L53-Angular Momentum Eigenfunctions — skipped.

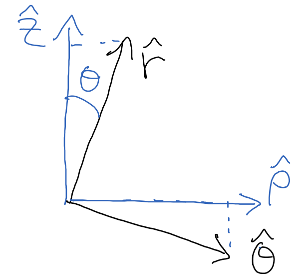
Wednesday, January 20, 2016 08:25

* co-ordinate representation of $\vec{L}, L^2$ operators

$$\vec{L} = \vec{r}\times\vec{p} = -i\hbar\,\vec{r}\times\nabla = -i\hbar(\hat{r}r)\times\left(\hat{r}\frac{\partial}{\partial r} + \hat{\theta}\frac{1}{r}\frac{\partial}{\partial\theta} + \hat{\phi}\frac{1}{r\sin\theta}\frac{\partial}{\partial\phi}\right)$$

$$= -i\hbar\left(\hat{\phi}\frac{\partial}{\partial\theta} - \hat{\theta}\frac{1}{\sin\theta}\frac{\partial}{\partial\phi}\right)$$

$$(\hat{r}\ \hat{\theta}\ \hat{\phi}) = (\hat{x}\ \hat{y}\ \hat{z})\begin{pmatrix} s_\theta c_\phi & c_\theta c_\phi & -s_\phi \\ s_\theta s_\phi & c_\theta s_\phi & c_\phi \\ c_\theta & -s_\theta & 0 \end{pmatrix}$$




$$\vec{L} = -i\hbar\left((-\hat{x}s_\phi + \hat{y}c_\phi)\frac{\partial}{\partial\theta} - (\hat{x}c_\theta c_\phi + \hat{y}c_\theta s_\phi - \hat{z}s_\theta)\frac{1}{s_\theta}\frac{\partial}{\partial\phi}\right)$$

$$L_x = -i\hbar\left(-s_\phi\frac{\partial}{\partial\theta} - c_\phi\cot\theta\frac{\partial}{\partial\phi}\right)$$

$$L_y = -i\hbar\left(c_\phi\frac{\partial}{\partial\theta} - s_\phi\cot\theta\frac{\partial}{\partial\phi}\right)$$

$$L_z = -i\hbar\frac{\partial}{\partial\phi}$$

note: $\frac{\partial}{\partial\phi}$ also in $L_x, L_y$.

$$L_\pm = L_x \pm iL_y = \hbar\left(\underbrace{(is_\phi \pm c_\phi)}_{\pm(c_\phi \pm is_\phi) = \pm e^{i\phi}}\frac{\partial}{\partial\theta} + i\underbrace{(c_\phi \pm is_\phi)}_{e^{i\phi}}\cot\theta\frac{\partial}{\partial\phi}\right)$$

$$= \pm\hbar e^{\pm i\phi}\left(\frac{\partial}{\partial\theta} \pm i\cot\theta\frac{\partial}{\partial\phi}\right)$$

$$L_\pm L_\mp = \pm\hbar e^{\pm i\phi}\left(\frac{\partial}{\partial\theta} \pm i\cot\theta\frac{\partial}{\partial\phi}\right)\cdot\mp\hbar e^{\mp i\phi}\left(\frac{\partial}{\partial\theta} \mp i\cot\theta\frac{\partial}{\partial\phi}\right)$$ [Prob# 4.21]

$$= -\hbar^2\left(\frac{\partial^2}{\partial\theta^2} \mp i\left(\underbrace{\frac{\partial\cot\theta}{\partial\theta}}_{-\csc^2\theta} + \cancel{\cot\theta\frac{\partial}{\partial\theta}}\right)\frac{\partial}{\partial\phi} \pm i\cot\theta\,\underbrace{e^{\pm i\phi}\frac{\partial}{\partial\phi}e^{\mp i\phi}}_{\mp i + \cancel{\frac{\partial}{\partial\phi}}}\frac{\partial}{\partial\theta} + \cot^2\theta\cdot\underbrace{e^{\pm i\phi}\frac{\partial}{\partial\phi}e^{\mp i\phi}}_{\mp i + \frac{\partial}{\partial\phi}}\frac{\partial}{\partial\phi}\right)$$

$$= -\hbar^2\left(\frac{\partial^2}{\partial\theta^2} + \cot\theta\frac{\partial}{\partial\theta} \pm i\underbrace{(\csc^2\theta - \cot^2\theta)}_{1}\frac{\partial}{\partial\phi} + \cot^2\theta\frac{\partial^2}{\partial\phi^2}\right)$$

