

## L35-Spherical Harmonics and Angular Momentum

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\* Angular Momentum - recall 3 nested S.-L. problems

$$\hat{T} = \frac{\hat{P}_m^2}{2m} = \frac{-\hbar^2}{2m} \nabla^2 = \frac{-\hbar^2}{2m} \cdot \frac{1}{r^2} \left( \frac{\partial}{\partial r} r^2 \frac{\partial}{\partial r} + \frac{1}{\sin\theta} \left( \frac{\partial}{\partial \theta} \sin\theta \frac{\partial}{\partial \theta} + \frac{1}{\sin\theta} \frac{\partial^2}{\partial \phi^2} \right) \right)$$

$$= \underbrace{\left( -\hbar^2 \cdot \frac{1}{r} \frac{\partial^2}{\partial r^2} r \right)}_{\hat{T}_r = \frac{\hat{P}_r^2}{2m}} / 2m + \left[ \underbrace{\frac{-\hbar^2}{\sin\theta} \left( \frac{\partial}{\partial \theta} \sin\theta \frac{\partial}{\partial \theta} \right)}_{\hat{L}_\theta^2} + \underbrace{\frac{1}{\sin\theta} \left( -\hbar^2 \frac{\partial^2}{\partial \phi^2} \right)}_{\hat{L}_\phi^2} \right] / 2mr^2$$

\*  $\hat{L}_z = -i\hbar \partial_\phi \quad \hat{L}_z e^{im\phi} = \hbar m e^{im\phi}$

periodic boundary conditions:  $\Psi(0) = \Psi(2\pi) \quad \Psi'(0) = \Psi'(2\pi)$   
 $\Rightarrow m = 0, \pm 1, \pm 2, \dots \in \mathbb{Z}$

\*  $\hat{L}^2 = \hbar^2 \cdot \frac{-1}{\sin\theta} \left( \frac{\partial}{\partial \theta} \frac{\sin^2\theta}{\sin\theta} \frac{\partial}{\partial \theta} - \frac{m^2}{\sin\theta} \right) \quad \hat{L}^2 \Psi(\theta) = \hbar^2 l(l+1) \Psi(\theta)$

$$= \hbar^2 \left( \frac{\partial}{\partial x} (1-x^2) \frac{\partial}{\partial x} + \frac{m^2}{1-x^2} \right)$$

$$\hat{L}^2 P_l^{(m)}(x) = \hbar^2 l(l+1) \underbrace{P_l^{(m)}(x)}_{\text{Associated Legendre functions}}$$

let  $x = \cos\theta = c_\theta$   
 $dx = -\sin\theta d\theta$   
 $\sqrt{1-x^2} = \sin\theta = s_\theta$

\* if  $m=0$  (azimuthal symmetry,  $L_z=0$ )  $P_l^{(0)}(x) = P_l(x)$  Legendre Polynomial

$$P_l(x) \equiv \frac{1}{2^l l!} \left( \frac{d}{dx} \right)^l (x^2 - 1)^l \quad \text{Rodrigues' formula} \quad P_0 = 1 \quad P_1 = x$$

$$(l+1)P_{l+1}(x) = (2l+1)x P_l(x) - l P_{l-1}(x) \quad \text{Bonnet's recursion formula}$$

$$\int_1^1 dx P_l(x) P_m(x) = \delta_{lm} \frac{2}{2l+1} \quad \text{Orthogonality \& normalization (S.-L.)}$$

\* if  $m > 0$   $P_l^m(x) = \underbrace{(1-x^2)^{m/2} \left( \frac{d}{dx} \right)^{lm}}_{S_l^{(m)}} P_l(x)$  some authors add factor  $(-1)^m$   $m > 0$   
"Condon-Shortley phase"

$$(l-m+1)P_{l+1}^m(x) = (2l+1)x P_l^m(x) - (l+m)P_{l-1}^m(x) \quad \text{Recursion formula}$$

$$P_m^m(x) = \underbrace{(-1)^m}_{\text{Condon-Shortley phase}} (2m-1)!! (1-x^2)^{m/2} \quad P_{m+1}^m(x) = x (2m+1) P_m^m(x) \quad \text{first 2 terms to start recursion.}$$

$$\int_1^1 dx P_l^{(m)}(x) P_l^{(m)}(x) = \delta_{mm} \frac{2}{2l+1} \frac{(l+|m|)!}{(l-|m|)!} \quad \text{Orthogonality \& normalization (S.-L.)}$$

$$P_0(x) = 1 \quad (m=0) \quad |m|=1: \quad l=1, 2, 3, \dots$$

$$P_1(x) = x = c_\theta \quad P_1'(x) = 1 \quad s_\theta$$

$$|m|=2: \quad l=2, 3, \dots$$

$$|m|=3: \quad l=3, 4, \dots$$

$$\begin{aligned}
 P_0(x) &= 1 \quad (m=0) & |m|=1: \quad l=1, 2, 3, \dots & |m|=2: \quad l=2, 3, \dots & |m|=3: \quad l=3, 4, \dots \\
 P_1(x) &= x = C_0 & P_1^1(x) &= (-) S_0 \\
 P_2(x) &= \frac{1}{2}(3C_0^2 - 1) & P_2^1(x) &= (-) 3 S_0 C_0 & P_2^2(x) = 3 S_0^2 \\
 P_3(x) &= \frac{1}{2}(5C_0^3 - 3C_0) & P_3^1(x) &= (+) \frac{3}{2} S_0 (5C_0^2 - 1) & P_3^2(x) = 15 S_0^2 C_0 & P_3^3(x) = (-) 15 S_0^3
 \end{aligned}$$

