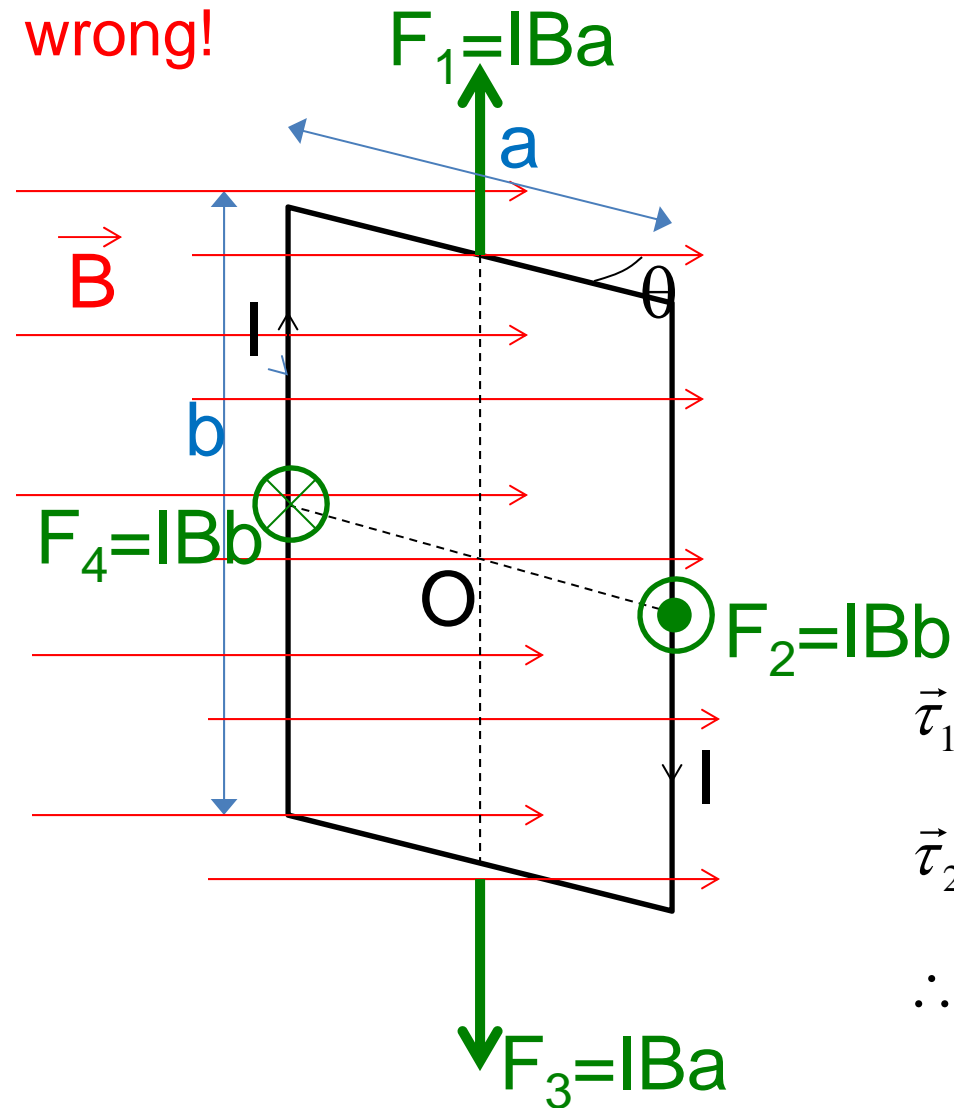


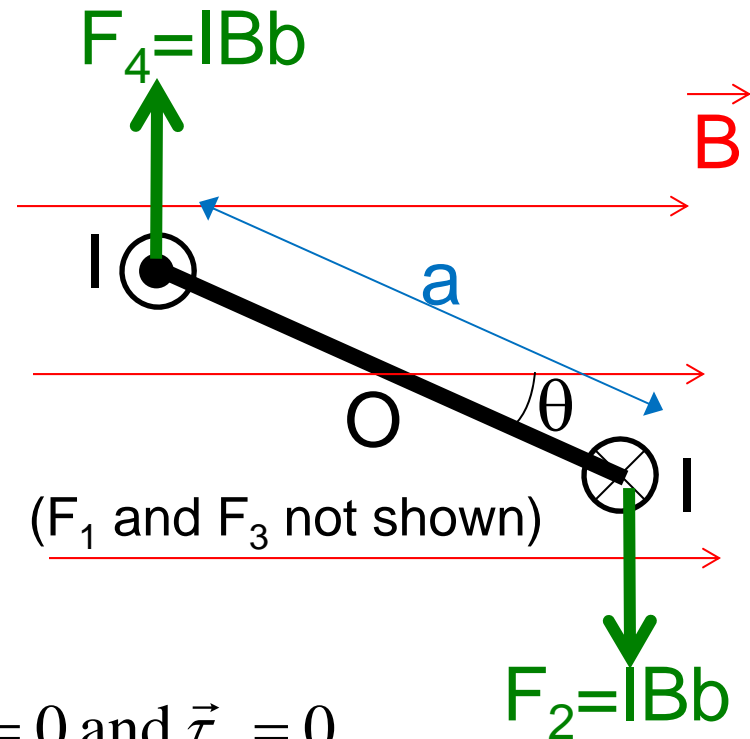
Class 28 Biot – Savart Law and Ampere's Law

The way we define θ here does not follow convention, but nothing wrong!

Magnetic torque on a Rectangular Loop



Top view



$$\vec{\tau}_1 = 0 \text{ and } \vec{\tau}_3 = 0$$

