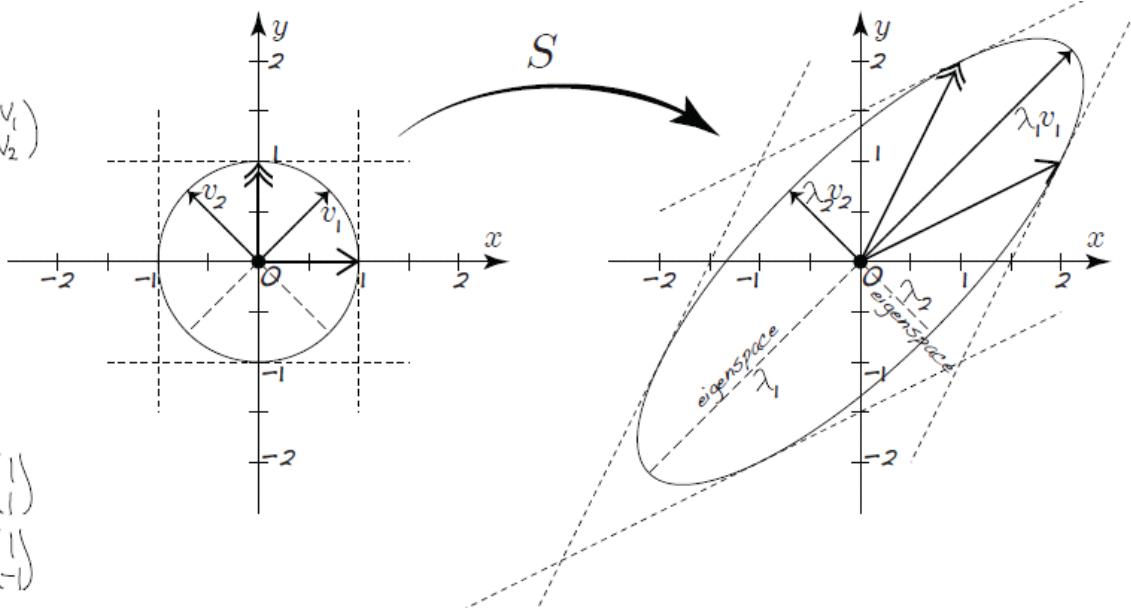
$$\begin{pmatrix} 2 & 1 \\ 1 & 2 \end{pmatrix} \begin{pmatrix} v_1 \\ v_2 \end{pmatrix} = \lambda \begin{pmatrix} v_1 \\ v_2 \end{pmatrix}$$

$$\begin{pmatrix} 2 & 1 \\ 1 & 2 \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \begin{pmatrix} 2 \\ 1 \end{pmatrix}$$

$$\begin{pmatrix} 2 & 1 \\ 1 & 2 \end{pmatrix} \begin{pmatrix} 0 \\ 1 \end{pmatrix} = \begin{pmatrix} 1 \\ 2 \end{pmatrix}$$

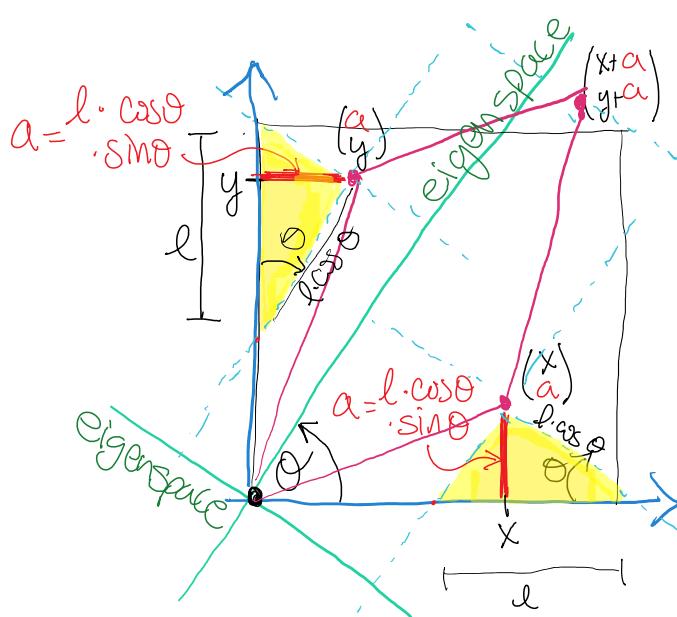
$$\begin{pmatrix} 2 & 1 \\ 1 & 2 \end{pmatrix} \begin{pmatrix} 1 \\ 1 \end{pmatrix} = \begin{pmatrix} 3 \\ 3 \end{pmatrix} = 3 \begin{pmatrix} 1 \\ 1 \end{pmatrix}$$

