

## \* Example: Harmonic Oscillator Radiation

(charge  $q$  attached to spring  $k=m\omega^2$  in  $x$ -axis)

$$\vec{p} = q \langle n | \hat{x} | n' \rangle \hat{x} = q \sqrt{\frac{\hbar}{2m\omega}} (a_- + a_+) = q \sqrt{\frac{\hbar}{2m\omega}} (\underbrace{\sqrt{n'} \delta_{n,n'-1}}_{\text{emission}} + \underbrace{\sqrt{n} \delta_{n',n-1}}_{\text{absorption}})$$

$$\text{using } a_{\pm} = \frac{1}{\sqrt{2\hbar m\omega}} (m\omega \hat{x} \mp i\hat{p}) ; \quad a_+ |n\rangle = \sqrt{n+1} |n+1\rangle \quad a_- |n\rangle = \sqrt{n} |n-1\rangle$$

The selection rule is:  $\Delta n = 1$  states separated by one photon!Radiation frequency:  $\hbar\omega_0 = \Delta E = \hbar\omega \Delta n' = \hbar\omega$  Classical!

$$A = \frac{\omega^3 |\vec{p}|^2}{3\pi\epsilon_0 \hbar c^3} = \frac{\omega^3}{3\pi\epsilon_0 \hbar c^3} \cdot \frac{q^2 \hbar n}{2m\omega} = \frac{n q^2 \omega^2}{6\pi\epsilon_0 m c^3} = \mathcal{T} = \frac{1}{\tau_n}$$

$$P = A \cdot \hbar\omega = \frac{q^2 \omega^2}{6\pi\epsilon_0 m c^3} \left( E - \frac{1}{2} \hbar\omega \right) \quad \text{vs.} \quad \frac{q^2 \ddot{x}^2}{6\pi\epsilon_0 c^3} = \frac{q^2 \dot{x}^2 \omega^4}{12\pi\epsilon_0 c^3} = \frac{q^2 \omega^2}{6\pi\epsilon_0 m c^3} E$$

*quantum diff.*

Satisfies the correspondence principle as expected.

\* Selection Rules: when does  $\vec{p} = \langle \psi_b | q \vec{r} | \psi_a \rangle = 0$ ?for  $\nabla^2$  symmetric,  $\psi = f(r) Y_{lm}(\theta, \phi) \rightarrow \vec{p} = q \langle n'l'm' | \vec{r} | nlm \rangle$ + rules for  $m, m'$ :  $[L_z, x] = i\hbar y$ ,  $[L_z, y] = -i\hbar x$ ,  $[L_z, z] = 0$ 

$$a) \quad 0 = \langle n'l'm' | [L_z, z] | nlm \rangle = \langle n'l'm' | L_z z - z L_z | nlm \rangle$$

$$= \hbar \langle n'l'm' | m' z - z m | nlm \rangle = \hbar (m - m') \langle n'l'm' | z | nlm \rangle$$

$$\text{thus } p_z = 0 \quad \text{unless } m \neq m'$$

$$b) \quad [L_z, x_{\pm}] = \pm \hbar x_{\pm} \quad x_{\pm} = x \pm iy$$

$$\langle n'l'm' | [L_z, x_{\pm}] | nlm \rangle = \langle n'l'm' | L_z x_{\pm} - x_{\pm} L_z | nlm \rangle$$

$$\langle n'l'm' | \pm \hbar x_{\pm} | nlm \rangle = \pm \hbar (m' - m) \langle n'l'm' | x_{\pm} | nlm \rangle$$

thus  $p_x \pm i p_y = 0$  unless  $m' = m \pm 1$

selection rules for  $m$ :  $\Delta m = \pm 1, 0$

+ rules for  $l$ :  $[L^2, [L^2, \vec{r}]] = 2\hbar^2 (\vec{r} L^2 + L^2 \vec{r})$

$$\langle n'l'm' | [L^2, [L^2, \vec{r}]] | nlm \rangle = \hbar^2 [l'(l'+1) - l(l+1)] \langle n'l'm' | [L^2, \vec{r}] | nlm \rangle$$

$$\| = \hbar^4 [l'(l'+1) - l(l+1)]^2 \langle n'l'm' | \vec{r} | nlm \rangle$$

$$2\hbar^2 \langle n'l'm' | \vec{r} L^2 + L^2 \vec{r} | nlm \rangle = 2\hbar^4 [l'(l'+1) + l(l+1)] \langle n'l'm' | \vec{r} | nlm \rangle$$

$$\hbar^4 \left\{ \underbrace{[l'(l'+1) - l(l+1)]^2}_{(l'+l+1)(l'-l)} - 2 \underbrace{[l'(l'+1) + l(l+1)]}_{(l'+l+1)^2 + (l'-l)^2 - 1} \right\} \langle n'l'm' | \vec{r} | nlm \rangle = 0$$

$a^2 b^2 - (a^2 + b^2 - 1) = 0$   
 $[a^2 - 1][b^2 - 1] = 0$

selection rule for  $l$ :  $\underbrace{[(l'+l+1)^2 - 1]}_{\text{can't be 0}} [(l'-l) - 1] = 0 \Rightarrow \Delta l = \pm 1$

+ rules for  $m_s$ :  $\Delta m_s = 0$