

* Example: Harmonic Oscillator Radiation

(charge q attached to spring $k=m\omega^2$ in x -axis)

$$\vec{p} = q \langle n | \hat{x} | n' \rangle \hat{x} = q \sqrt{\frac{\hbar}{2m\omega}} (a_- + a_+) = q \sqrt{\frac{\hbar}{2m\omega}} (\underbrace{\sqrt{n'} \delta_{n,n'-1}}_{\text{emission}} + \underbrace{\sqrt{n} \delta_{n',n-1}}_{\text{absorption}})$$

using $a_{\pm} = \frac{1}{\sqrt{2\hbar m\omega}} (m\omega \hat{x} \mp i\hat{p})$; $a_+ |n\rangle = \sqrt{n+1} |n+1\rangle$ $a_- |n\rangle = \sqrt{n} |n-1\rangle$

The selection rule is: $\Delta n = 1$ states separated by one photon!Radiation frequency: $\hbar\omega_0 = \Delta E = \hbar\omega \Delta n' = \hbar\omega$ Classical!

$$A = \frac{\omega^3 |\vec{p}|^2}{3\pi\epsilon_0 \hbar c^3} = \frac{\omega^3}{3\pi\epsilon_0 \hbar c^3} \cdot \frac{q^2 \hbar n}{2m\omega} = \frac{n q^2 \omega^2}{6\pi\epsilon_0 m c^3} = \mathcal{T} = \frac{1}{\tau_n}$$

$$P = A \cdot \hbar\omega = \frac{q^2 \omega^2}{6\pi\epsilon_0 m c^3} \left(E - \frac{1}{2} \hbar\omega \right) \text{ vs. } \frac{q^2 \ddot{x}^2}{6\pi\epsilon_0 c^3} = \frac{q^2 \omega^2}{12\pi\epsilon_0 c^3} E = \frac{q^2 \omega^2}{6\pi\epsilon_0 m c^3} E$$

quantum diff.

Satisfies the correspondence principle as expected.

* Selection Rules: when does $\vec{p} = \langle \psi_b | q \vec{r} | \psi_a \rangle = 0$?for ∇^2 symmetric, $\psi = f(r) Y_{lm}(\theta, \phi) \rightarrow \vec{p} = q \langle n'l'm' | \vec{r} | nlm \rangle$ + rules for m, m' : $[L_z, x] = i\hbar y$, $[L_z, y] = -i\hbar x$, $[L_z, z] = 0$

$$a) 0 = \langle n'l'm' | [L_z, z] | nlm \rangle = \langle n'l'm' | L_z z - z L_z | nlm \rangle$$

$$= \hbar \langle n'l'm' | m' z - z m | nlm \rangle = \hbar (m - m') \langle n'l'm' | z | nlm \rangle$$

thus if $m \neq m'$ then $p_z = 0$

$$b) [L_z, X_{\pm}] = \pm \hbar X_{\pm} \quad X_{\pm} = X \pm iy$$

$$\langle n'l'm' | [L_z, x_{\pm}] | nlm \rangle = \langle n'l'm' | L_z x_{\pm} - x_{\pm} L_z | nlm \rangle$$

$$\langle n'l'm' | \pm \hbar x_{\pm} | nlm \rangle = \pm \hbar (m' - m) \langle n'l'm' | x | nlm \rangle$$

thus $p_x \pm i p_y = 0$ unless $m' = m \pm 1$

selection rules for m : $\Delta m = \pm 1, 0$

+ rules for l : $[L^2, [L^2, \vec{r}]] = 2\hbar^2 (\vec{r} L^2 + L^2 \vec{r})$

$$\langle n'l'm' | [L^2, [L^2, \vec{r}]] | nlm \rangle = \hbar^2 [l'(l'+1) - l(l+1)] \langle n'l'm' | [L^2, \vec{r}] | nlm \rangle$$

$$\| = \hbar^4 [l'(l'+1) - l(l+1)]^2 \langle n'l'm' | \vec{r} | nlm \rangle$$

$$2\hbar^2 \langle n'l'm' | \vec{r} L^2 + L^2 \vec{r} | nlm \rangle = 2\hbar^4 [l'(l'+1) + l(l+1)] \langle n'l'm' | \vec{r} | nlm \rangle$$

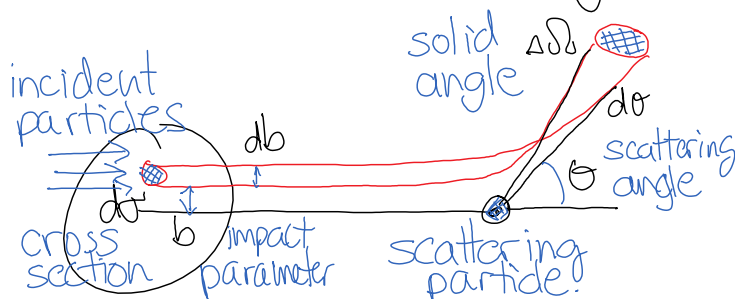
$$\hbar^4 \left\{ \underbrace{[l'(l'+1) - l(l+1)]^2}_{(l'+l+1)(l'-l)} - 2 \underbrace{[l'(l'+1) + l(l+1)]}_{(l'+l+1)^2 + (l'-l)^2 - 1} \right\} \langle n'l'm' | \vec{r} | nlm \rangle = 0$$

$a^2 b^2 - (a^2 + b^2 - 1) = 0$
 $[a^2 - 1][b^2 - 1] = 0$

selection rule for l : $\underbrace{[(l'+l+1)^2 - 1]}_{\text{can't be 0}} [(l'-l) - 1] = 0 \Rightarrow \Delta l = \pm 1$

+ rules for m_s : $\Delta m_s = 0$

* Introduction to Scattering Cross Sections



$$dN = \underbrace{\mathcal{L}}_{\text{beam}} \frac{d\sigma}{d\Omega} \cdot \underbrace{\Delta\Omega}_{\text{detector}} \quad [\text{expt.}]$$

$$\frac{d\sigma}{d\Omega} = \frac{2\pi}{2\pi} \cdot \underbrace{b}_{\text{physics}} \cdot \frac{db}{d\theta} \quad [\text{theory}]$$

"Differential cross sections $d\sigma \rightarrow \frac{d\sigma}{d\Omega}$ are the meeting point between experimentalists and theorists."