

Class 38: Electromagnetic radiation

Applying Maxwell's Third Equation to Plane Electromagnetic Waves

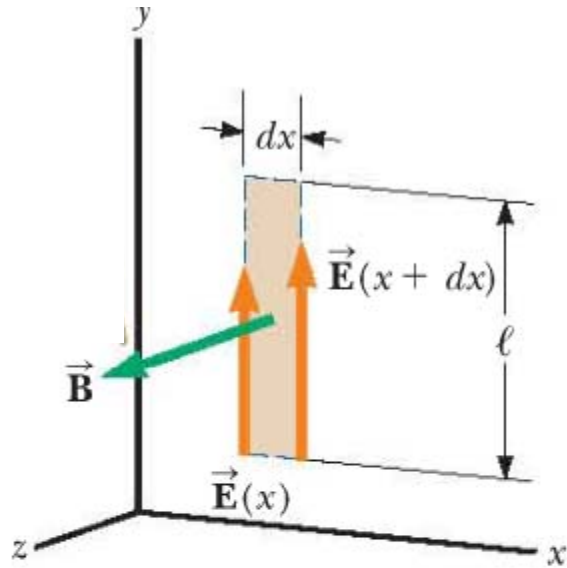


Figure 34.6 At an instant when a plane wave moving in the positive x direction passes through a rectangular path of width dx lying in the xy plane, the electric field in the y direction varies from $\vec{E}(x)$ to $\vec{E}(x + dx)$.

$$\oint \vec{E} \cdot d\vec{s} = - \frac{d}{dt} \iint \vec{B} \cdot d\vec{A}$$

$$\begin{aligned} \oint \vec{E} \cdot d\vec{s} &= E(x + dx)\ell + 0 - E(x)\ell + 0 \\ &= \ell[E(x + dx) - E(x)] \\ &= \ell \frac{\partial E}{\partial x} \cdot dx \end{aligned}$$

$$\begin{aligned} - \frac{d}{dt} \iint \vec{B} \cdot d\vec{A} &= - \frac{\partial}{\partial t} (B \cdot \ell dx) \\ &= - \ell \frac{\partial B}{\partial t} \cdot dx \end{aligned}$$

$$\therefore \frac{\partial E}{\partial x} = - \frac{\partial B}{\partial t}$$

Applying Maxwell's Fourth Equation to Plane Electromagnetic Waves

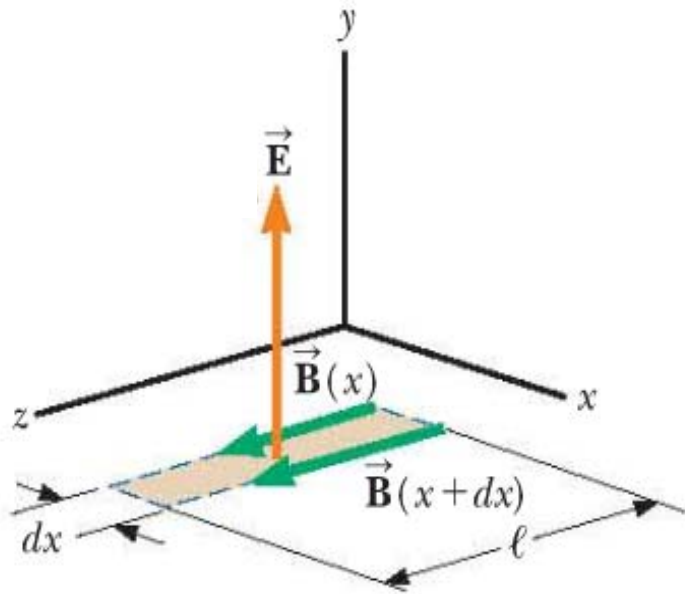


Figure 34.7 At an instant when a plane wave passes through a rectangular path of width dx lying in the xz plane, the magnetic field in the z direction varies from $\vec{B}(x)$ to $\vec{B}(x+dx)$.