$$\begin{pmatrix} 2 & 1 \\ 1 & 2 \end{pmatrix} \begin{pmatrix} 1 \\ -1 \end{pmatrix} = \begin{pmatrix} 3 \\ -3 \end{pmatrix} = 1 \begin{pmatrix} 1 \\ -1 \end{pmatrix}$$



- principle axes of an ellipse are normal (perpendicular)
- symmetric matrices are pure stretches

$$(x-x_0, y-y_0) \begin{pmatrix} \frac{1}{a^2} & \varepsilon \\ \varepsilon & \frac{1}{b^2} \end{pmatrix} \begin{pmatrix} x-x_0 \\ y-y_0 \end{pmatrix} = 1$$



$$M = \begin{pmatrix} x & a \\ a & y \end{pmatrix}$$

- any shape of trapezoid can be made by flattening a square  $\Delta$  along  $\theta$  line
- we still need to rotate to get an arbitrary trapezoid  $M = RS$

## \* Eigensystem notation

- if  $M^{\dagger} = M$  (Hermitian) then there exists a complete set of eigenvalues  $\lambda_i$  and eigenvectors  $\vec{v}_i$  such that  $M\vec{v}_i = \lambda_i \vec{v}_i$  and  $\vec{v}_i^T \vec{v}_j = \delta_{ij}$ .
- we can put  $\vec{v}_i$  together to form a unitary transformation matrix  $V = (\vec{v}_1 | \vec{v}_2 | \dots | \vec{v}_n)$ ,  
 so  $MV = VD$      $D = \begin{pmatrix} \lambda_1 & & \\ & \lambda_2 & \\ & & \ddots \end{pmatrix}$  (diagonalization)
 
$$\begin{pmatrix} m_{xx} & m_{xy} \\ m_{yx} & m_{yy} \end{pmatrix} \begin{pmatrix} v_{1x} & v_{2x} \\ v_{1y} & v_{2y} \end{pmatrix} = \begin{pmatrix} v_{1x} & v_{2x} \\ v_{1y} & v_{2y} \end{pmatrix} \begin{pmatrix} \lambda_1 & \\ & \lambda_2 \end{pmatrix}$$
- $V$  is unitary:  $V^T V = (\vec{v}_1 \vec{v}_2 \vec{v}_3) \cdot (\vec{v}_1 \vec{v}_2 \vec{v}_3)^T = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} = I$   
 (orthonormal)
- $V$  is also closed:  
 $VV^T = (\vec{v}_1 \vec{v}_2) \begin{pmatrix} \vec{v}_1 \\ \vec{v}_2 \end{pmatrix} \cdot = \underbrace{\vec{v}_1 \vec{v}_1^*}_{P_1} + \underbrace{\vec{v}_2 \vec{v}_2^*}_{P_2} = I$
- $D$  = matrix elements of  $M$  in basis  $V$ :  

$$D = V^T M V = \langle v | M | v \rangle$$
  
 $I = V^T V$
- eigenvalue decomposition of  $M$ :  

$$M = V D V^T = \sum_i \lambda_i P_i$$
  
 $I = V V^T = \sum_i P_i$

## \* Characterization of linear operators

- a) ANY linear function has a singular value decomposition (SVD)
- $$M = \underbrace{R S}_{\text{symmetric}} = \underbrace{R}_{\text{eig.}} \underbrace{V W V^T}_{U^T U = I_B} = U W V^T \quad \text{where}$$