$$\vec{\tau}_2 = \vec{\tau}_4 = IBb \cdot \frac{a}{2} \cos \theta \otimes$$

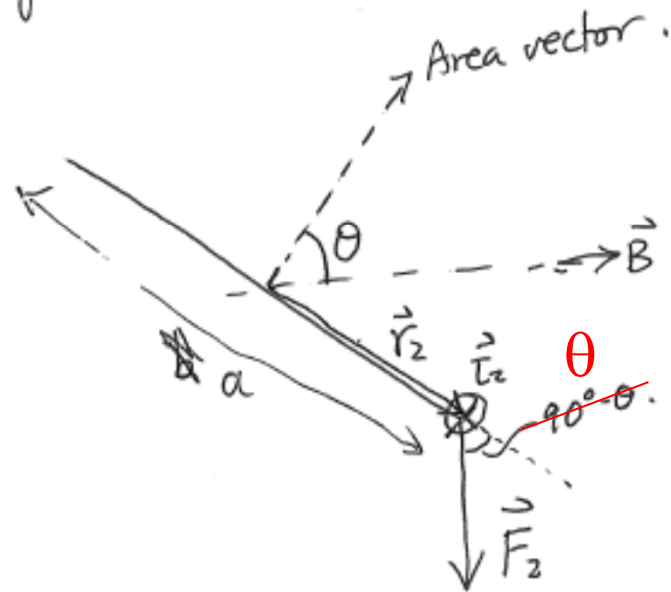
$$\begin{aligned} \therefore \vec{\tau} &= IBa \cdot b \cos \theta \otimes = IBA \cos \theta \otimes \\ &= I\vec{A} \times \vec{B} \end{aligned}$$