$$\oint \vec{B} \cdot d\vec{s} = \mu_0 (I_{\text{in}} + \varepsilon_0 \frac{d}{dt} \oiint \vec{E} \cdot d\vec{A}) \quad (I_{\text{in}} = 0)$$

$$\oint \vec{B} \cdot d\vec{s} = B(x)\ell + 0 - B(x+dx)\ell + 0$$

$$= -\ell[B(x+dx) - B(x)]$$

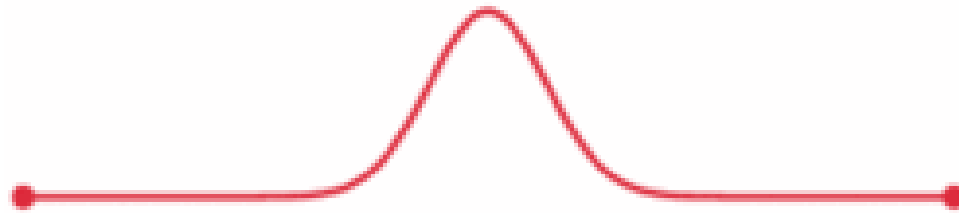
$$= -\ell \frac{\partial B}{\partial x} \cdot dx$$

$$\mu_0 \varepsilon_0 \frac{d}{dt} \oiint \vec{E} \cdot d\vec{A} = \varepsilon_0 \mu_0 \frac{\partial}{\partial t} (E \cdot \ell dx)$$

$$= \varepsilon_0 \mu_0 \ell \frac{\partial E}{\partial t} \cdot dx$$

$$\therefore -\frac{\partial B}{\partial x} = \varepsilon_0 \mu_0 \frac{\partial E}{\partial t}$$

Classical wave equation



$$\frac{1}{v^2} \frac{\partial^2 \psi}{\partial t^2} = \frac{\partial^2 \psi}{\partial x^2}$$

Electric and magnetic fields in wave motion

$$\text{Third equation} \Rightarrow \frac{\partial E}{\partial x} = -\frac{\partial B}{\partial t} \Rightarrow \begin{cases} \frac{\partial^2 E}{\partial x^2} = -\frac{\partial^2 B}{\partial x \partial t} & \text{--- (1a)} \\ \frac{\partial^2 E}{\partial t \partial x} = -\frac{\partial^2 B}{\partial t^2} & \text{--- (1b)} \end{cases}$$

$$\text{Fourth equation} \Rightarrow -\frac{\partial B}{\partial x} = \epsilon_0 \mu_0 \frac{\partial E}{\partial t} \Rightarrow \begin{cases} \frac{\partial^2 B}{\partial x^2} = -\epsilon_0 \mu_0 \frac{\partial^2 E}{\partial x \partial t} & \text{--- (2a)} \\ \frac{\partial^2 B}{\partial t \partial x} = -\epsilon_0 \mu_0 \frac{\partial^2 E}{\partial t^2} & \text{--- (2b)} \end{cases}$$

$$(1a) \text{ and } (2b) \Rightarrow \frac{\partial^2 E}{\partial x^2} = \epsilon_0 \mu_0 \frac{\partial^2 E}{\partial t^2} \Rightarrow \epsilon_0 \mu_0 \frac{\partial^2 E}{\partial t^2} = \frac{\partial^2 E}{\partial x^2}$$

$$(1b) \text{ and } (2a) \Rightarrow -\frac{1}{\epsilon_0 \mu_0} \frac{\partial^2 B}{\partial x^2} = -\frac{\partial^2 B}{\partial t^2} \Rightarrow \epsilon_0 \mu_0 \frac{\partial^2 B}{\partial t^2} = \frac{\partial^2 B}{\partial x^2}$$

Both E and B follow the classical wave equation :

