

* family of Gaussian integrals: $I_n \equiv \int_0^\infty dx x^n e^{-\alpha x^2}$

$$I_1 = \int_0^\infty x e^{-\alpha x^2} dx = \frac{1}{2\alpha} \int_0^\infty e^{-u} du = \frac{1}{2\alpha} \quad \begin{matrix} u = 2x^2 \\ du = 2 \cdot 2x dx \end{matrix}$$

$$I_0^2 = \int_0^\infty \int_0^\infty dx dy e^{-\alpha(x^2+y^2)} = \int_0^\infty s ds e^{-\alpha s^2} \int_0^{2\pi} d\phi = \frac{\pi}{2} I_1 = \frac{\pi}{4\alpha}$$

$$I_{n+2} = \int_0^\infty x^{n+2} dx e^{-\alpha x^2} = -\frac{\partial}{\partial \alpha} I_n = \frac{n+1}{2\alpha} I_n \quad I_2 = -\frac{\partial}{\partial \alpha} \frac{1}{2\alpha} = \frac{1}{\alpha^2}$$

n	0	1	2	3	4	5	6
$\frac{1}{2}(n+1)$	$1/2$	1	$3/2$	2	$5/2$	3	$7/2$
$2I_n$	$\sqrt{\frac{\pi}{\alpha}}$	$\frac{1}{\alpha}$	$\frac{1}{2}\sqrt{\frac{\pi}{\alpha^3}}$	$\frac{1}{\alpha^2}$	$\frac{3}{4}\sqrt{\frac{\pi}{\alpha^5}}$	$\frac{2}{\alpha^3}$	$\frac{15}{8}\sqrt{\frac{\pi}{\alpha^7}}$

$$I_{2k} = \left(\frac{2k-1}{2\alpha}\right) \left(\frac{2k-3}{2\alpha}\right) \cdots \left(\frac{1}{2\alpha}\right) I_0 = \frac{(2k-1)!!}{(2\alpha)^k} \frac{1}{2} \sqrt{\frac{\pi}{\alpha}} = \frac{(2k)!}{k! (2\alpha)^{k+1} \sqrt{\pi}} = \frac{\Gamma(k+\frac{1}{2})}{2 \alpha^{k+\frac{1}{2}}}$$

$$I_{2k+1} = \left(\frac{2k}{2\alpha}\right) \left(\frac{2k-2}{2\alpha}\right) \cdots \left(\frac{2}{2\alpha}\right) I_1 = \frac{k!}{\alpha^k} \frac{1}{2\alpha} = \frac{\Gamma(k+1)}{2 \alpha^{k+1}} \xrightarrow{n=2k+1} I_n = \frac{\Gamma(\frac{n+1}{2})}{2 \alpha^{\frac{n+1}{2}}}$$

$$\Gamma(t) \equiv \int_0^\infty u^{t-1} e^{-u} du \quad \text{let } u = 2x^2 \quad du = 2 \cdot 2x dx$$

$$= \int_0^\infty (2x^2)^{t-1} e^{-2x^2} \cdot 2 \cdot 2x dx$$

$$= 2\alpha^t \int_0^\infty x^{2t-1} e^{-\alpha x^2} dx$$

$$= 2\alpha^t I_{2t-1} \quad \begin{matrix} n=2t-1 \\ t = \frac{1}{2}(n+1) \end{matrix}$$

$$\text{so } \boxed{2I_n = \frac{\Gamma(\frac{1}{2}(n+1))}{2^{\frac{1}{2}(n+1)}}}$$

$$\Gamma(t+1) = \int_0^\infty u^t e^{-u} du$$

$$= \int_0^\infty u^t d(e^{-u})$$

$$= -u^t e^{-u} \Big|_0^\infty + \int_0^\infty e^{-u} du^t$$

$$= t \int_0^\infty u^{t-1} e^{-u} du$$

$$\left\{ \begin{array}{l} \Gamma(t+1) = t \Gamma(t) \\ \Gamma(n+1) = n! \\ \Gamma(\frac{1}{2}) = \sqrt{\pi} \end{array} \right\} \text{properties of } \Gamma(t)$$

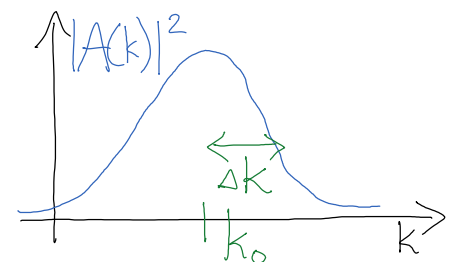
* Moments of the Gaussian wave packet (following Gasiorowicz 2-6)

$$A(k) = N e^{-\frac{1}{4} \left(\frac{k-k_0}{\Delta k}\right)^2} \quad |A(k)|^2 \text{ normal distribution}$$