$$M = \underbrace{RS}_{\substack{\text{rotation} \\ A \rightarrow B}} = \underbrace{R}_{\substack{\text{symmetric} \\ 2 \text{ rotations}}} \underbrace{VWV^T}_{\substack{\text{diagonal} \\ \text{combined rotation}}} = UWV^T \quad \text{where } \begin{aligned} U^T U &= I_B \\ V^T V &= I_A \\ W &\text{ diagonal} \end{aligned}$$

eigenvalues  $\lambda_i$  of  $W$ : singular values  
 $V, U$  are eigen-basis of  $A, B$  respectively.

- b) ANY linear operator has  $n$  complex eigenvalues:  $\lambda_i$   
(Fundamental Theorem of Algebra on characteristic equation)
- $\gamma_i$  = geometric multiplicity = number of each eigenvalue
  - $m_i$  = algebraic multiplicity = dimension of eigen space
  - if  $m_i = \gamma_i$  then the matrix is diagonalizable
  - otherwise it still has a Jordan decomposition  $M = V \begin{pmatrix} \lambda_1 & & \\ & \ddots & \\ & & \lambda_n \end{pmatrix} V^{-1}$   
note that  $V$  is not necessarily orthogonal or normal!  $\rightarrow V^{-1}$ , not  $V^T$
  - this is a similarity transform

- c) ANY normal matrix  $M$  (ie.  $M^T M = M M^T$ )  
has a unitary diagonalization  
(orthonormal eigenvectors).  $V^T V = I$

- d) [anti] Hermitian operators have [imaginary] real eigenvalues.

$$\lambda_i^* v_i^T v_j = v_i^T H^T v_j = v_i^T H v_j = v_i^T v_j \lambda_j$$

$$\text{if } i=j : (\lambda_i - \lambda_i^*) v_i^T v_i = 0 \rightarrow \lambda_i \text{ real}$$

$$\text{if } \lambda_i \neq \lambda_j : (\lambda_i - \lambda_j^*) v_i^T v_j = 0 \rightarrow v_i^T v_j = 0 \text{ (unitary)}$$

\* What does  $^T$  have to do with operators?

Symmetric / antisymmetric vs. Symmetric / orthogonal decomposition

~ recall complex numbers  $U = \rho + i\phi$   $\rho^* = \rho$   $(i\phi)^* = -i\phi$

$$e^U = e^{\rho+i\phi} = r e^{i\phi} \quad |e^{i\phi}|^2 = e^{-i\phi} e^{i\phi} = e^{i0} = 1$$

~ similar behaviour of symmetric / antisymmetric matrices

$$M = \begin{pmatrix} a & b \\ c & d \end{pmatrix} = \begin{pmatrix} a & (b+c)/2 \\ (b+c)/2 & d \end{pmatrix} + \begin{pmatrix} 0 & (b-c)/2 \\ (c-b)/2 & 0 \end{pmatrix} = T + A$$

$$K = e^M = I + M + \frac{1}{2!}M^2 + \frac{1}{3!}M^3 + \dots = e^{T+A} \neq e^T e^A \text{ in general.}$$

M	arbitrary matrix
T	symmetric
A	antisymmetric
S	symmetric
R	orthogonal

$$S = e^T = e^{VWV^{-1}} = V e^W V^{-1} \quad R = e^A \quad R^T R = (e^A)^T e^A = e^{A^T + A} = e^0 = I$$

$$\det(e^{\lambda_1} e^{\lambda_2} \dots) = e^{\lambda_1} \cdot e^{\lambda_2} \dots = e^{\lambda_1 + \lambda_2 + \dots} = e^{\text{tr}(\lambda_1 \lambda_2 \dots)} \quad \det e^A = e^{\text{tr} A} = e^0 = I$$

if  $e^{T+A} = e^T \cdot e^A$  [ie T, A commute], then

$K = S \cdot R$  polar decomposition