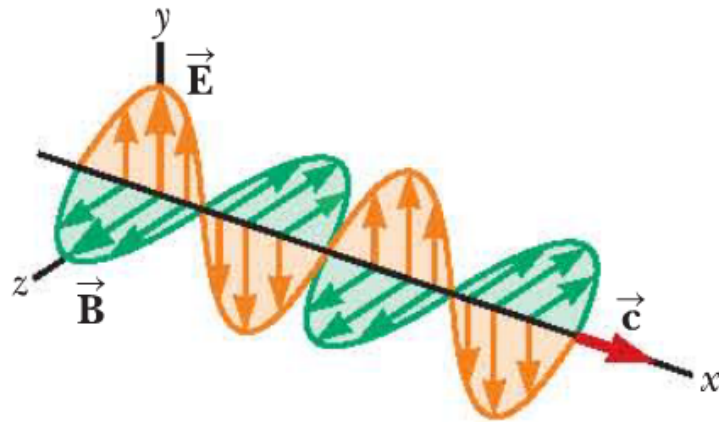
$$\boxed{\frac{\partial^2 \psi}{\partial t^2} = \frac{1}{v^2} \frac{\partial^2 \psi}{\partial x^2}}$$

$$\text{with } v = \frac{1}{\sqrt{\epsilon_0 \mu_0}}$$

$$\epsilon_0 = 8.8542 \times 10^{-12} \text{ C}^2 \text{N}^{-1} \text{m}^{-2}$$

$$\mu_0 = 4\pi \times 10^{-7} \text{ TmA}^{-1}$$

Electromagnetic plane waves



ACTIVE FIGURE 34.8

A sinusoidal electromagnetic wave moves in the positive x direction with a speed c .

$$\frac{1}{c^2} \frac{\partial^2 E}{\partial t^2} = \frac{\partial^2 E}{\partial x^2} \quad \text{and} \quad \frac{1}{c^2} \frac{\partial^2 B}{\partial t^2} = \frac{\partial^2 B}{\partial x^2}$$

Sinusoidal solution :

$$E = E_{\max} \cos(kx - \omega t)$$

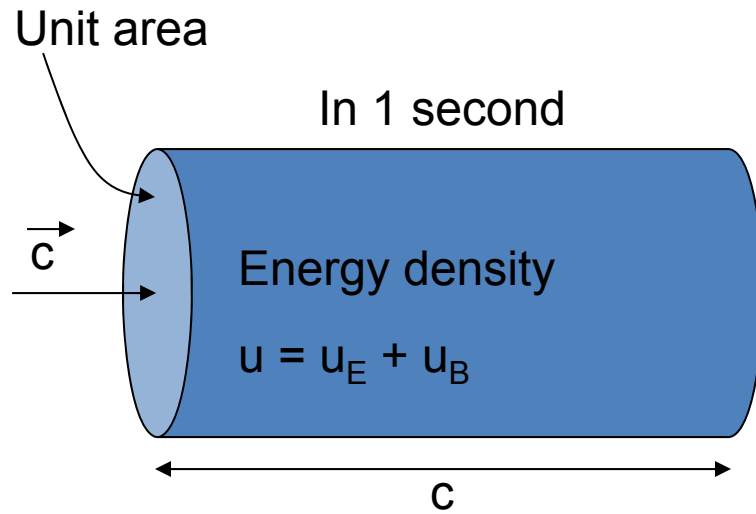
$$B = B_{\max} \cos(kx - \omega t)$$

In here : $k = \frac{2\pi}{\lambda}$ and $\omega = 2\pi f$

$$\omega = ck, \quad E_{\max} = cB_{\max}, \quad \text{and} \quad E = cB$$



$$c = f\lambda$$



Intensity

Intensity I is defined as:

$$I = S_{\text{avg}}$$

From pointing vector:

$$I = S_{\text{avg}}$$

$$\text{But } S = c \left(\frac{1}{\mu_0} B^2 \right) \text{ or } c (\epsilon_0 E^2)$$

$$\therefore I = c \left(\frac{1}{\mu_0} B^2 \right) \text{ or } c (\epsilon_0 E^2)$$

$$= c \left(\frac{1}{\mu_0} \langle B^2 \rangle \right) \text{ or } c (\epsilon_0 \langle E^2 \rangle)$$

$$= c \left(\frac{1}{2\mu_0} B_{\text{max}}^2 \right) \text{ or } c \left(\frac{\epsilon_0}{2} E_{\text{max}}^2 \right)$$