$$= -\hbar^2\left[\frac{1}{\sin\theta}\frac{\partial}{\partial\theta}\sin\theta\frac{\partial}{\partial\theta} + \frac{\cos^2\theta}{\sin^2\theta}\frac{\partial^2}{\partial\phi^2}\right] \pm \underbrace{\hbar\cdot -i\hbar\frac{\partial}{\partial\phi}}_{\hbar L_z = \frac{1}{2}[L_+, L_-]}$$

$$= L_x^2 + L_y^2 \pm \hbar L_z$$

$$L^2 = (L_+L_- - \hbar L_z) + L_z^2 = -\hbar^2\left[\frac{1}{\sin\theta}\frac{\partial}{\partial\theta}\sin\theta\frac{\partial}{\partial\theta} + \frac{1}{\sin^2\theta}\frac{\partial^2}{\partial\phi^2}\right]$$

* top rung eigenfunction [Prob#4.22]

$$\langle \Omega | L_+ | \ell\ell\rangle = \hbar e^{i\phi}\left(\frac{\partial}{\partial\theta} + i\cot\theta\,\frac{\partial}{\partial\phi}\right) Y_{\ell\ell}(\theta,\phi) = 0$$

let $Y_{\ell\ell} = f(\theta)\cdot e^{i\ell\phi}$ so that $L_z e^{i\ell\phi} = -i\hbar\partial_\phi e^{i\ell\phi} = \hbar\ell\, e^{i\ell\phi}$

$$\frac{df}{d\theta} - \ell\cot\theta\cdot f = 0 \qquad \frac{df}{f} = \ell\,\frac{\cos\theta\, d\theta}{\sin\theta} = \ell\,\frac{d\sin\theta}{\sin\theta}$$

$$\ln f = \ln(\sin\theta)^\ell \qquad f(\theta) = \sin^\ell\theta$$

thus $Y_{\ell\ell}(\theta,\phi) = N\sin^\ell\theta\, e^{i\ell\phi}$. Now determine N:

$$\langle\ell\ell|\ell\ell\rangle = \int d\Omega\, Y^*_{\ell\ell}(\theta,\phi)\,Y_{\ell\ell}(\theta,\phi) = |N|^2\int_0^\pi \sin\theta\, d\theta\int_0^{2\pi} d\phi\, \sin^{2\ell}\theta\, e^{i\cancelto{0}{(\ell-\ell)}\phi} = 2\pi|N|^2 I_\ell = 1$$

where $I_\ell = \int_0^\pi \sin^{2\ell+1}\theta\, d\theta = \int_0 \underbrace{\sin^{2\ell}\theta}_{u}\cdot d\underbrace{(-\cos\theta)}_{v}$

$$= \cancelto{0}{-\sin^{2\ell}\theta\cdot\cos\theta\Big|_0^\pi} + \int_0^\pi \cos\theta\, d(\sin^{2\ell}\theta)$$

$$= \int_0^\pi 2\ell\cdot\sin^{2\ell-1}\theta\cdot\underbrace{\cos^2\theta}_{1-\sin^2\theta}\,d\theta = 2\ell\,(I_{\ell-1} - I_\ell)$$

thus $I_\ell = \frac{2\ell}{2\ell+1} I_{\ell-1} = \frac{2\cdot4\cdots(2\ell)}{3\cdot5\cdots(2\ell+1)}\, I_0$ ; $I_0 = \int_0^\pi \sin\theta\, d\theta = 2$

$$|N|^2 = \frac{1}{2\pi I_\ell} = \frac{1}{4\pi}\,\frac{1\cdot3\cdot5\cdot7\cdots2\ell+1}{2\cdot4\cdot6\cdots2\ell} = \frac{1}{4\pi}\,\frac{1\cdot2\cdot3\cdot4\cdots2\ell+1}{2^2\cdot4^2\cdots(2\ell)^2} = \frac{(2\ell+1)!}{4\pi\,(2^\ell\ell!)^2}$$

$$\boxed{Y_{\ell\ell}(\theta,\phi) = \frac{(-1)^\ell}{2^\ell\ell!}\sqrt{\frac{(2\ell+1)!}{4\pi}}\;\underbrace{\sin^\ell\theta\, e^{i\ell\phi}}_{\sin\theta(\cos\phi+i\sin\phi) = \left(\frac{x+iy}{\rho}\right)^\ell}}$$

up to an arbitrary phase $(-1)^\ell$

* lower rungs: use $L_\pm|\ell m\rangle = \hbar\sqrt{(\ell\mp m)(\ell\pm m+1)}\;|\ell, m\pm1\rangle$