Correction

Class 25

Written notes pg.2

$$\text{Torque } \vec{\tau} = \vec{r} \times \vec{F}$$



$$|\vec{\tau}_2| = \frac{a}{2} \cdot IBb \cdot \sin(90^\circ - \theta).$$

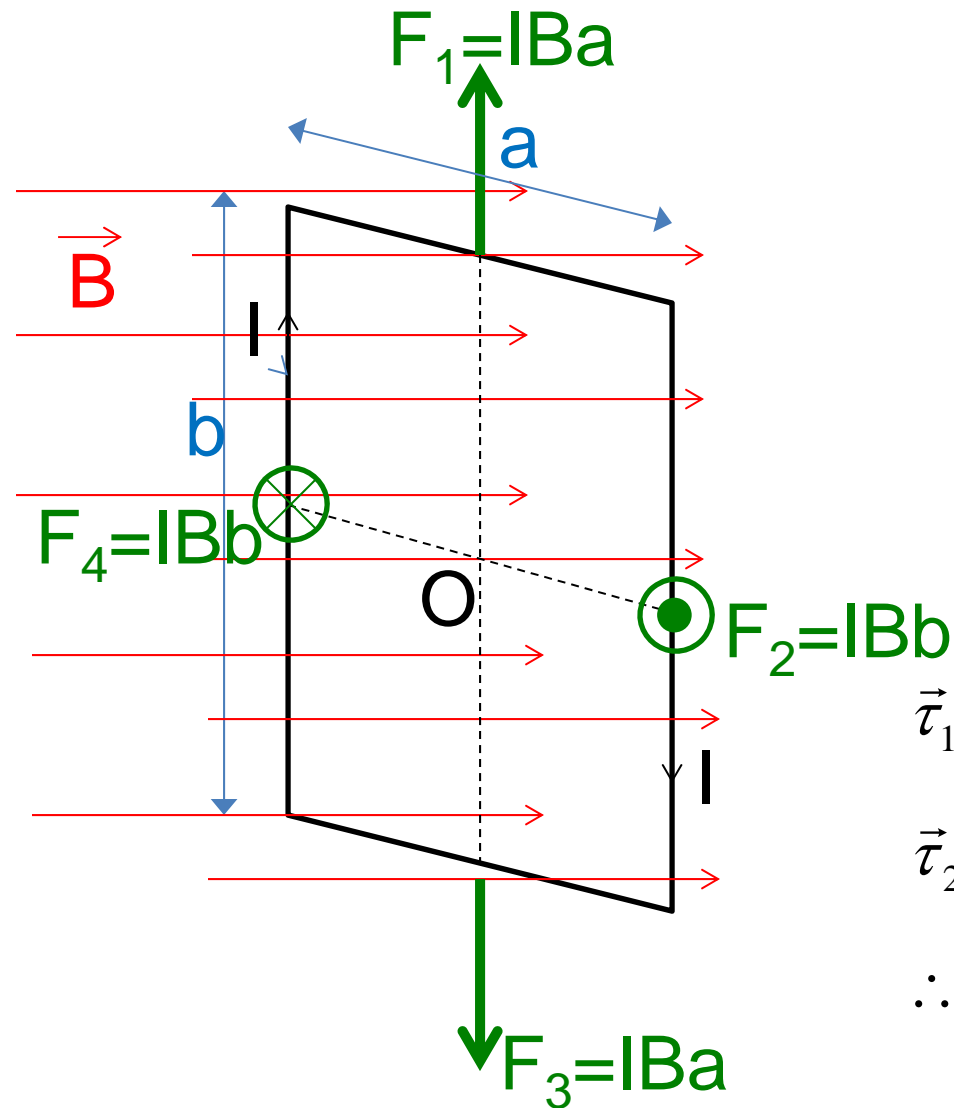
$$= \frac{a}{2} IBb \cos \theta.$$

$$= \frac{A}{2} \cdot IB \cos \theta.$$

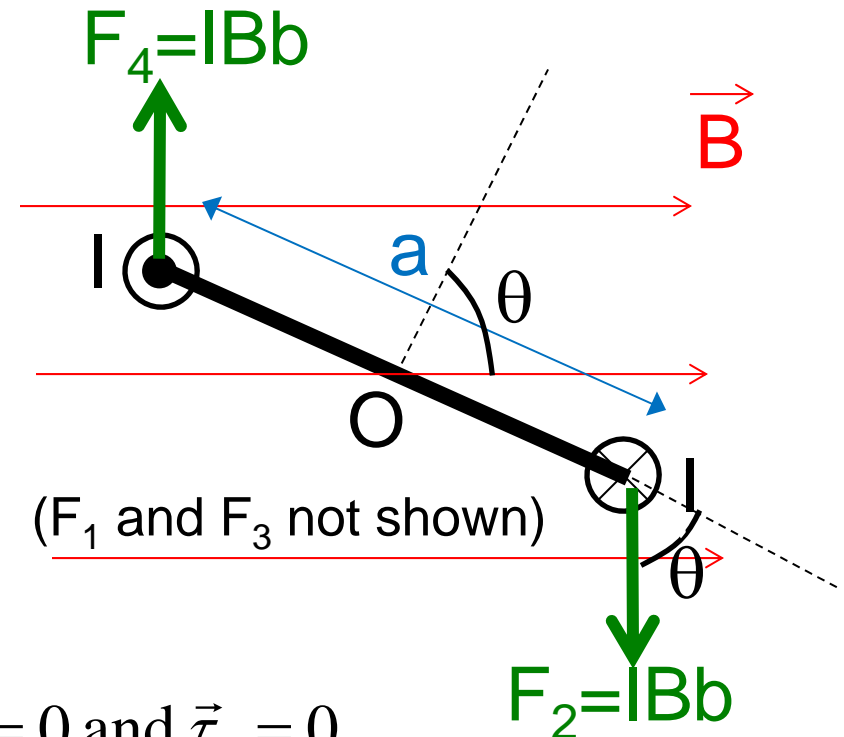
Correction Magnetic torque on a Rectangular Loop

