

L32-Uncertainty Principle: Time

Friday, November 20, 2015 07:28

* time dependence of expected value:

$$\frac{d}{dt} \langle Q \rangle = \frac{d}{dt} \langle \psi | \hat{Q} | \psi \rangle$$

$$= \langle \frac{\partial \psi}{\partial t} | \hat{Q} | \psi \rangle + \langle \psi | \frac{\partial \hat{Q}}{\partial t} | \psi \rangle + \langle \psi | \hat{Q} | \frac{\partial \psi}{\partial t} \rangle$$

$$= \langle \frac{1}{i\hbar} \hat{H} \psi | \hat{Q} | \psi \rangle + \langle \psi | \frac{\partial \hat{Q}}{\partial t} | \psi \rangle + \langle \psi | \hat{Q} | \frac{1}{i\hbar} \hat{H} \psi \rangle$$

$$\langle \frac{dQ}{dt} \rangle = \langle \psi | \frac{1}{i\hbar} [\hat{H}, \hat{Q}] + \frac{\partial \hat{Q}}{\partial t} | \psi \rangle \quad \text{even true without } \langle \psi | \dots | \psi \rangle$$

as an operator equation

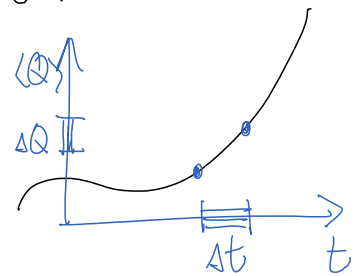
- if $[\hat{H}, \hat{Q}] = 0$ then $\hat{Q}$ is conserved ($\langle Q \rangle$ constant)

* how much time does it take $\langle Q \rangle$ to change?

$$\sigma_H \sigma_Q \geq | \frac{1}{2i} \langle [\hat{H}, \hat{Q}] \rangle | = | \frac{1}{2i} \hbar \frac{d\langle Q \rangle}{dt} | = \frac{\hbar}{2} | \frac{d\langle Q \rangle}{dt} |$$

let $\Delta E = \sigma_H$, $\Delta t \equiv \frac{\sigma_Q}{|d\langle Q \rangle/dt|}$ time for $\langle Q \rangle$ to change σ_Q

then $\Delta E \Delta t \geq \frac{\hbar}{2}$



- a pure energy state never changes. $\sigma_H = 0 \Rightarrow \frac{\Delta Q}{\Delta t} \rightarrow \infty$

- sudden changes require infinite energy range $\Delta t \rightarrow 0 \Rightarrow \sigma_H \rightarrow \infty$

* Breit-Wigner resonances http://quantummechanics.ucsd.edu/ph130a/130_notes/node428.html

- damped, undriven classical harmonic oscillator

$$kx + b\dot{x} + m\ddot{x} = 0$$

$$\text{let } x = e^{\lambda t}$$

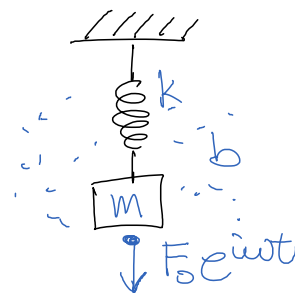
$$\frac{1}{1 \pm i}$$

$$kx + b\dot{x} + m\ddot{x} = 0$$

$$\text{let } x = e^{\lambda t}$$

$$k + b\lambda + m\lambda^2 = 0$$

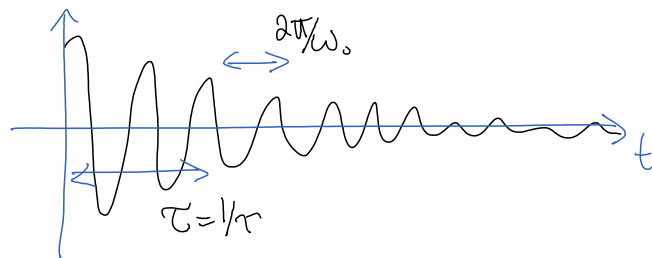
$$\lambda = \frac{-b}{2m} \pm \sqrt{\left(\frac{b}{2m}\right)^2 - \frac{k}{m}}$$



$$x = (A_1 e^{i\omega t} + A_2 e^{-i\omega t}) e^{-\tau/2 t} \quad \text{where}$$

$$\omega_0^2 = \tau^2 + \frac{k}{m} \quad \text{resonant frequency}$$

$$\tau = \frac{b}{m} = 1/\tau \quad \text{lifetime}$$

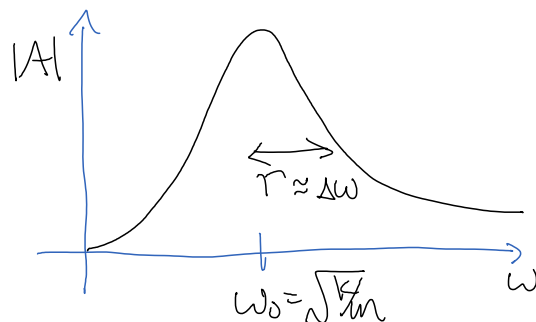


- drive it with an external frequency $F_{\text{ext}} = F_0 e^{i\omega t}$

$$F_{\text{ext}} - kx - b\dot{x} = m\ddot{x}$$

$$F_0 e^{i\omega t} = (-m\omega^2 + ib\omega + k) A e^{i\omega t}$$

$$A = \frac{F_0}{(k - m\omega^2) + ib\omega} \quad |A|^2 = \frac{F_0^2}{(k - m\omega^2)^2 + b^2\omega^2}$$



- transition of energy E with lifetime $\tau = 1/\tau$

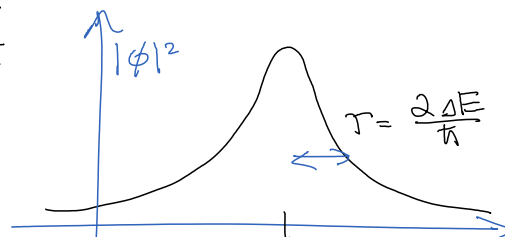
$$\Psi(x, t) = \psi(x) e^{-iEt/\hbar} e^{-\tau t/2} \quad \text{so that } P = \int dx |\Psi|^2 \sim e^{-\tau t}$$

$$\phi(x, \omega) = \mathcal{F}\{\Psi(x, t)\} = \int dt e^{i\omega t} [\psi(x) e^{-iEt/\hbar} e^{-\tau t/2}]$$

$$= \int_0^\infty dt e^{i(\omega - \omega_0 + i\frac{\Gamma}{2})t} = \frac{i}{(\omega - \omega_0) + i\frac{\Gamma}{2}} \quad E = \hbar\omega_0$$

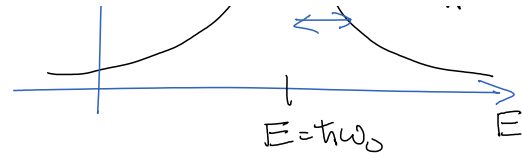
$$\text{thus } I(\omega) = \|\phi(x, \omega)\|^2 = \frac{1}{(\omega - \omega_0)^2 + \frac{\Gamma^2}{4}}$$

"Breit-Wigner" line shape. Fullwidth Γ



such higher are sharp. further

note: $\Delta E \cdot \Delta t = \hbar \frac{\Gamma}{2} \cdot \tau = \hbar \frac{\Gamma}{2}$



- the longer a transition takes, the more oscillations go into the wave packet, and the better defined the frequency is: $\Delta t \Delta \omega \gg \frac{1}{2}$ (just like $\Delta x \cdot \Delta k \gg \frac{1}{2}$)