- apply $L_-$ repeatedly to $Y_{\ell\ell}$ to get all $Y_{\ell m}$ for $\ell$.

$$\boxed{Y_{\ell m}(\theta,\phi) = (-1)^m\sqrt{\frac{2\ell+1}{4\pi}\frac{(\ell-m)!}{(\ell+m)!}}\;P_\ell^m(\cos\theta)\,e^{im\phi}}$$

$$L_\pm[Y_{\ell m}] = \pm\hbar e^{\pm i\phi}\left(\frac{\partial}{\partial\theta} \pm i\cot\theta\,\frac{\partial}{\partial\phi}\right)\left[(-1)^m\sqrt{\frac{2\ell+1}{4\pi}\frac{(\ell-m)!}{(\ell+m)!}}\;P_\ell^m e^{im\phi}\right]$$

$$= \hbar\sqrt{(l\mp m)(l\pm m+1)} \cdot \left[(-1)^{m\pm1}\sqrt{\frac{2l+1}{4\pi}\frac{(l-m\mp1)!}{(l+m\pm1)!}}\; P_l^{m\pm1}\, e^{i(m\pm1)\phi}\right]$$

- use this equation to construct recurrence relations for $P_l^m$

$$\left[\frac{d}{d\xi}(1-\xi^2)\frac{d}{d\xi} - \frac{m^2}{1-\xi^2} + l(l+1)\right] P_l^m(\xi) = 0 \qquad \text{Sturm-Liouville ODE}$$

$$P_l^m(\xi) = \underbrace{(1-\xi^2)^{m/2}}_{\sin^m\theta}\left(\frac{d}{d\xi}\right)^m P_l(\underbrace{\xi}_{\cos\theta}) \qquad \text{where} \qquad P_l(\xi) = \frac{1}{2^l l!}\left(\frac{d}{d\xi}\right)^l(\xi^2-1)^l$$

$$\frac{dP}{d\xi} = \frac{d}{d\xi}\left[(1-\xi^2)^{m/2}\left(\frac{d}{d\xi}\right)^m P_l(\xi)\right]$$

$$= \frac{m}{2}(1-\xi^2)^{m/2-1}(-2\xi)\frac{d}{d\xi}^m P_l(\xi) + (1-\xi^2)^{m/2}\left(\frac{d}{d\xi}\right)^{m+1}P_l(\xi)$$

$$= -\frac{1}{\sqrt{1-\xi^2}}P_l^{m+1} - \frac{m\xi}{1-\xi^2}P_l^m$$

- use $L_\pm = L_x \pm iL_y$ to calculate $L_x|lm\rangle$ or $L_y|lm\rangle$

* summary: $H, L^2, L_z$ are a C.S.C.O (Complete Set of Commuting Observables)

which completely characterize the eigenfunctions

of Hamiltonians with rotational symmetry

(eg. central potentials like H-atom, 3-d oscillator, ...)

$$\mathcal{H} = \frac{p^2}{2m}+V = -\frac{\hbar^2}{2m}\nabla^2+V = -\hbar^2\frac{1}{r}\frac{\partial^2}{\partial r^2}r + \underbrace{\frac{L^2}{2I}}_{-\frac{\hbar^2 l(l+1)}{2mr^2}} + V(r)$$

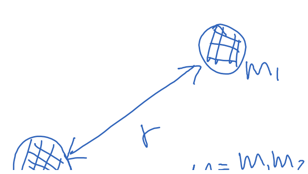
"centrifugal potential"

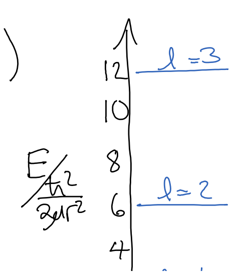
Energy is degenerate in "m" (rotational sym.)

$$\psi(r,\theta,\phi) = R(r)\,Y_{lm}(\theta,\phi) \qquad \begin{aligned} L^2 Y_{lm} &= \hbar^2 l(l+1)\,Y_{lm} \\ L_z Y_{lm} &= \hbar m\; Y_{lm}\end{aligned}$$

* Example: diatomic molecule: (rigid rotor)

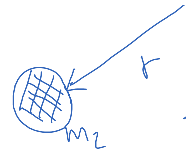
if $r$ is fixed then




$$\mathcal{H} = \frac{L^2}{\ldots} = \frac{\hbar^2 l(l+1)}{\ldots}$$

$$\mathcal{H} = \frac{L^2}{2I} = \frac{\hbar^2 \ell(\ell+1)}{2\mu r^2}$$




$$\mu = \frac{m_1 m_2}{m_1 + m_2}$$