Class 25 slide 4

Class 27 slide 3



Top view



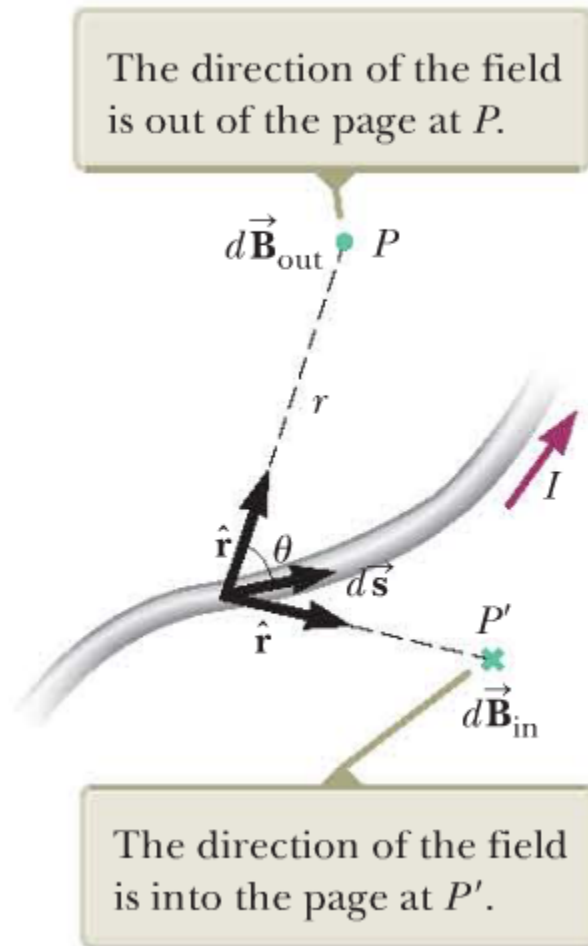
$$\vec{\tau}_1 = 0 \text{ and } \vec{\tau}_3 = 0$$

$$\vec{\tau}_2 = \vec{\tau}_4 = IBb \cdot \frac{a \sin \theta}{2 \cos \theta} \otimes$$

$$\therefore \vec{\tau} = IBa \cdot b \frac{\sin \theta}{\cos \theta} \otimes = IBa \sin \theta \otimes$$

$$= I \vec{A} \times \vec{B}$$

Biot-Savart Law



Magnetic field at point P due to the infinitesimal element ds :

$$d\vec{B} = \frac{\mu_0 I}{4\pi} \frac{d\vec{s} \times \hat{r}}{r^2} \quad \text{or} \quad \frac{\mu_0 I}{4\pi} \frac{d\vec{s} \times \vec{r}}{r^3}$$

Magnetic field due to the whole wire:

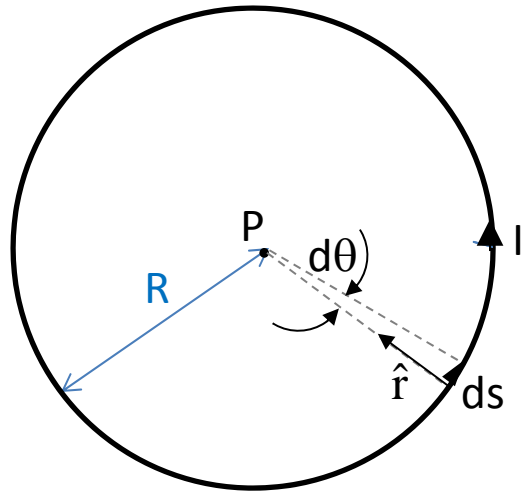
$$\vec{B} = \int_{\text{wire}} \frac{\mu_0 I}{4\pi} \frac{d\vec{s} \times \hat{r}}{r^2}$$

μ_0 is a constant called permeability of free space:

$$\mu_0 = 4\pi \times 10^{-7} \text{ TmA}^{-1}$$

In the calculation of magnetic field, Biot-Savart Law play the same role as the Coulomb's Law in electric field.

Magnetic Field at the Center of a Circular Current Loop



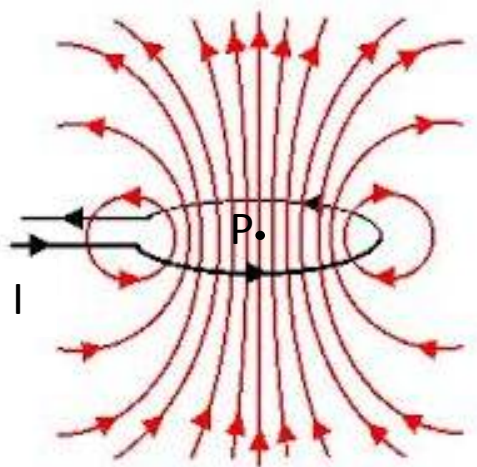
$$\begin{aligned} d\vec{B} &= \frac{\mu_0 I}{4\pi} \frac{d\vec{s} \times \hat{r}}{r^2} \\ &= \frac{\mu_0 I}{4\pi} \frac{ds}{R^2} \odot \end{aligned}$$

$$ds = R d\theta$$

$$\begin{aligned} d\vec{B} &= \frac{\mu_0 I}{4\pi} \frac{R d\theta}{R^2} \odot \\ &= \frac{\mu_0 I}{4\pi R} d\theta \odot \end{aligned}$$

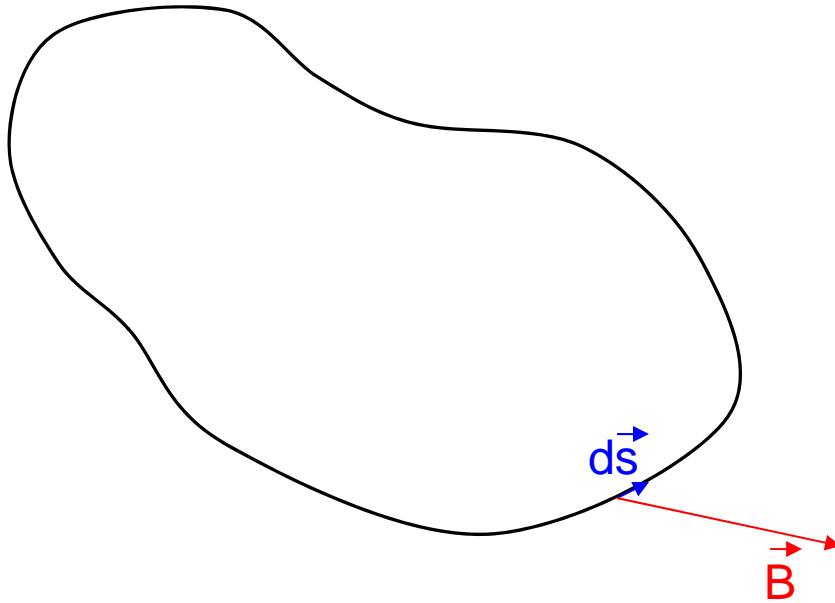
$$\begin{aligned} \therefore \vec{B} &= \int_0^{2\pi} \frac{\mu_0 I}{4\pi R} d\theta \odot = \odot \frac{\mu_0 I}{4\pi R} \int_0^{2\pi} d\theta = \frac{\mu_0 I}{4\pi R} [\theta]_0^{2\pi} \odot \\ &= \frac{\mu_0 I}{4\pi R} \cdot 2\pi \odot \end{aligned}$$

$$\therefore \vec{B} = \frac{\mu_0 I}{2R} \odot$$



Ampere's Law

Path Integral of Magnetic Field



Path Integral:

$$\oint \vec{B} \cdot d\vec{s}$$

Do you remember:

$$\oint \vec{E} \cdot d\vec{s} = ? \quad (\text{for stationary case})$$

Gauss's Law (Maxwell's first equation)