From above picture:

$$I = \langle \text{energy in the blue cylinder} \rangle$$

$$= c (\langle u_E \rangle + \langle u_B \rangle)$$

$$= c \left(\frac{1}{2} \epsilon_0 \langle E^2 \rangle + \frac{1}{2\mu_0} \langle B^2 \rangle \right)$$

$$= c \left(\frac{1}{4} \epsilon_0 E_{\text{max}}^2 + \frac{1}{4\mu_0} B_{\text{max}}^2 \right)$$

$$= c \left(\frac{1}{4} \epsilon_0 (cB)_{\text{max}}^2 + \frac{1}{4\mu_0} B_{\text{max}}^2 \right) \text{ or } c \left(\frac{1}{4} \epsilon_0 E_{\text{max}}^2 + \frac{1}{4\mu_0} \left(\frac{E_{\text{max}}}{c} \right)^2 \right)$$

$$= c \left(\frac{1}{4\mu_0} B_{\text{max}}^2 + \frac{1}{4\mu_0} B_{\text{max}}^2 \right) \text{ or } c \left(\frac{1}{4} \epsilon_0 E_{\text{max}}^2 + \frac{1}{4} \epsilon_0 E_{\text{max}}^2 \right)$$

$$= c \left(\frac{1}{2\mu_0} B_{\text{max}}^2 \right) \text{ or } c \left(\frac{1}{2} \epsilon_0 E_{\text{max}}^2 \right)$$

Poynting vector

At any point, knowing \vec{E} and \vec{B} , we can define Poynting vector \vec{S} as:

$$\vec{S} = \frac{1}{\mu_0} \vec{E} \times \vec{B}$$

Poynting vector gives the energy passes through a unit surface area perpendicular to the direction of wave propagation. \vec{S} is along the direction of wave propagation and has unit W/m².

For plane wave:

$$\begin{aligned} \vec{S} = \frac{1}{\mu_0} \vec{E} \times \vec{B} &\Rightarrow S = \frac{EB}{\mu_0} = \frac{c}{\mu_0} B^2 \text{ or } \frac{1}{c\mu_0} E^2 \\ &\Rightarrow S = c\left(\frac{1}{\mu_0} B^2\right) \text{ or } c(\epsilon_0 E^2) \end{aligned}$$

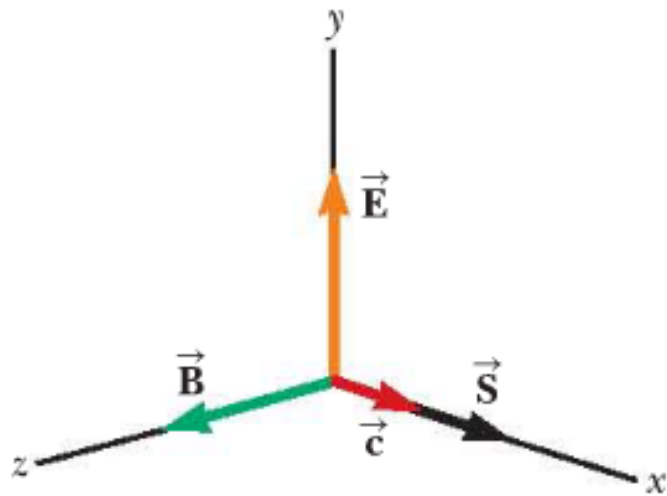
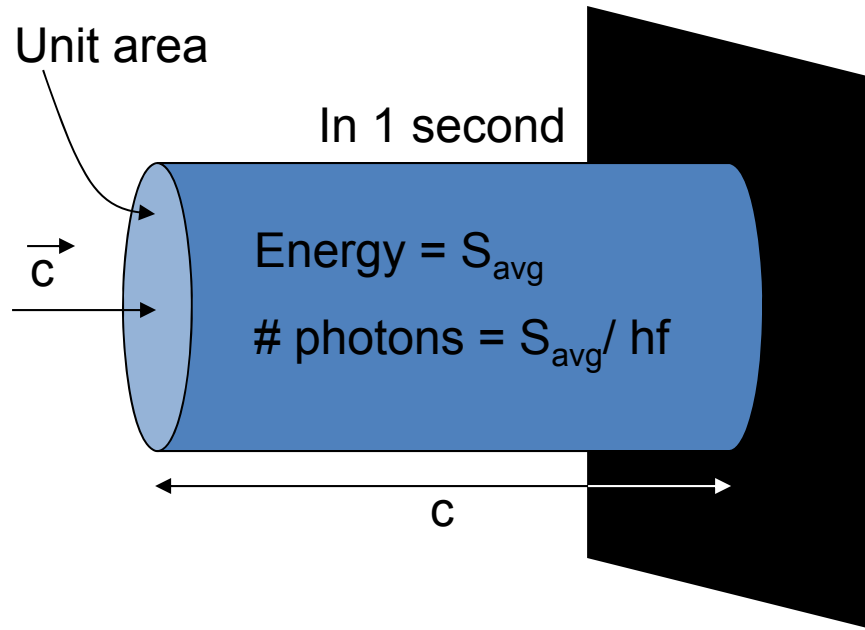