$$\text{Normalization: } \int |\Psi(x)|^2 dx = \int |A(k)|^2 dk = 1$$

$$\int_{-\infty}^{\infty} N^2 e^{-\frac{1}{2} \left(\frac{k-k_0}{\Delta k}\right)^2} dk \quad x = \frac{k-k_0}{\Delta k} \quad dk = \Delta k dx$$

$$= N^2 \Delta k \int_{-\infty}^{\infty} e^{-x^2/2} dx = N^2 \Delta k \cdot \sqrt{2\pi} = 1 \quad \text{so } N = \frac{1}{\sqrt{2\pi \Delta k}}$$



$$= N^2 \Delta k \int_{-\infty}^{\infty} e^{-x^2/2} dx = N^2 \Delta k \cdot \sqrt{2\pi} = 1 \quad \text{so } N = \frac{1}{\sqrt{2\pi\Delta k}}$$

$$\text{Mean: } \langle k \rangle \equiv \int_{-\infty}^{\infty} N^2 e^{-\frac{1}{2}(\frac{k-k_0}{\Delta k})^2} k dk = N^2 \Delta k \int_{-\infty}^{\infty} e^{-x^2/2} (\Delta k x + k_0) dx = k_0$$

$$\begin{aligned} \text{RMS: } \langle k^2 \rangle &\equiv \int_{-\infty}^{\infty} N^2 e^{-\frac{1}{2}(\frac{k-k_0}{\Delta k})^2} k^2 dk = N^2 \Delta k \int_{-\infty}^{\infty} e^{-x^2/2} (\Delta k x + k_0)^2 dx \\ &= N^2 \Delta k^3 \int_{-\infty}^{\infty} e^{-x^2/2} x^2 dx + k_0^2 = \Delta k^2 + \langle k \rangle^2 \quad \text{so } \sigma_k = \Delta k \end{aligned}$$

* Fourier Transform: to obtain the wave function $\Psi(x)$

$$\text{Note the factor } \frac{1}{\sqrt{2\pi}} \text{ ensures that } \int_{-\infty}^{\infty} |\Psi(x)|^2 dx = \int_{-\infty}^{\infty} |A(k)|^2 dk$$

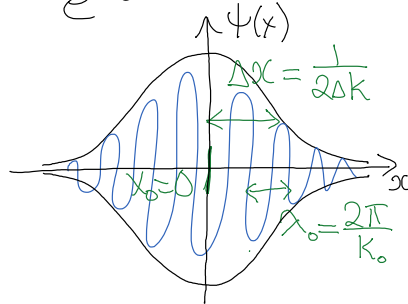
$$\begin{aligned} \Psi(x) &\equiv \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} A(k) e^{ikx} dk = \frac{N}{\sqrt{2\pi}} \int_{-\infty}^{\infty} e^{-\frac{1}{4}x^2 + i(\Delta k x + k_0)x} dk \\ &= \frac{N}{\sqrt{2\pi}} \Delta k \int_{-\infty}^{\infty} du e^{-\frac{1}{4}u^2} e^{-\Delta k^2 x^2} e^{ik_0 x} \end{aligned}$$

$A(k) = N e^{-\frac{1}{4}x^2}$
 $k = \Delta k \cdot x + k_0$
 • complete the square:
 $-\frac{1}{4}x^2 + i\Delta k x = -\frac{1}{4}u^2 - \Delta k^2 x^2$
 where $u = x - 2i\Delta k x$
 $dk = \Delta k dx = \Delta k du$

$$\Psi(x) = \underbrace{\frac{1}{\sqrt{2\pi\Delta k}}}_{\text{norm}} \underbrace{e^{-\frac{1}{4}(\frac{x}{\Delta k})^2}}_{\text{envelope}} \underbrace{e^{ik_0 x}}_{\text{phase}}$$

compare:

$$A(k) = \underbrace{\frac{1}{\sqrt{2\pi\Delta k}}}_{\text{norm}} \underbrace{e^{-\frac{1}{4}(\frac{k-k_0}{\Delta k})^2}}_{\text{envelope}}$$



Question: How does $\Psi(x)$ change if $A(k) \rightarrow A(k)e^{-ikx_0}$?

Note: We use the inverse Fourier transform for $\Psi(x) \rightarrow A(k)$
 The Gaussian wave packet is symmetric in x vs. k -space!

* Heisenberg uncertainty principle: the more frequency components of a wave packet, i.e. the larger the bandwidth Δk , the easier it is to localize $\Psi(x)$, i.e. the smaller Δx .

The Gaussian is the "minimum uncertainty wave packet" for which $\Delta x \Delta k = \frac{1}{2}$. Any other shape has larger $\Delta k \Delta x$.

$$\text{Thus H.U.P. } \Delta x \Delta k \geq \frac{1}{2} \quad \text{or } \boxed{\Delta x \Delta p \geq \hbar/2}$$