Group activity: 1) use recursion formulas or derivatives to calculate  $P_l^m$   
 2) verify that  $L^2 P_l^m(C_0) = h^2 l(l+1) P_l^{l+m}(C_0)$

\* Spherical harmonics:  $Y_{lm}(\theta, \phi) \equiv \sqrt{\frac{2l+1}{4\pi}} \frac{(l-m)!}{(l+m)!} P_l^{l+m}(C_0) e^{im\phi}$ ,  $\varepsilon = \begin{cases} (-1)^m & m > 0 \\ 1 & m < 0 \end{cases}$

They are eigenfunctions of both  $L^2, L_z$ :  $L^2 Y_{lm} = h^2 l(l+1) Y_{lm}$   $L_z Y_{lm} = hm Y_{lm}$

They are normalized:  $\langle lm' | lm \rangle = \int d\Omega Y_{lm'}^*(\theta, \phi) Y_{lm}(\theta, \phi) = \delta_{ll'} \delta_{mm'}$

They are complete:  $\sum_{lm} |lm\rangle \langle lm| = I$   $\sum_{lm} Y_{lm}(\theta, \phi) Y_{lm}^*(\theta', \phi') = \delta(C_0 - C_0') \delta(\phi - \phi')$

\* Application: solution of  $\nabla^2 V = 0$  in spherical coordinates: solid harmonics

$$\left[ \frac{1}{r} \frac{\partial^2}{\partial r^2} r - \frac{l(l+1)}{r^2} \right] R(r) = 0 \quad \text{let } r R(r) = u(r)$$

$$u'' = \frac{l(l+1)}{r^2} u = \frac{(-l-1)(-l)}{r^2} u \quad u = \underbrace{r^{-l-1}}_{\text{nonsingular at } r \rightarrow 0} \text{ or } r^{-l}$$

$$V(r, \theta, \phi) = \sum_{lm} r^l P_l^{l+m}(C_0) e^{im\phi}$$

$$= \sum_{lm} r^{l-l-m!} (a C_0^{l-m!} + b C_0^{l-m!-2} + \dots) r^{l+m!} S_0^{l+m!} (C_\phi + i S_\phi)^m$$

$$= \sum_{lm} (a z^{l-m!} + b z^{l-m!-2} r^2 + \dots) (x + iy)^m$$

$$= \sum_{\substack{i+j=k \\ i,j,k \geq 2}} a_{ijk} x^i y^j z^k \quad \text{multinomial in } x, y, z! \quad (\text{but not all terms})$$

$$\begin{aligned}
 \nabla^2 V &= \sum_{ijk} a_{ijk} (\delta_x^2 + \delta_y^2 + \delta_z^2) x^i y^j z^k \\
 &= \sum_{ijk} a_{ijk} [i(i-1) x^{i-2} y^j z^k + \dots]
 \end{aligned}$$

$$\begin{aligned}
 &= \sum_{ijk} \underbrace{\left( + i+2 \cdot i+1 \cdot a_{ijk} + j+2 \cdot j+1 \cdot a_{ij2k} + k+2 \cdot k+1 \cdot a_{ijk+2} \right)}_{=0 \quad \forall i,j,k} x^i y^j z^k = 0
 \end{aligned}$$

Example:  $a_{200} + a_{020} + a_{002} = 0$  excludes  $l=0$ .  $r^2 = x^2 + y^2 + z^2$ , but not

$$l=2: [ m=2: x^2-y^2, 2xy, \quad m=1: 2xz, 2yz, \quad m=0: \frac{1}{2}(3z^2-r^2) ]$$

\* Atomic orbitals:

	$ m =0$	$ m =1$	$ m =2$	$ m =3$
$l=0$ (s)sharp	$S = 1$ 			
$l=1$ (p)principal	$P_z = z$ 	$P_x = x$ 	$P_y = y$ 	
$l=2$ (d)iffuse	$d_{2z}r^2 = z^2 - \frac{1}{2}(x^2+y^2)$ 	$d_{xz} = 2xz$ 	$d_{yz} = 2yz$ 	$d_{xy} = 2xy$ $d_{x^2-y^2} = x^2-y^2$ 
$l=3$ (f)ine	$f_{5z}3zr^2 = z^3 - \frac{3}{2}(x^2y^2)z$ 	$f_{5z^2}x^2r^2 = 6z^2x - \frac{3}{2}(x^3+y^3)$ $f_{5z^2}y^2r^2 = 6z^2y - \frac{3}{2}(y^3+x^3)$ 	$f_{xyz} = 30xyz$ $f_{x^2-y^2z} = 15(x^2z-y^2z)$ 	$f_{x^2z^2} = 15(x^2z^2-y^2z^2)$ $f_{y^2z^2} = 15(y^2z^2-x^2z^2)$ 