For *any* closed surface,

From Class 5

$$\epsilon_0 \Phi_E = q_{\text{in}} \quad \text{or} \quad \epsilon_0 \oint \vec{E} \cdot d\vec{A} = q_{\text{in}}$$

Two types of problems that involve Gauss's Law:

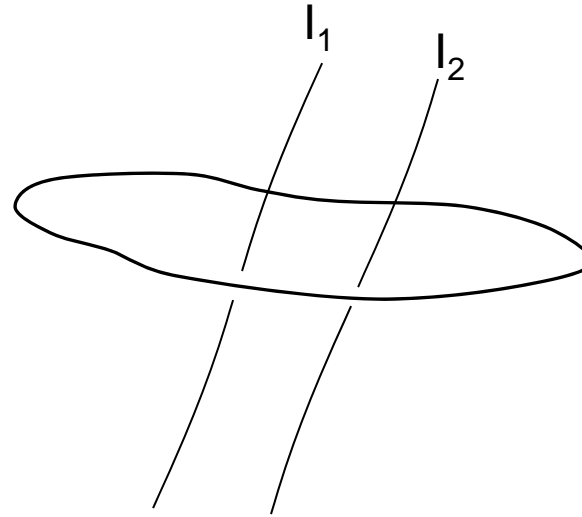
1. Give you left hand side (i.e. flux through a given surface), calculate the right hand side (i.e. charge enclosed by that surface).
2. Give you right hand side (i.e. a charge distribution) , calculate the left hand side (i.e. flux and the electric field).

Ampere's Law (Maxwell's third equation - partial)

For *any* closed loop,

$$\oint \vec{B} \cdot d\vec{s} = \mu_0 I_{\text{in}}$$

$$\mu_0 = 4\pi \times 10^{-7} \text{ TmA}^{-1}$$



Two types of problems that involve Ampere's Law:

1. Give you left hand side (i.e. line integral of a given loop), calculate the right hand side (i.e. current enclosed by that loop).
2. Give you right hand side (i.e. current) , calculate the left hand side (i.e. the line integral and the magnetic field).

Calculating Magnetic Field Using Ampere's Law

1. By symmetry argument, construct a loop so that the path integral $\oint \vec{B} \cdot d\vec{s}$ can be easily calculated. In most cases when Ampere's Law is applicable,

$$\oint \vec{B} \cdot d\vec{s} = BL$$

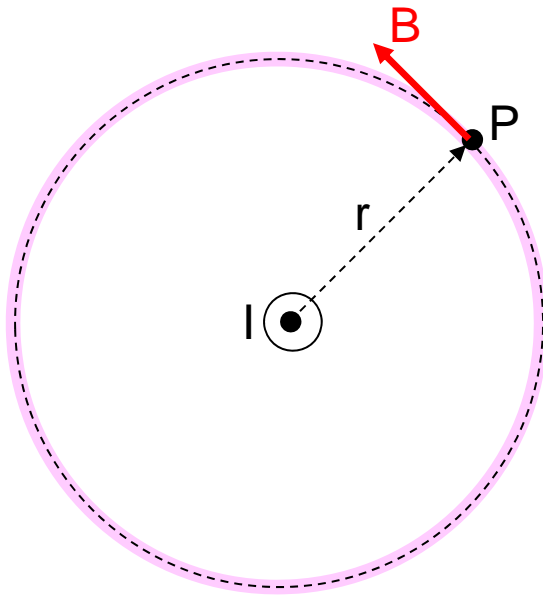
where L is the length of the loop.