Figure 34.10 The Poynting vector \vec{S} for a plane electromagnetic wave is along the direction of wave propagation.

Momentum



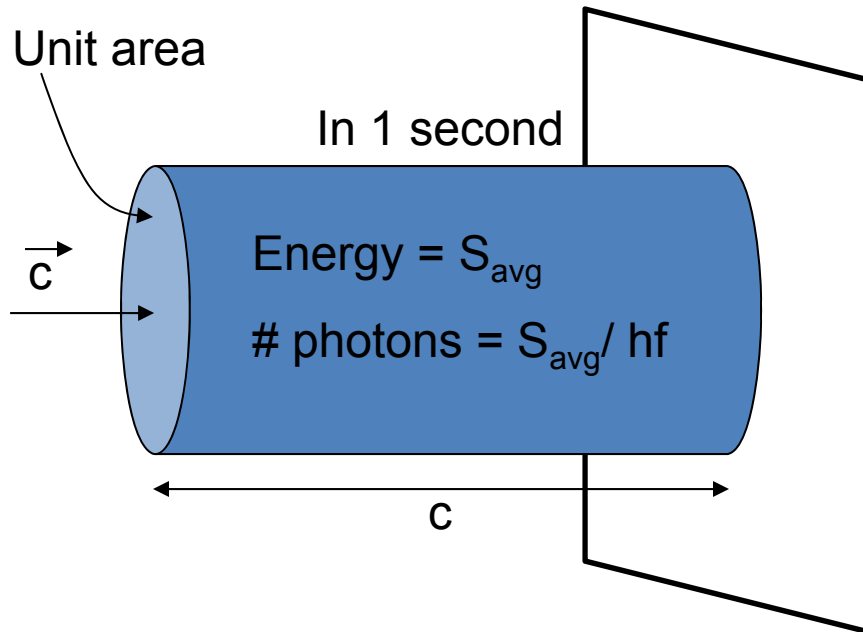
If all these photons are absorbed by the surface:

$$\text{In one second, } \Delta p = \frac{S_{\text{avg}}}{c}$$

$$\therefore F_{\text{unit area}} = \frac{\Delta p}{\Delta t} = \frac{S_{\text{avg}}}{c}$$

But this is on unit area

$$\therefore P = \frac{S_{\text{avg}}}{c}$$



If all these photons are reflected by the surface:

$$\text{In one second, } \Delta p = 2 \cdot \frac{S_{\text{avg}}}{c}$$

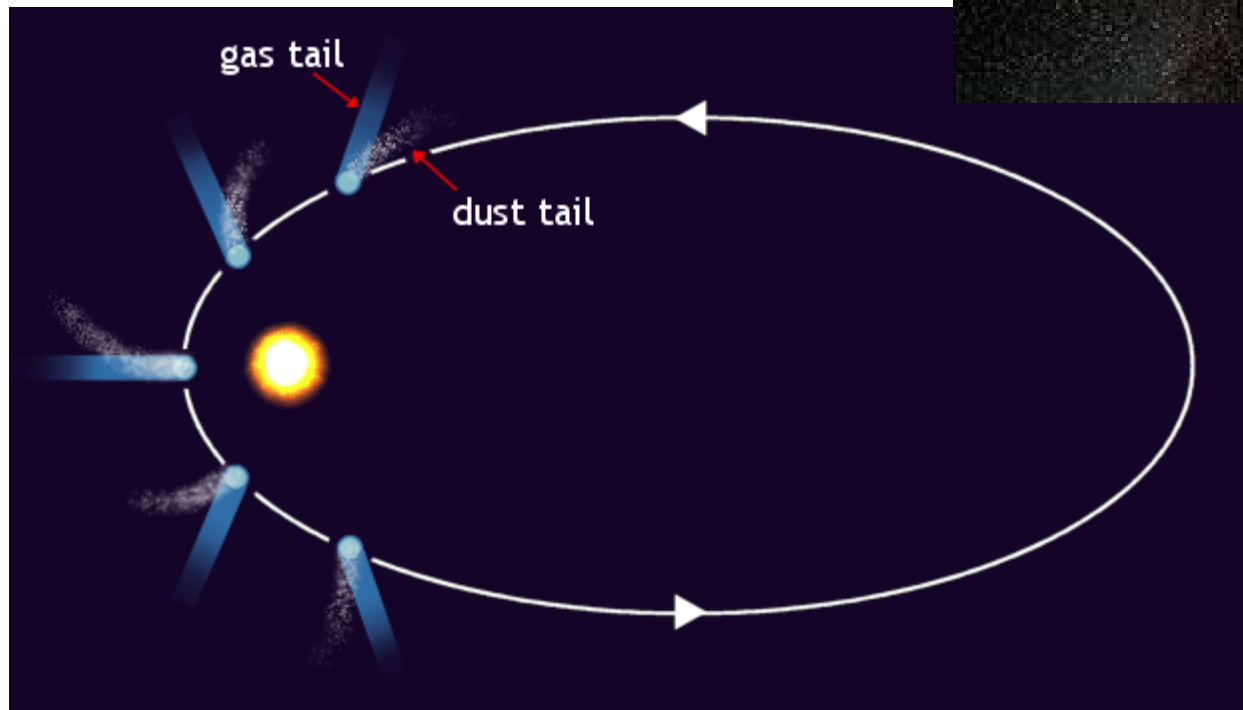
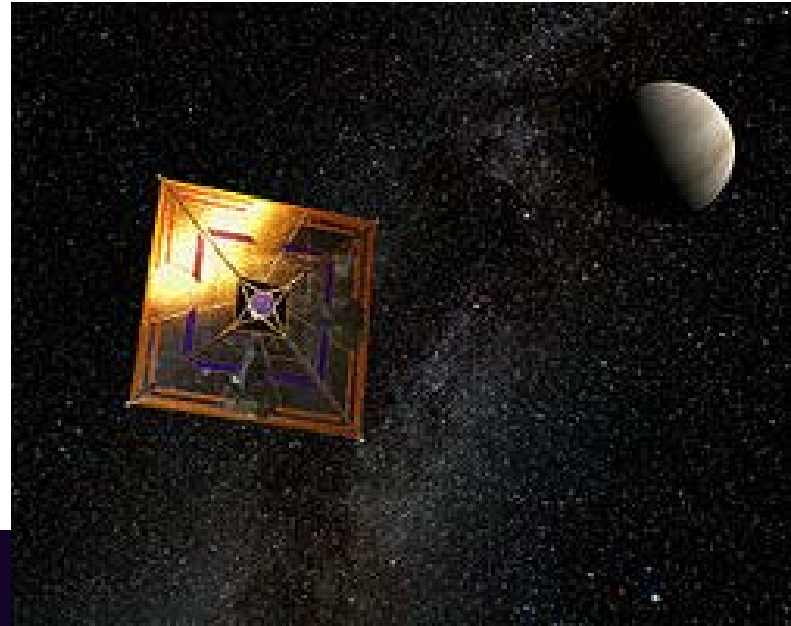
$$\therefore F_{\text{unit area}} = \frac{\Delta p}{\Delta t} = \frac{2S_{\text{avg}}}{c}$$

But this is on unit area

$$\therefore P = \frac{2S_{\text{avg}}}{c}$$

Radiation Pressure

IKAROS



Comet

Electromagnetic spectrum