2. You can then apply Ampere's Law and solve for B :

$$\oint \vec{B} \cdot d\vec{s} = \mu_0 I_{\text{in}} \Rightarrow BL = \mu_0 I_{\text{in}} \Rightarrow B = \frac{\mu_0 I_{\text{in}}}{L}$$

3. Two common cases when Ampere's Law can be used to calculate magnetic field: infinite long wire and infinite long solenoid / toroid.

Magnetic field due to a long wire



Want to calculate the magnetic field B at point P .

By symmetry argument, B is in the plane of the paper (infinite long wire), has the same magnitude for all points on the dotted circular loop (azimuthal symmetry), and tangent to the circular loop (so $\cos \theta = 1$).

$$\therefore \oint \vec{B} \cdot d\vec{s} = B \cdot 2\pi r$$

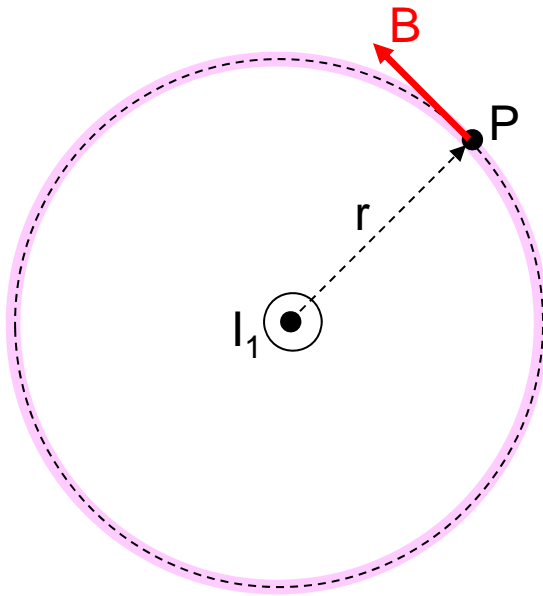
$$\text{Ampere's Law: } \oint \vec{B} \cdot d\vec{s} = \mu_0 I \Rightarrow B \cdot 2\pi r = \mu_0 I$$

$$\Rightarrow B = \frac{\mu_0 I}{2\pi r}$$

Magnetic Force Between Two Parallel Long Wires

Magnetic field at point P due to I_1 :

$$B = \frac{\mu_0 I_1}{2\pi r}$$

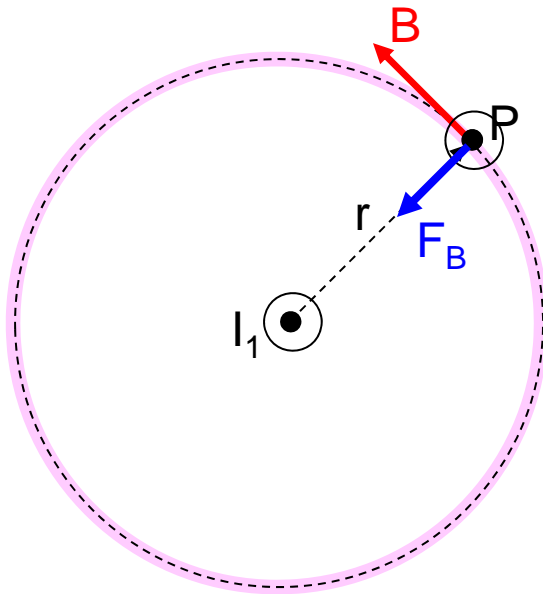


If another current I_2 parallel to I_1 is passing through point P, it will experience a force because of the field there.

Magnetic Force Between Two Parallel Long Wires

Magnetic field at point P due to I_1 :

$$B = \frac{\mu_0 I_1}{2\pi r}$$



If another current I_2 parallel to I_1 is passing through point P, it will experience a force because of the field there.

$$\vec{F}_B = I_2 \vec{L} \times \vec{B} \Rightarrow F_B = I_2 B L \sin 90^\circ = I_2 B L$$

$$\Rightarrow F_B = I_2 L \cdot \frac{\mu_0 I_1}{2\pi r}$$

$$\Rightarrow \frac{F_B}{L} = \frac{\mu_0 I_1 I_2}{2\pi r}$$

Force is attractive if the two currents are in the same direction, repulsive if the two currents are in opposite direction.